



Automation of Turbulence Simulation

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Some References

- * J.Hoffman, Computation of mean drag for bluff body problems using Adaptive DNS/LES, *SIAM J. Sci. Comput. Vol.27(1)*, 2005.
- * J.Hoffman and C.Johnson, A new approach to Computational Turbulence Modeling, to appear in *Comput. Methods Appl. Mech. Engrg.*, 2005.
- * J.Hoffman, Simulation of turbulent flow past bluff bodies on coarse meshes using General Galerkin methods: drag crisis and turbulent Euler solutions, to appear in *Springer Journal of Computational Mechanics*, 2005.
- * J.Hoffman, Efficient computation of mean drag for the subcritical flow past a circular cylinder using Adaptive DNS/LES, in review (*Int. J. of Numerical Methods in Fluids*) .
- * J.Hoffman, Adaptive simulation of the turbulent flow due to a cylinder rolling along ground, in review (*Comput. Methods Appl. Mech. Engrg.*).
- * J.Hoffman, Adaptive simulation of the turbulent flow past a sphere, in review (*J. Fluid Mech.*).
- * J.Hoffman and C.Johnson, Adaptive FEM for Incompressible Turbulent Flow, Springer, 2006.

Automation of Turbulence Simulation

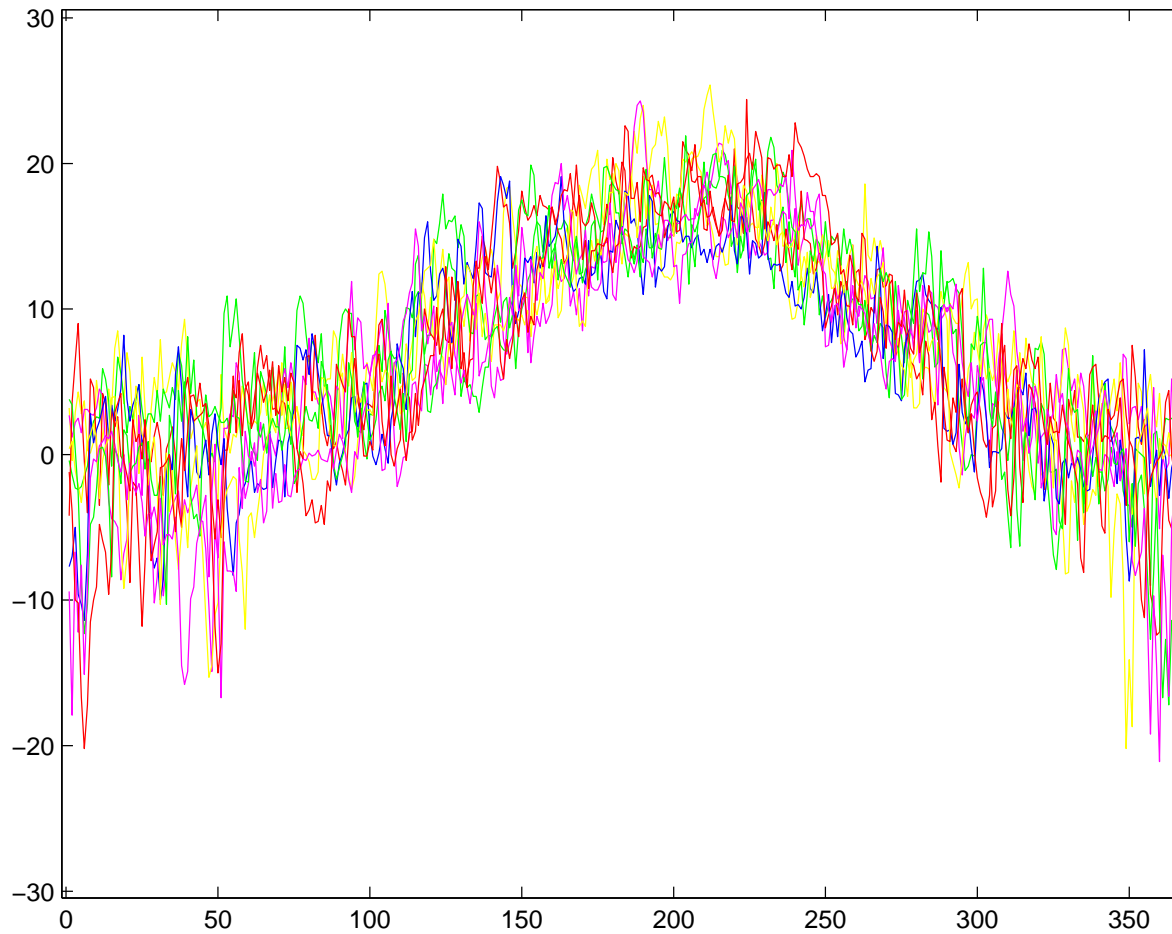
- ACMM: Generality in: Method and Implementation
- Method: Adaptive stabilized FEM: General Galerkin G2
- Implementation: FEniCS project

This talk: Test of Method: G2

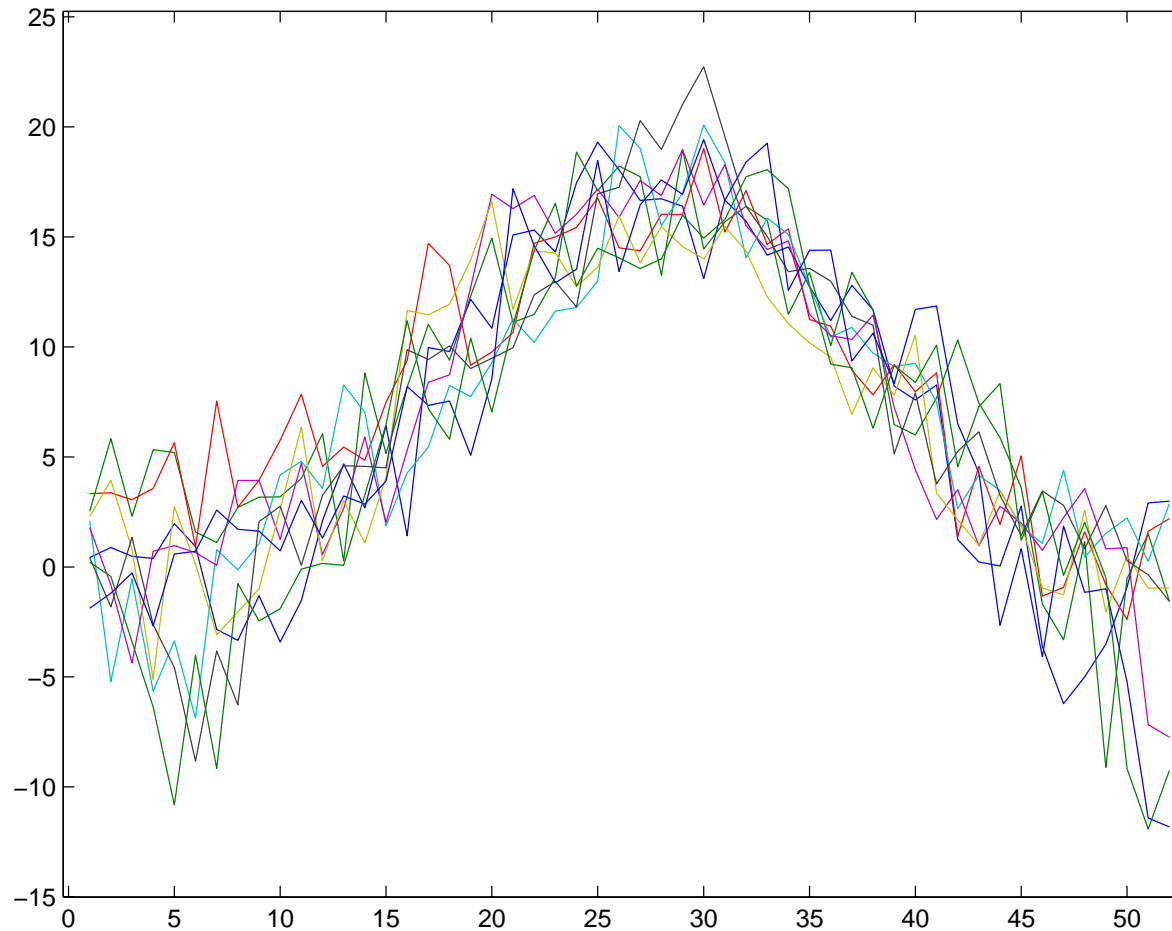
Turbulence Simulation; characterized by non-generality:

- discretization (geometry specific FD, spectral,...)
- turbulence modeling (parameters depend on data, numerics,...; different filters used in RANS/LES,...)
- wall modeling using RANS, boundary layer theory,...

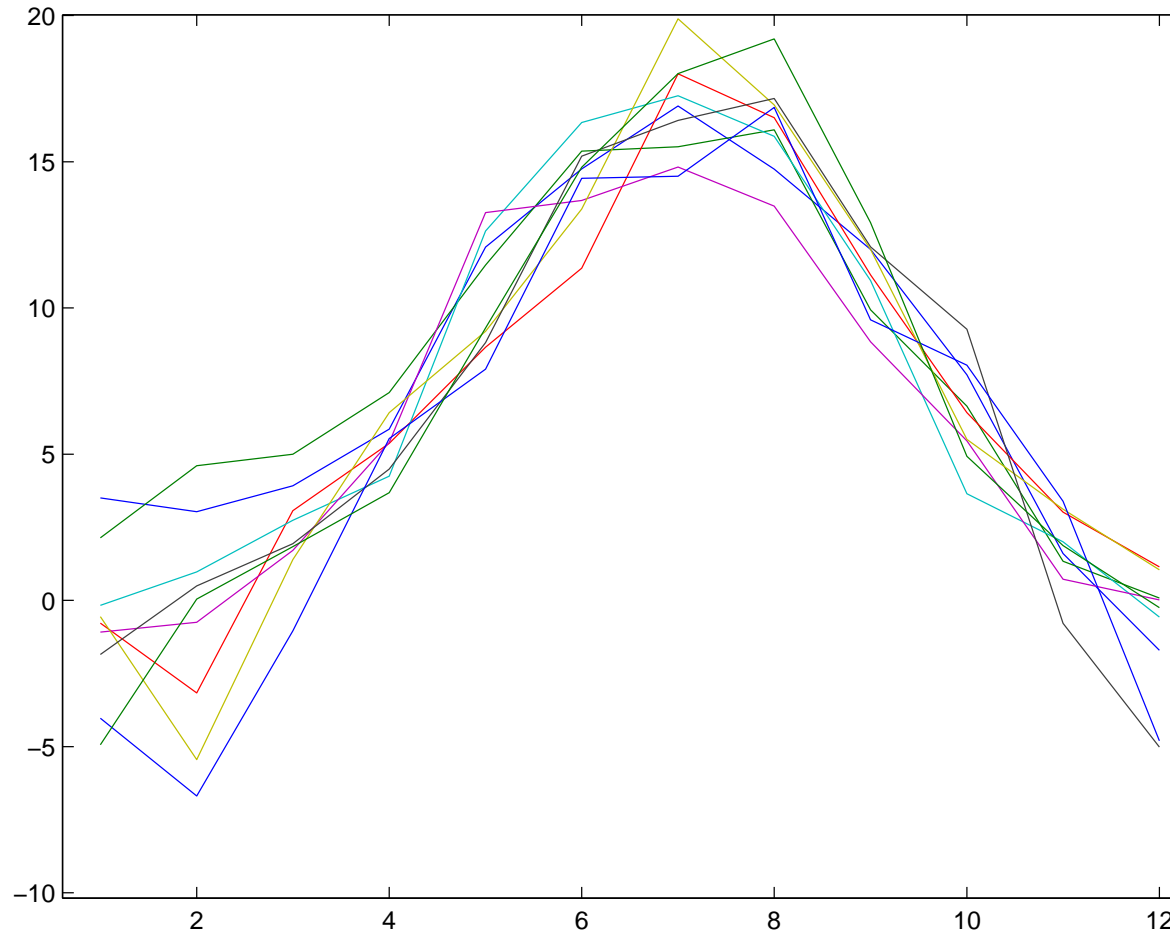
Målilla 88-95: 24-hr mean: $\pm 10^{\circ}\text{C}$



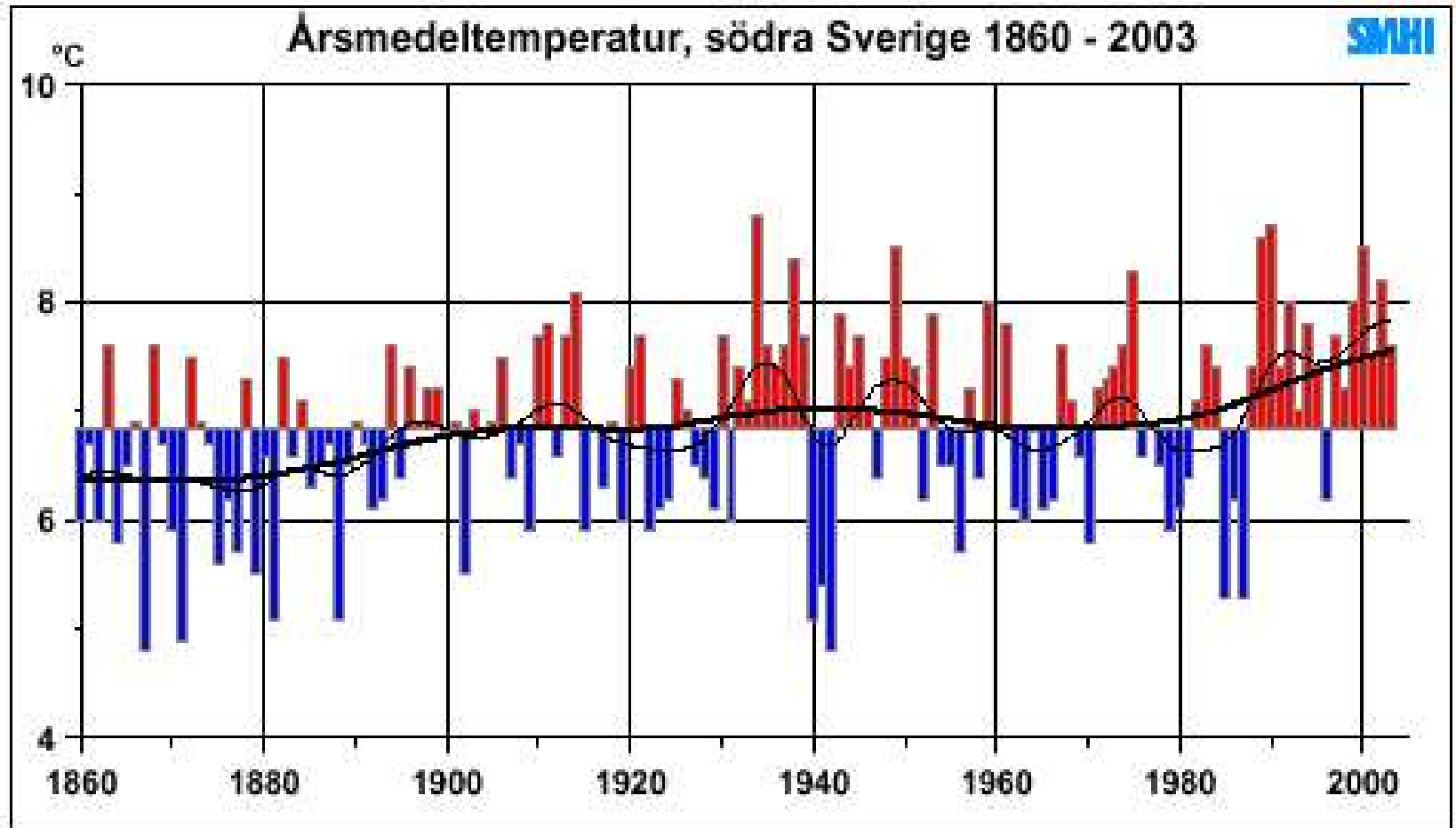
Målilla 88-95: weekly mean: $\pm 5^\circ\text{C}$



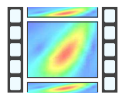
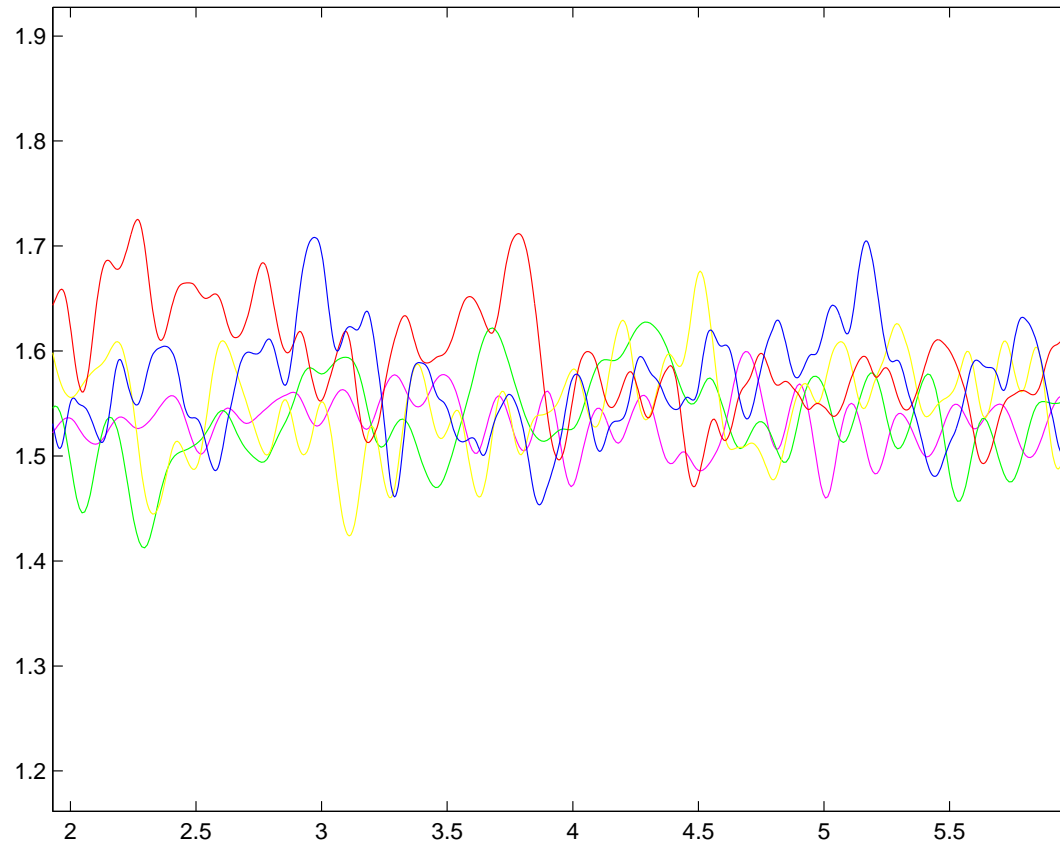
Målilla 88-95: monthly mean: $\pm 3^{\circ}\text{C}$



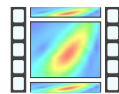
Yearly mean 1860-2003: $\pm 1^\circ\text{C}$



Cube: $c_D(t) \pm 10\%$, mean $c_D \pm 2\%$



cube: xy-plane



cube: xz-plane

Turbulence: Chaotic Dyn. System

- Pointwise quantities strongly sensitive to perturbations and thus unpredictable (to any tolerance of interest)
- Mean values in space/time moderately sensitive to perturbations and thus predictable up to a tolerance
- Ex: Weather prediction (Målilla):
24-hr mean temperature unpredictable ($\pm 10^\circ\text{C}$)
monthly mean temperature predictable ($\pm 3^\circ\text{C}$)
- Ex: Bluff body flow (experiments/computations):
pointwise drag $c_D(t)$ unpredictable ($> \pm 10\%$)
mean value drag $\int_I c_D(t) dt$ predictable ($\pm 2\%$)

Navier-Stokes Equations (NSE)

$$R(\hat{u}) = \begin{pmatrix} \dot{u} + u \cdot \nabla u - \nu \Delta u + \nabla p - f \\ \nabla \cdot u \end{pmatrix} = 0$$

$\hat{u} = (u, p)$: u velocity, p pressure, ν viscosity, $f = 0$

■ Pointwise existence & uniqueness unknown: $R(\hat{u}) = 0$
(Clay Institute \$1 million Prize Problem)

■ Existence (but not uniq.) of weak solution (Leray 1934):

Find $\hat{u} \in \hat{V}$: $(R(\hat{u}), \hat{v}) = 0 \quad \forall \hat{v} = (v, q) \in \hat{V}$

$\hat{V} \subset H^1(Q)$: $Q = \Omega \times I$ space-time domain, $(\cdot, \cdot) = (\cdot, \cdot)_Q$

$(R(\hat{u}), \hat{v}) \equiv (\dot{u}, v) + (u \cdot \nabla u, v) - (\nabla \cdot v, p) + (\nabla \cdot u, q) + (\nu \nabla u, \nabla v)$

Navier-Stokes Equations (NSE)

$$R(\hat{u}) = \begin{pmatrix} \dot{u} + u \cdot \nabla u - \nu \Delta u + \nabla p - f \\ \nabla \cdot u \end{pmatrix} = 0$$

$\hat{u} = (u, p)$: u velocity, p pressure, ν viscosity, $f = 0$

- Computational turbulence: Mean values computable ($< \pm 5\%$), point values not computable ($> \pm 10\%$).
- Computational cost (#dofs): $\sim Re^3$ ($Re = UL/\nu$)
(resolving all scales in Direct Num. Simulation DNS)

Limit today: $Re \approx 10^3$ (many industrial appl. $Re > 10^6$)

Analysis & Computation of NSE

Scientific goals today:

- Prove exist & uniq of pointwise NSE: $R(\hat{u}) = 0$
- Push limit of DNS (wrt Re^3 constraint)
- Turbulence modeling: find model for unresolved scales (filtering of NSE: Reynolds stresses, closure problem)

Alternative scientific goals:

- Approximate weak solution \hat{U} : weak uniqueness in (mean value) output $M(\hat{U})$: stability of \hat{U} wrt $M(\cdot)$
- Adaptive algorithm: min #dof : $error(M(\hat{U})) < TOL$

Existence of ϵ -Weak Solutions

- $W_\epsilon = \{\hat{u} \in \hat{V} : |(R(\hat{u}), \hat{v})| \leq \epsilon \|\hat{v}\|_{\hat{V}} \quad \forall \hat{v} \in \hat{V}\}$

(approximate weak solution: $\sim \|R(\hat{u})\|_{H^{-1}} \leq \epsilon$)

- Existence: Construction of W_ϵ : General Galerkin G2

- Find $\hat{U} \in \hat{V}_h$: $(R(\hat{U}), \hat{v}) + (hR(\hat{U}), R(\hat{v})) = 0 \quad \forall \hat{v} \in \hat{V}_h$

$$\frac{1}{2} \|U(t)\|^2 + \|\sqrt{\nu} \nabla U\|_Q^2 + \|\sqrt{h} R(\hat{U})\|_Q^2 = \frac{1}{2} \|U(0)\|^2$$

$$\begin{aligned} (R(\hat{U}), \hat{v}) &= (R(\hat{U}), \hat{v} - \pi_h \hat{v}) - (hR(\hat{U}), R(\pi_h \hat{v})) \\ &\leq (C + M_U) \|hR(\hat{U})\|_Q \|\hat{v}\|_{\hat{V}} \leq C\sqrt{h} \|\hat{v}\|_{\hat{V}} \end{aligned}$$

- $\hat{U} \in \mathbf{G2} \Rightarrow \hat{U} \in W_\epsilon \quad \epsilon = (C + M_U) \|hR(\hat{U})\|_Q$

Weak Uniqueness: Duality

■ Output (functional): $M(\hat{u}) \equiv (\hat{u}, \hat{\psi})$

■ Dual NSE: Find $\hat{\varphi} = (\varphi, \theta)$: $a(\hat{u}, \hat{w}; \hat{v}, \hat{\varphi}) = M(\hat{v}) \quad \forall \hat{v} \in \hat{V}_0$
 $v = (v, q) \in \hat{V}_0 \quad \hat{V}_0 = \{\hat{v} \in \hat{V} : v(\cdot, 0) = 0\}$

$$a(\hat{u}, \hat{w}; \hat{v}, \hat{\varphi}) \equiv (\dot{v}, \varphi) + (u \cdot \nabla v, \varphi) + (v \cdot \nabla w, \varphi) \\ - (\nabla \cdot \varphi, q) + (\nabla \cdot v, \theta) + (\nu \nabla v, \nabla \varphi)$$

■ Weak Uniqueness: $\hat{u}, \hat{w} \in W_\epsilon$, $S_\epsilon(\hat{\psi}) \equiv \max_{\hat{u}, \hat{w} \in W_\epsilon} \|\hat{\varphi}\|_{\hat{V}}$

$$|M(\hat{u}) - M(\hat{w})| = |a(\hat{u}, \hat{w}; \hat{u} - \hat{w}, \hat{\varphi})| \\ = |(R(\hat{u}), \hat{\varphi}) - (R(\hat{w}), \hat{\varphi})| \leq 2\epsilon S_\epsilon(\hat{\psi})$$

■ Exact solution ($\epsilon = 0$): stability information lost!

Weak Uniqueness: Duality

■ Weak Uniqueness: $\hat{u}, \hat{w} \in W_\epsilon$

$$|M(\hat{u}) - M(\hat{w})| \leq 2\epsilon S_\epsilon(\hat{\psi}) \quad \|R(\hat{u})\|_{H^{-1}}, \|R(\hat{w})\|_{H^{-1}} \leq \epsilon$$

■ Computability: $\hat{u} \in W_\epsilon$ and $\hat{U} \in G2$

$$|M(\hat{u}) - M(\hat{U})| \leq (\epsilon + \epsilon_{G2}) S_{\epsilon_{G2}}(\hat{\psi}) \quad \epsilon_{G2} = C \|hR(\hat{U})\|$$

■ Residual only needs to be small in a weak norm!!!
(for weak uniqueness)

■ $\|R(\hat{u})\|_{H^{-1}}$ & $\|hR(\hat{U})\|_{L_2}$ vs $\|R(\hat{u})\|_{L_2}$

■ Weak uniqueness characterized by stability factor $S_\epsilon(\hat{\psi})$

G2 for NSE: Adaptive DNS/LES

NSE: $R(\hat{u}) = 0$

G2: $\hat{U} \in \hat{V}_h : (R(\hat{U}), \hat{v}) + SD_\delta(\hat{U}; \hat{v}) = 0 \quad \forall \hat{v} \in \hat{V}_h$

Functional output: $M(\hat{u}) = (\hat{u}, \hat{\psi})$ (drag, lift,...)

Dual NSE: $a(\hat{u}, \hat{w}; \hat{v}, \hat{\varphi}) = M(\hat{v}) \quad \forall \hat{v} \in \hat{V}_0$

$$a(\hat{u}, \hat{w}; \hat{v}, \hat{\varphi}) \equiv (\dot{v}, \varphi) + (u \cdot \nabla v, \varphi) + (v \cdot \nabla w, \varphi) \\ - (\nabla \cdot \varphi, q) + (\nabla \cdot v, \theta) + (\nu \nabla v, \nabla \varphi)$$

Error identity (using duality and G2):

$$|M(\hat{u}) - M(\hat{U})| = |(R(\hat{U}), \hat{\varphi} - \hat{\Phi}) + SD_\delta(\hat{U}; \hat{\Phi})| \quad \forall \hat{\Phi} \in \hat{V}_h$$

G2 for NSE: Adaptive DNS/LES

$$|M(\hat{u}) - M(\hat{U})| \leq \sum_{K \in \mathcal{T}} \mathcal{E}_K = \sum_{K \in \mathcal{T}} \|hR(\hat{U})\| \|\hat{\varphi}_h\|_1 + |SD_\delta(\hat{U}; \hat{\varphi}_h)|$$

Galerkin discretization error + stabilization modeling error

Adaptive algorithm: From coarse mesh \mathcal{T}^0 do

(1) compute primal and dual problem on \mathcal{T}^k

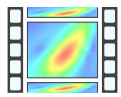
(2) if $\sum_{K \in \mathcal{T}^k} \mathcal{E}_K^k < \text{TOL}$ then STOP, else

(3) refine elements $K \in \mathcal{T}^k$ with largest $\mathcal{E}_K^k \rightarrow \mathcal{T}^{k+1}$

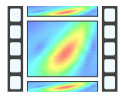
(4) set $k = k + 1$, then goto (1)

Bluff Body Benchmark Problems

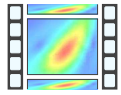
Drag Coeff. c_D = normalized mean drag force



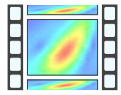
circular cylinder $Re=3900$



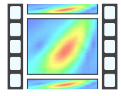
dual solution



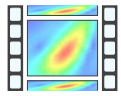
sphere $Re=10\ 000$



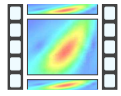
dual solution



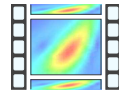
square cylinder $Re=22\ 000$



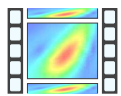
dual solution



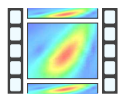
cube $Re=40\ 0000$: xy



cube xz

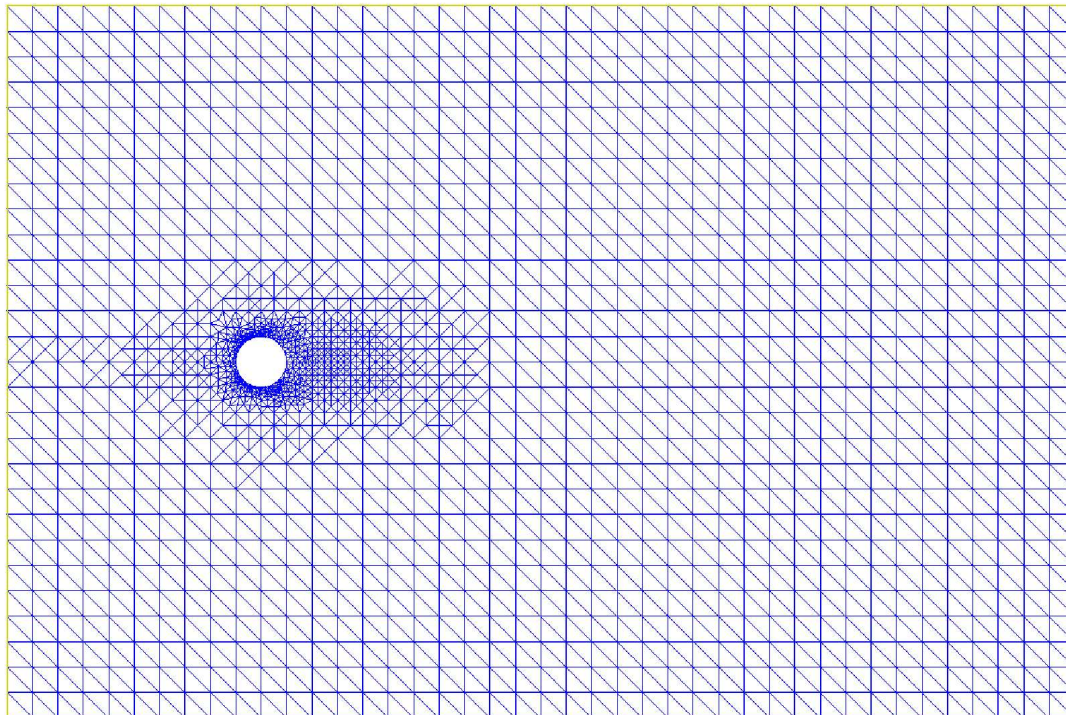
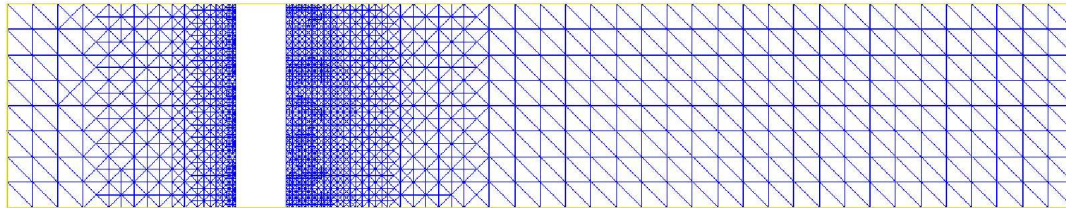


dual sol. xy

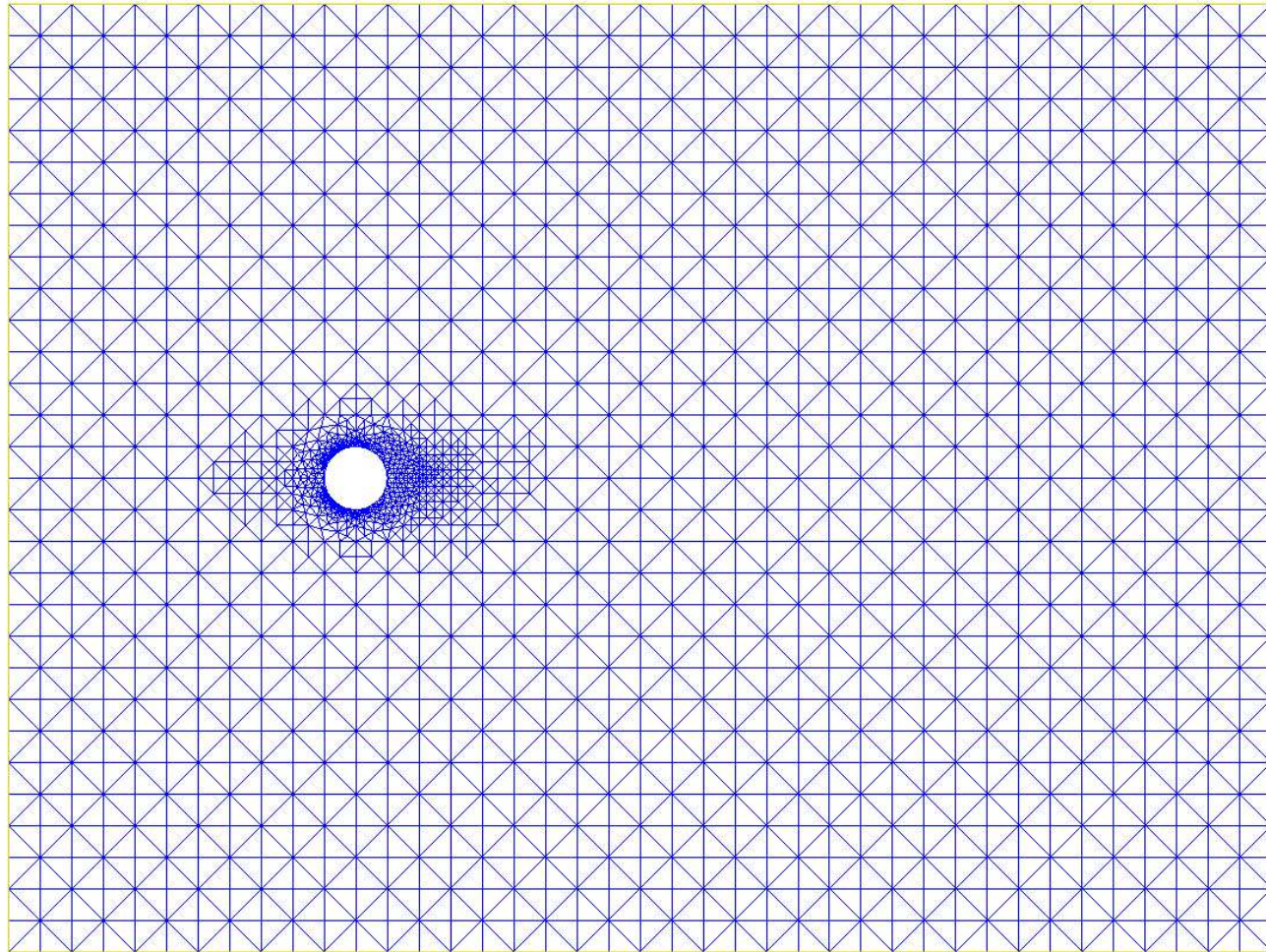


dual sol. xz

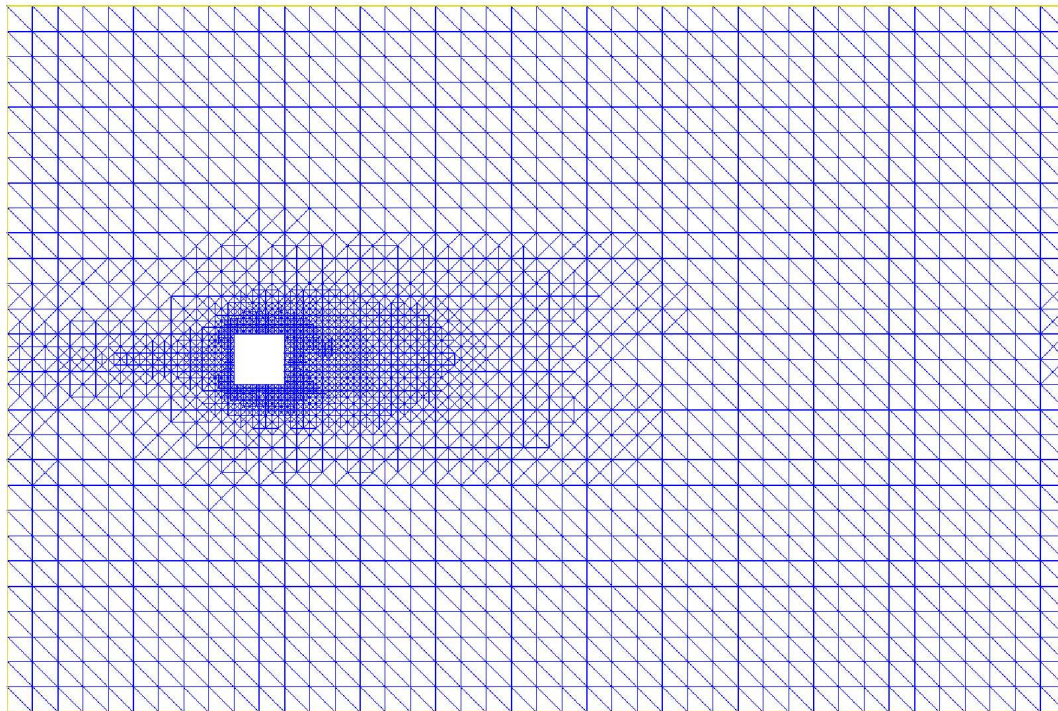
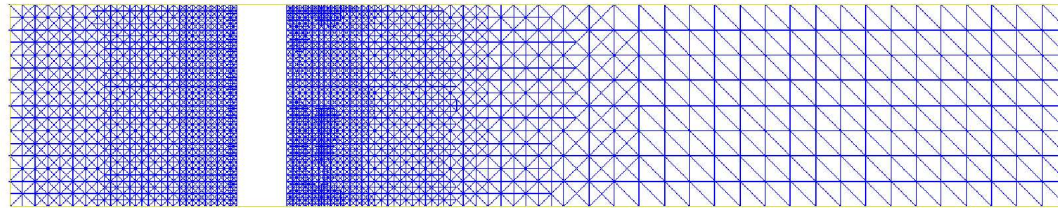
Ref. Mesh wrt C_D : circular cylinder



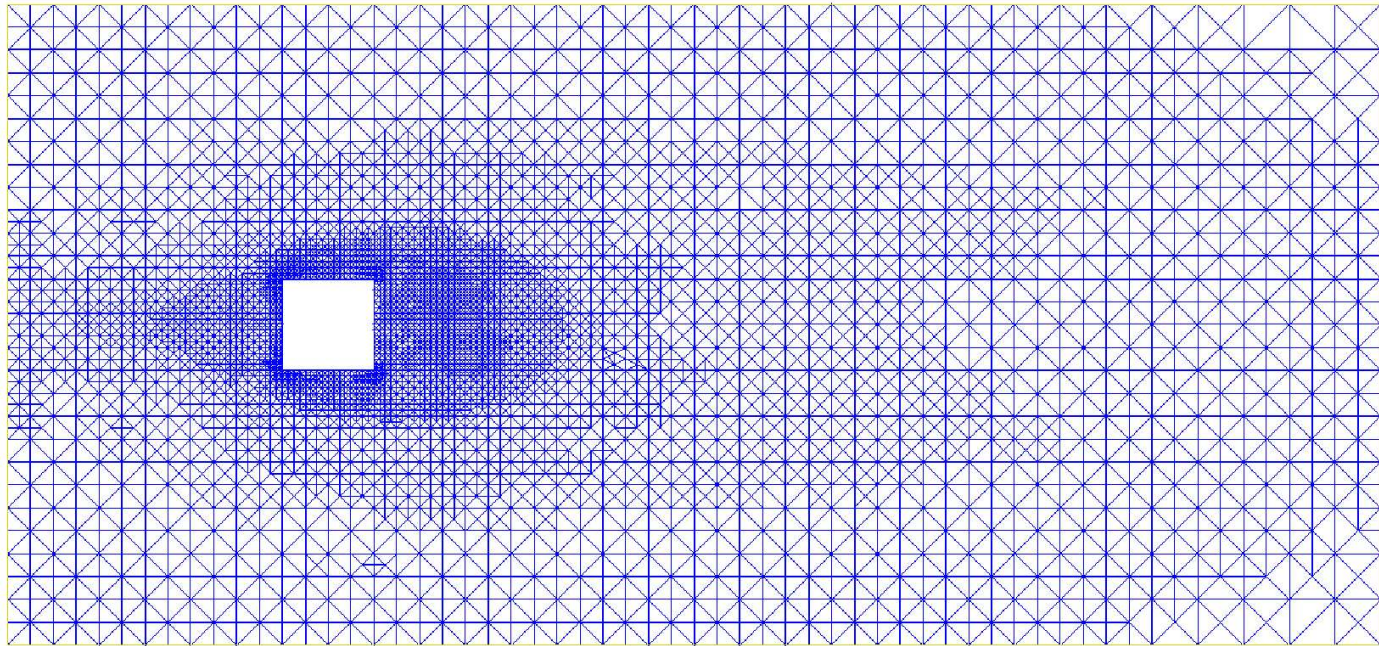
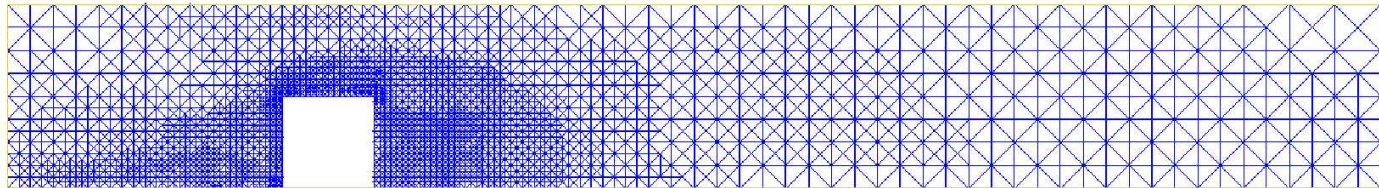
Ref. Mesh wrt C_D : sphere



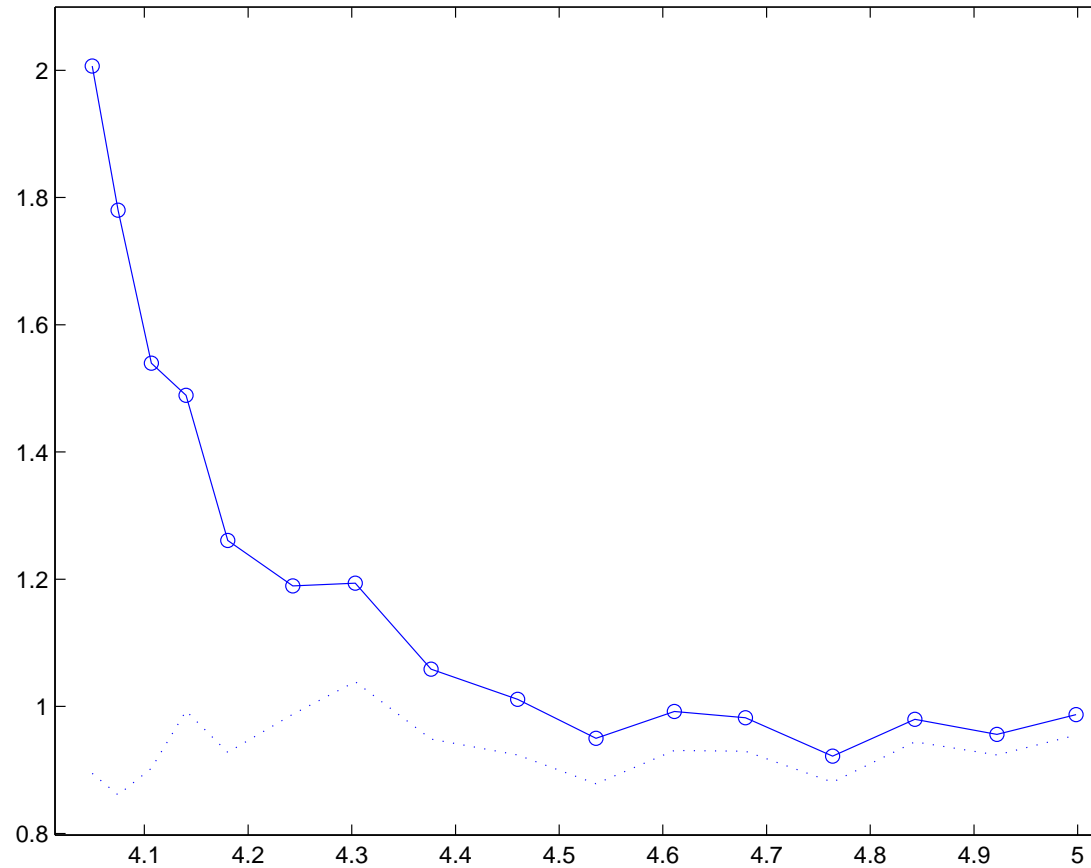
Ref. Mesh wrt C_D : square cylinder



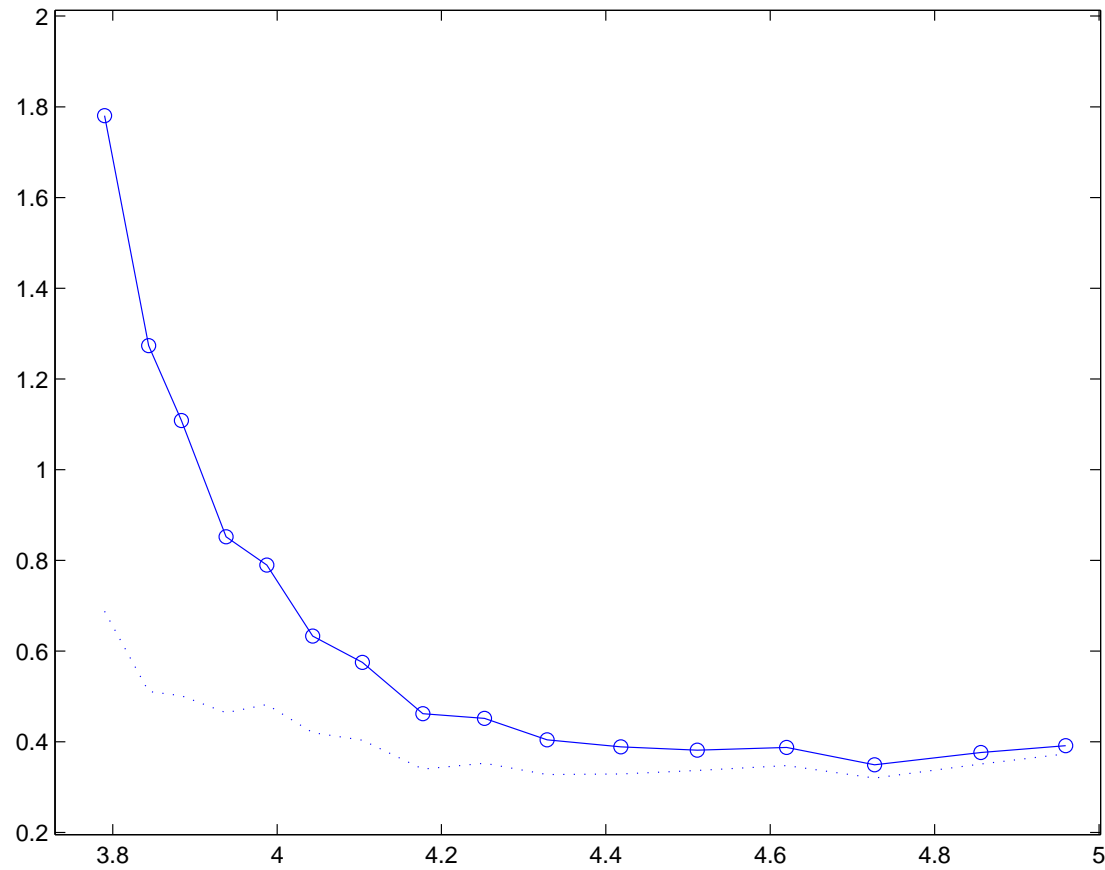
Ref. Mesh wrt C_D : surf mount cube



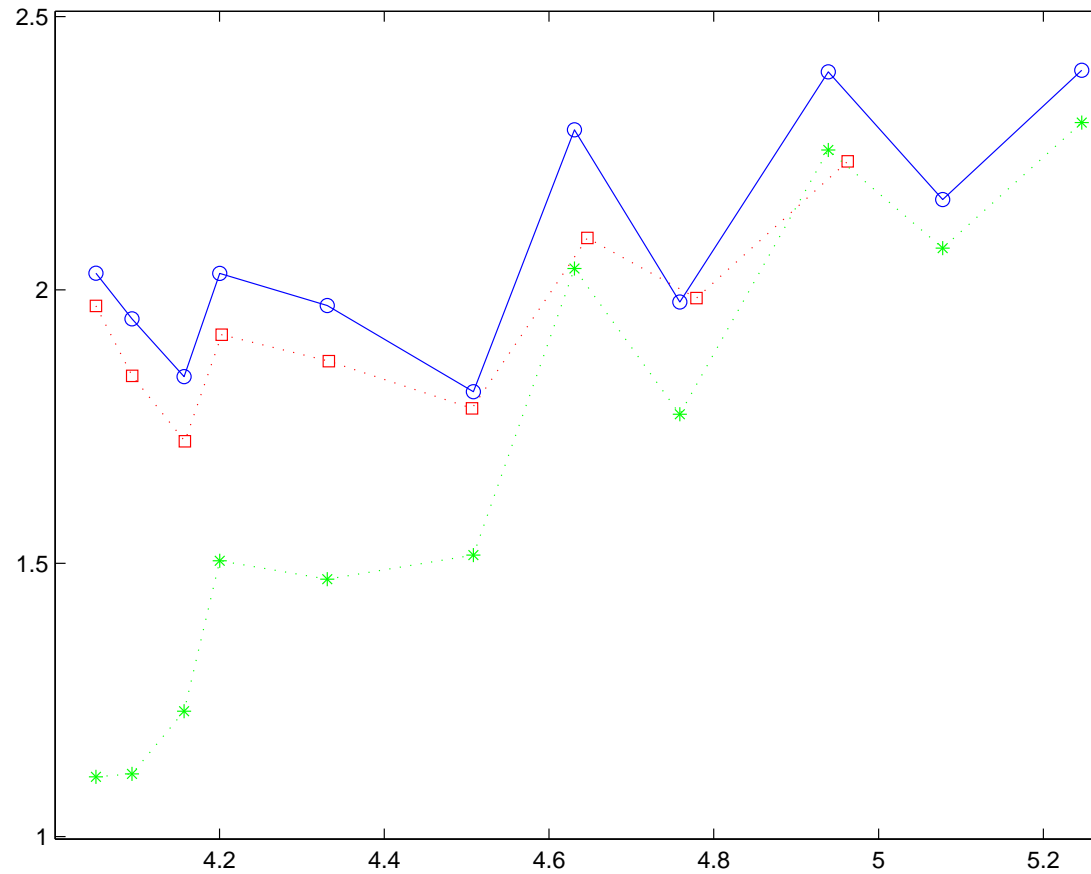
Circular cylinder: $c_D \approx 1.0$



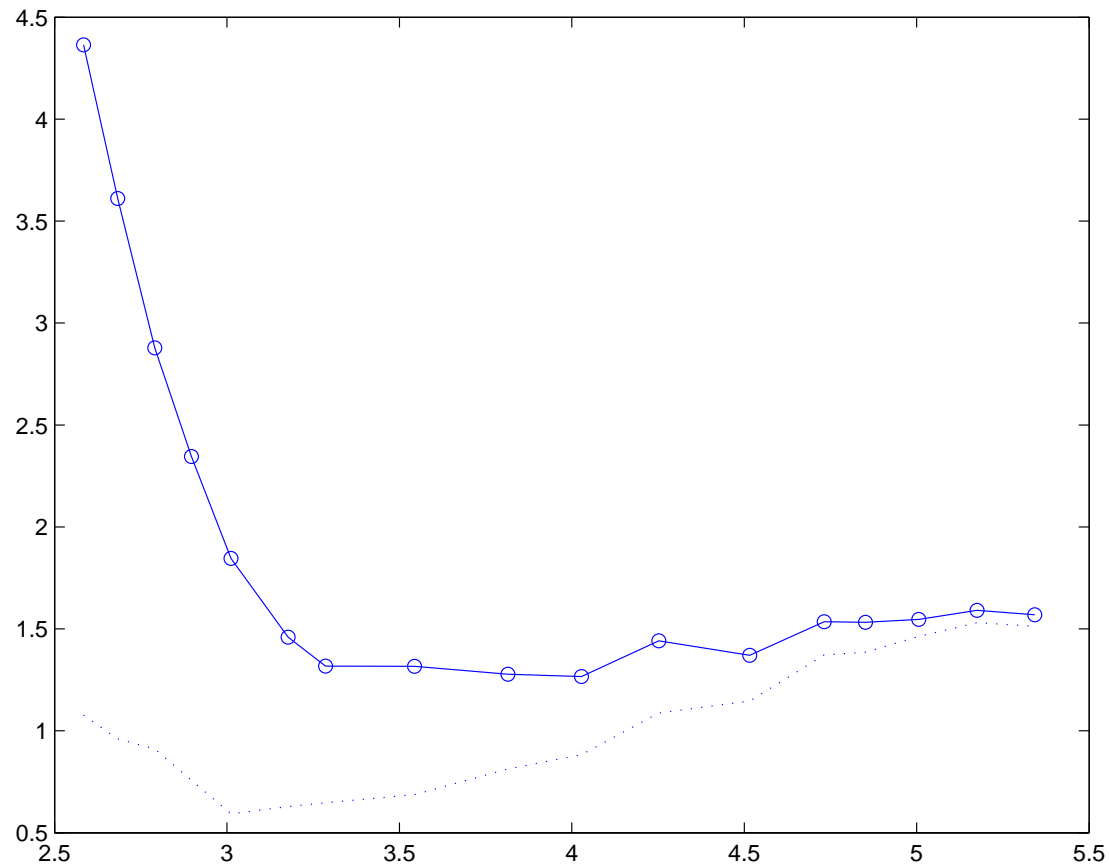
Sphere: $c_D \approx 0.4$



Square cylinder: $c_D \approx 2.2$



Surface mounted cube: $c_D \approx 1.5$

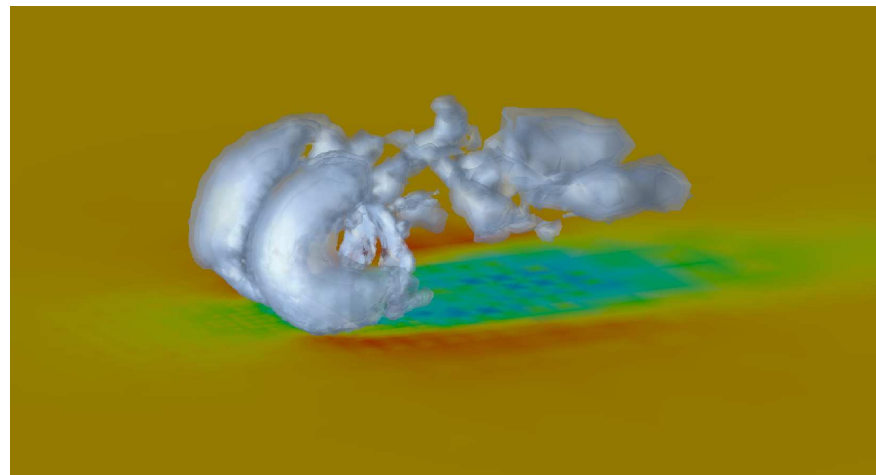


Summary Benchmark Problems

- G2 for NSE: No filtering. No Reynolds stresses. G2 automatic “turbulence model”.
- Adaptive algorithm captures separation points, and “correct” (finite limit) dissipation in the turbulent wake.
- Mean value output (drag, lift, frequencies, separation points, pressure coeff,...) computable up to a tolerance corresponding to experimental accuracy ($\approx 1-5\%$).
- About 10-100 times less mesh points needed to compute drag than in non-adaptive LES.
- All computations on a standard PC (2 GHz, 0.5Gb)

Ex: Rotating Cylinder $Re = 10^4$

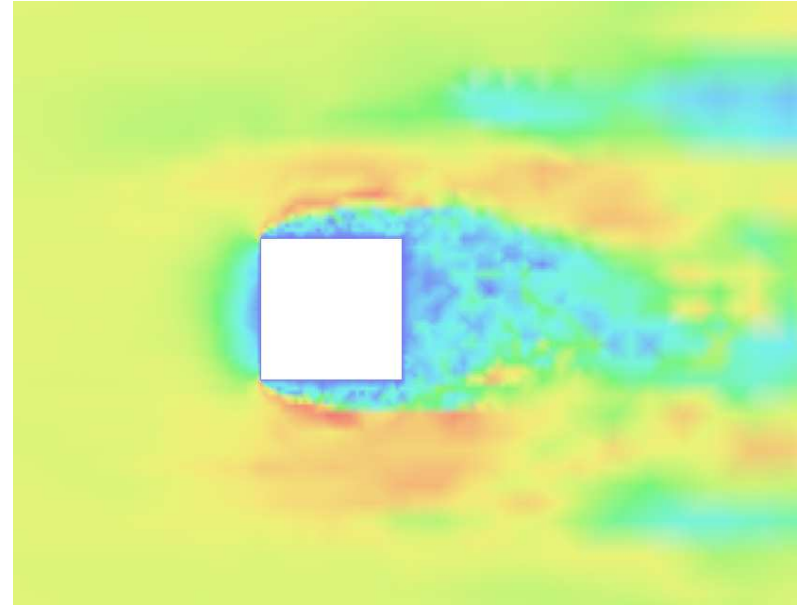
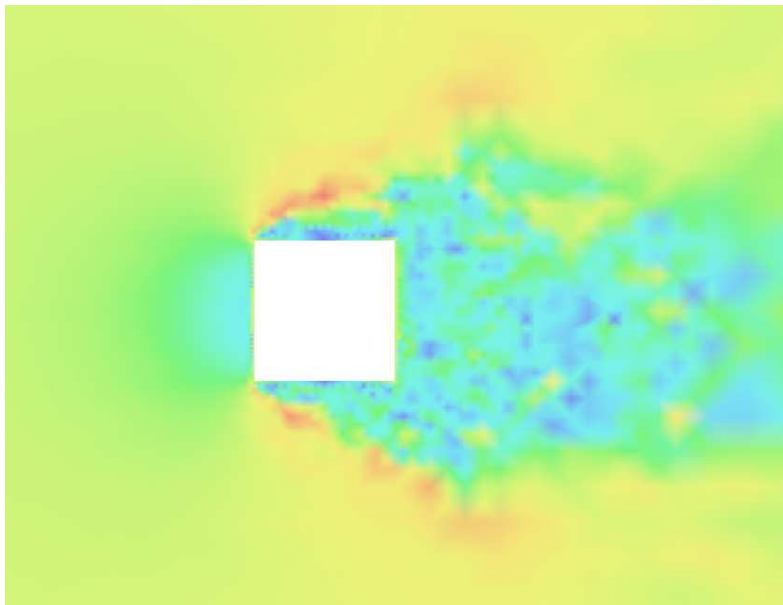
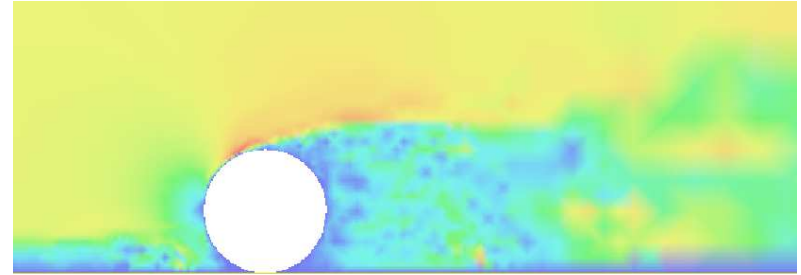
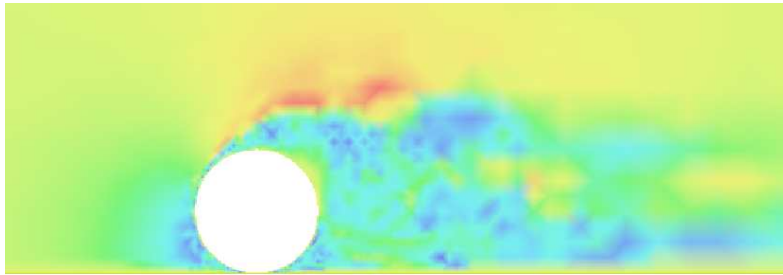
Rotating cylinder, ground moving at free stream speed



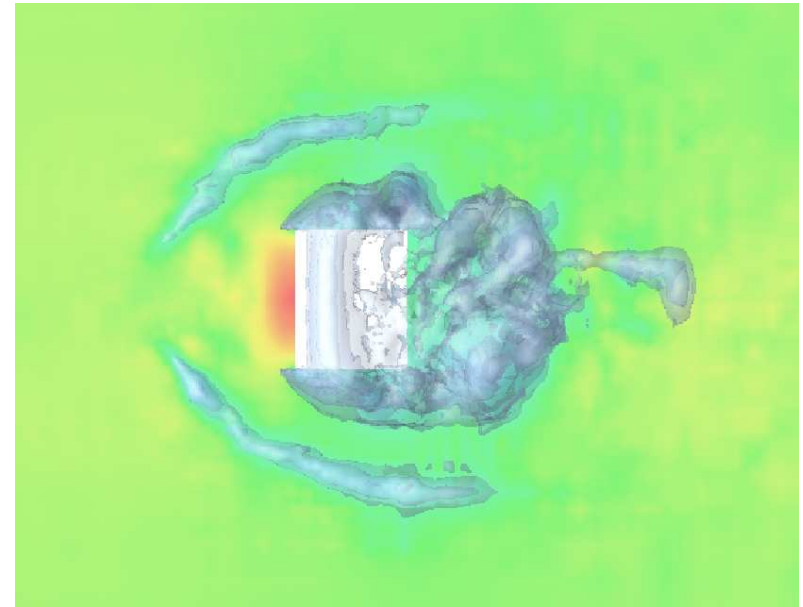
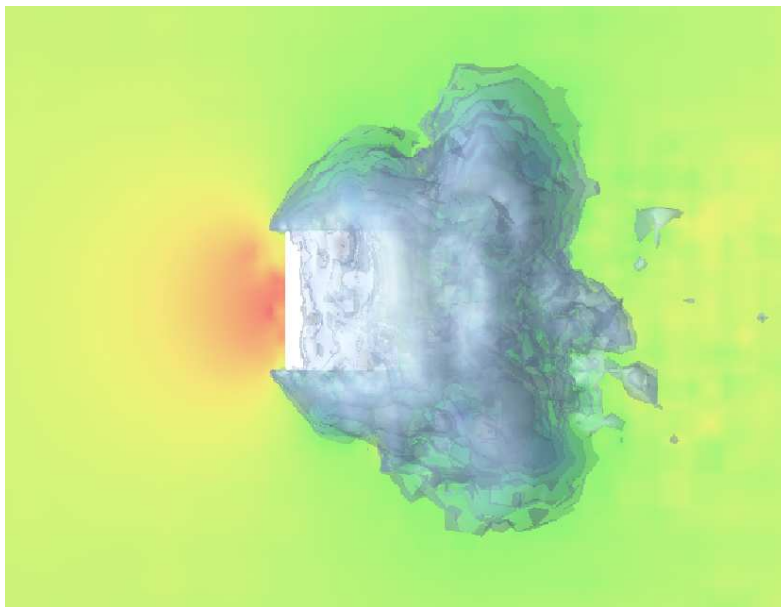
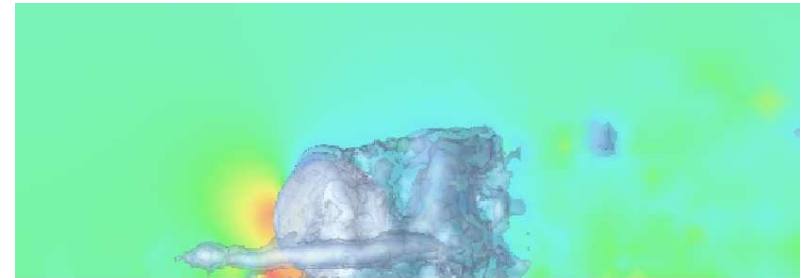
Model of rotating cyl on ground: moving coordinate frame
(Wheels of F1 racing car, airplane at take-off or landing,...)
Compare with stationary wheel (simple wind tunnel testing)

 *velocity xz*  *velocity xy*  *transversal velocities*

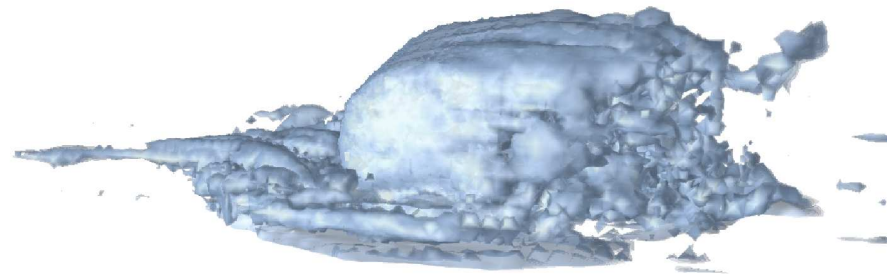
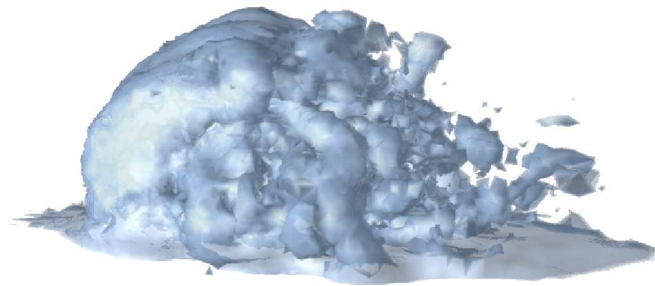
G2: rot vs stat: velocity



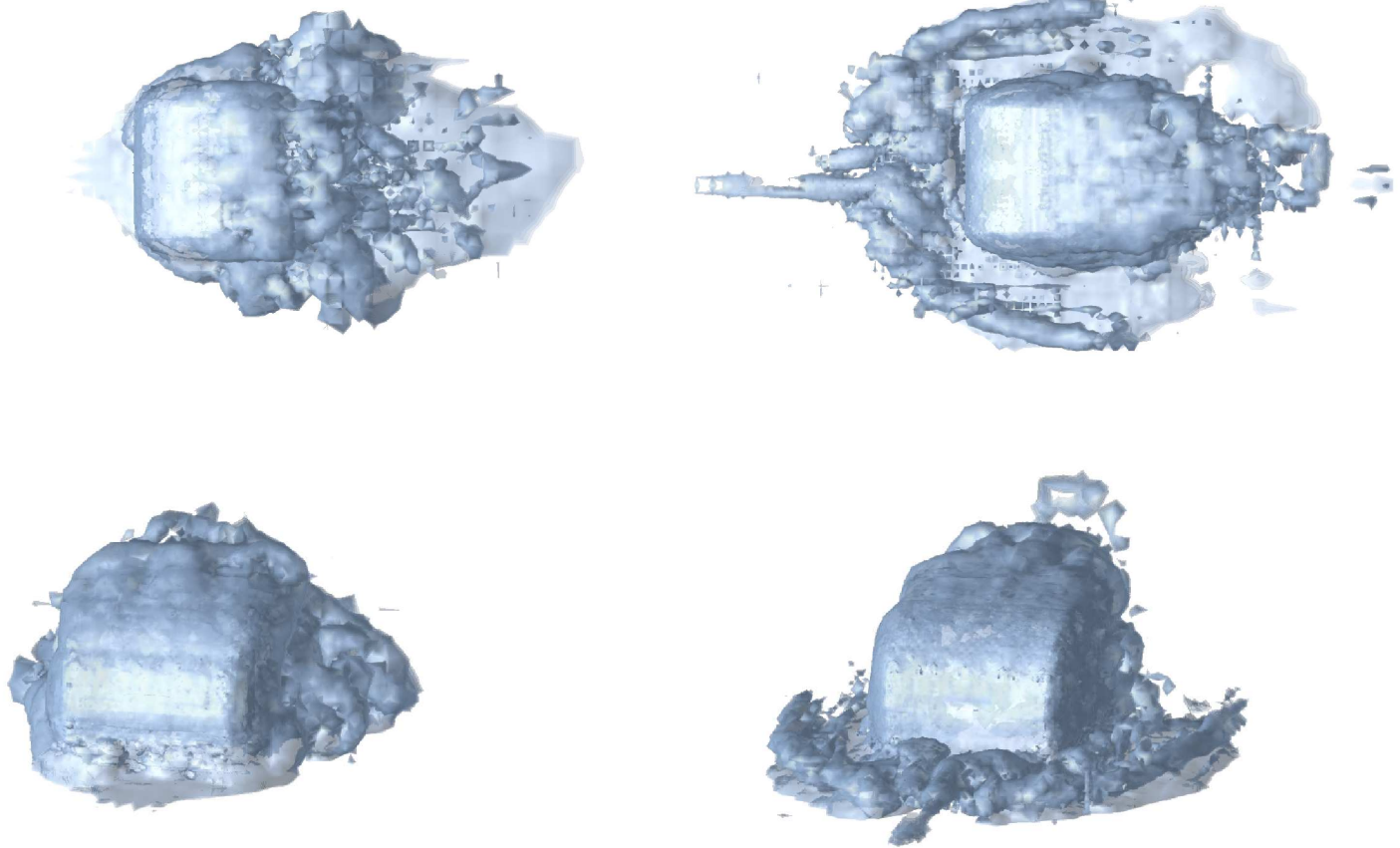
G2: rot vs stat: pressure



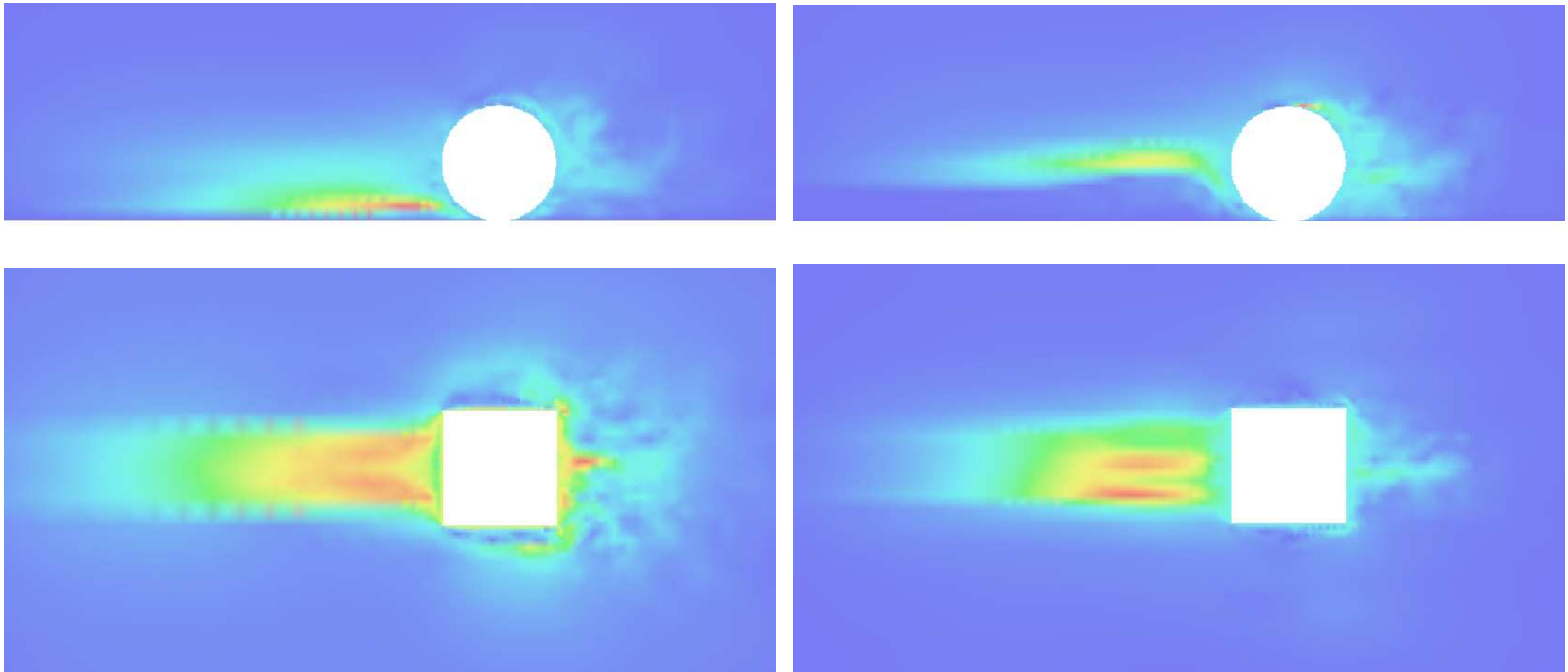
G2: rot vs stat: vorticity



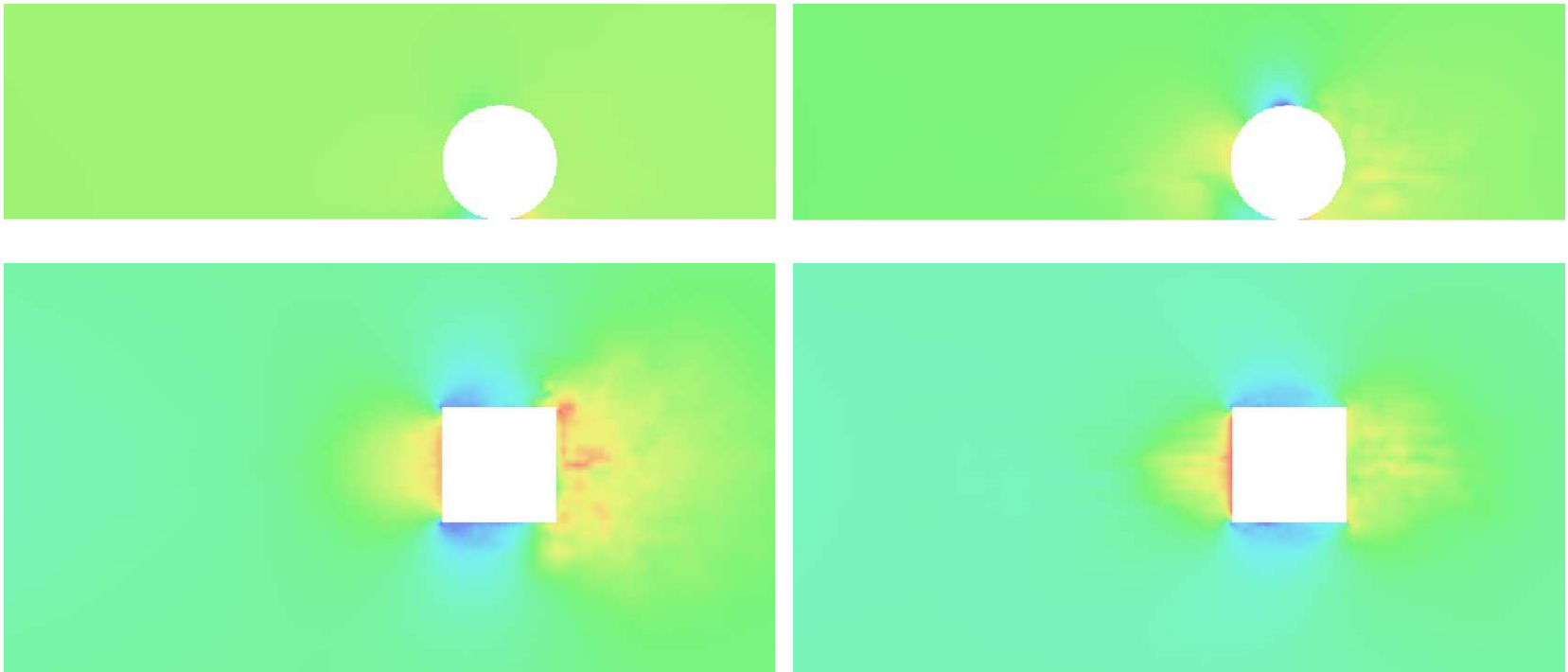
G2: rot vs stat: vorticity



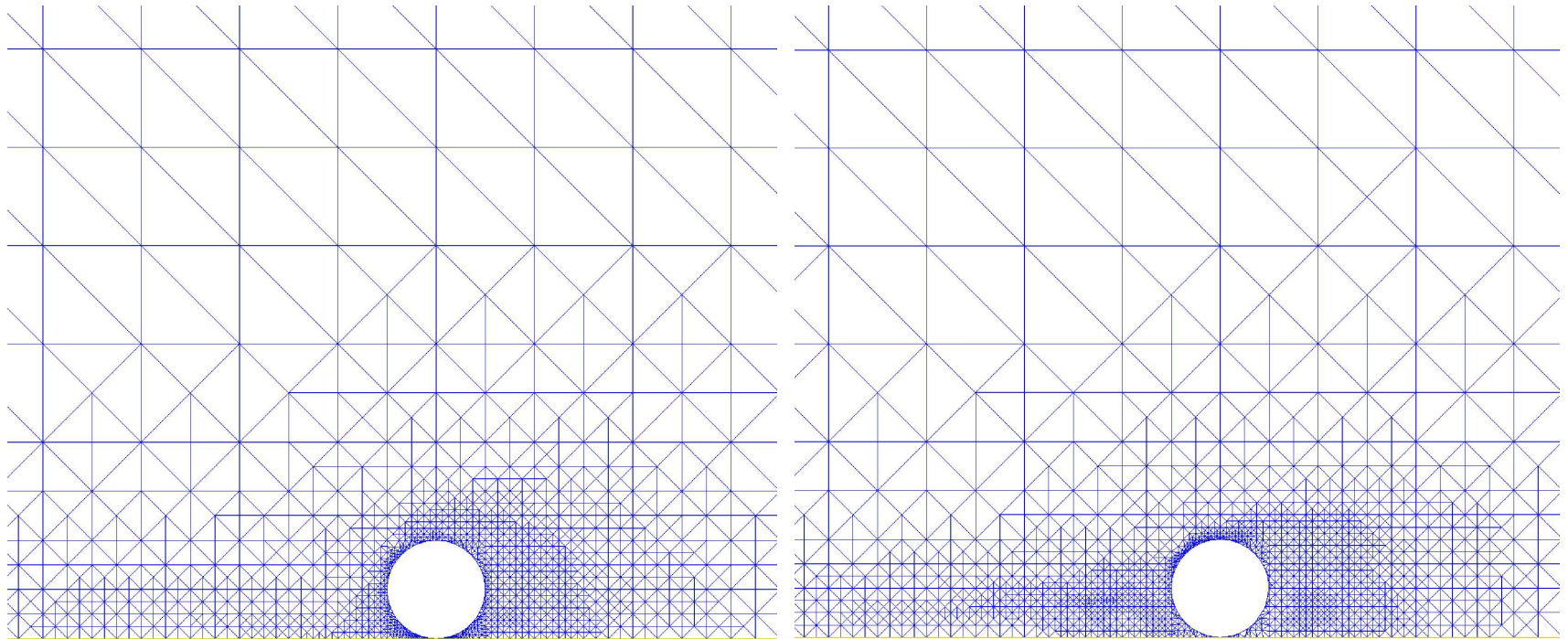
G2: rot vs stat: dual velocity



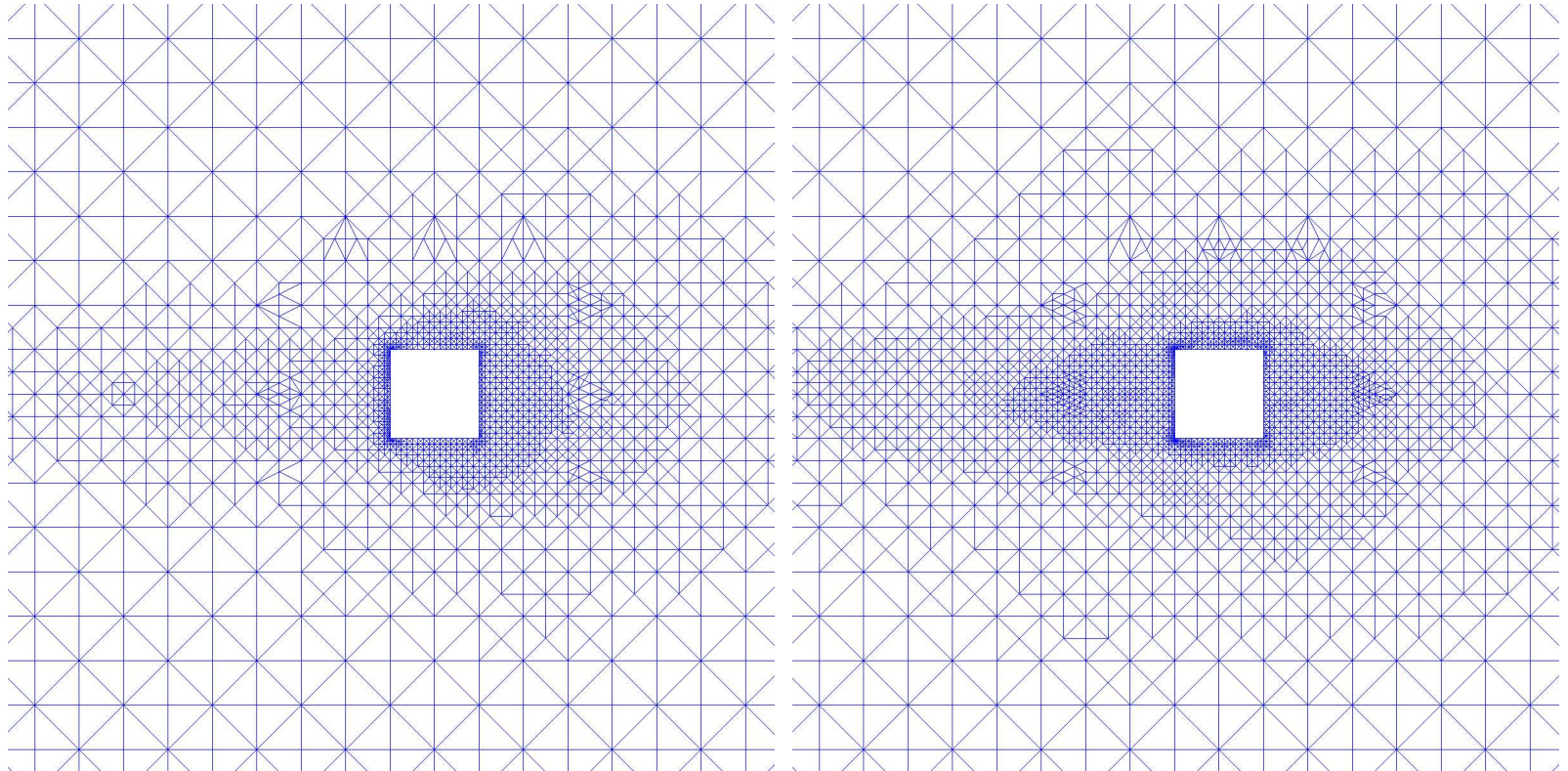
G2: rot vs stat: dual pressure



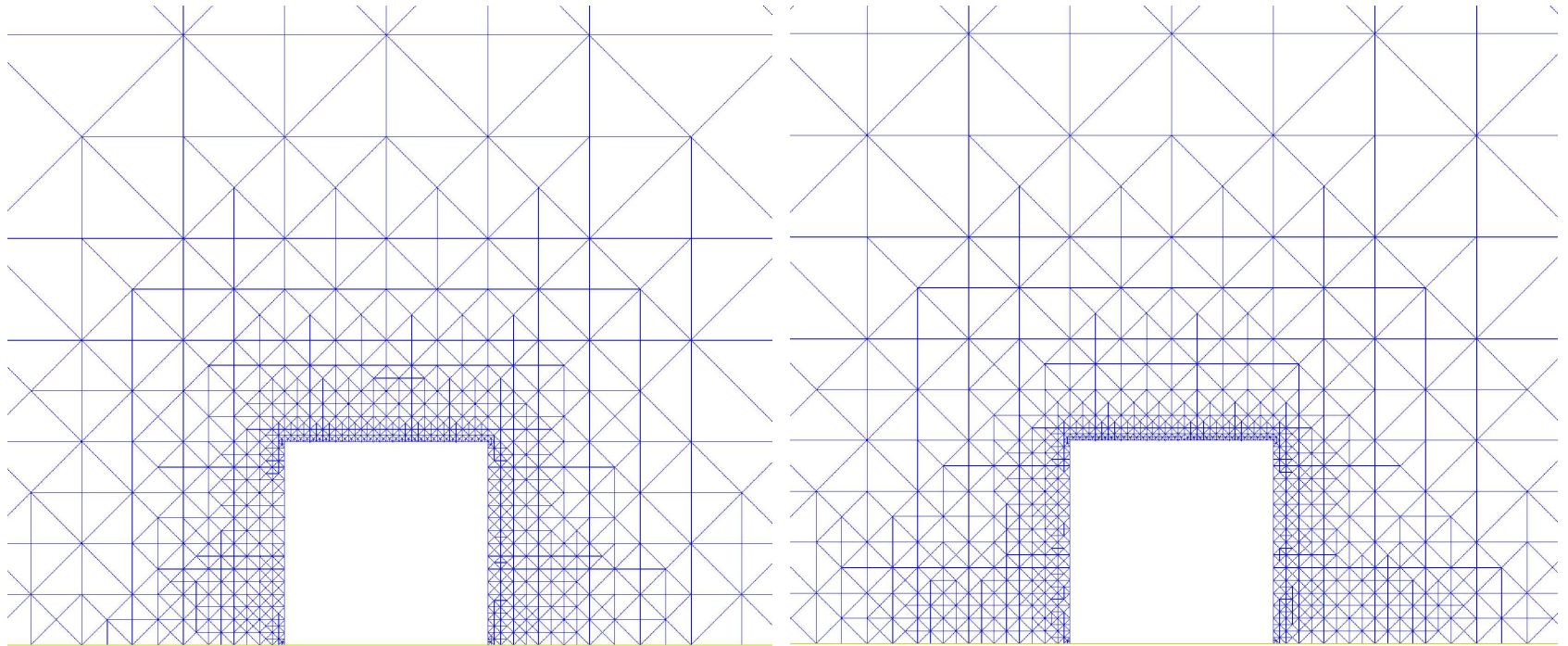
G2: rot vs stat: mesh



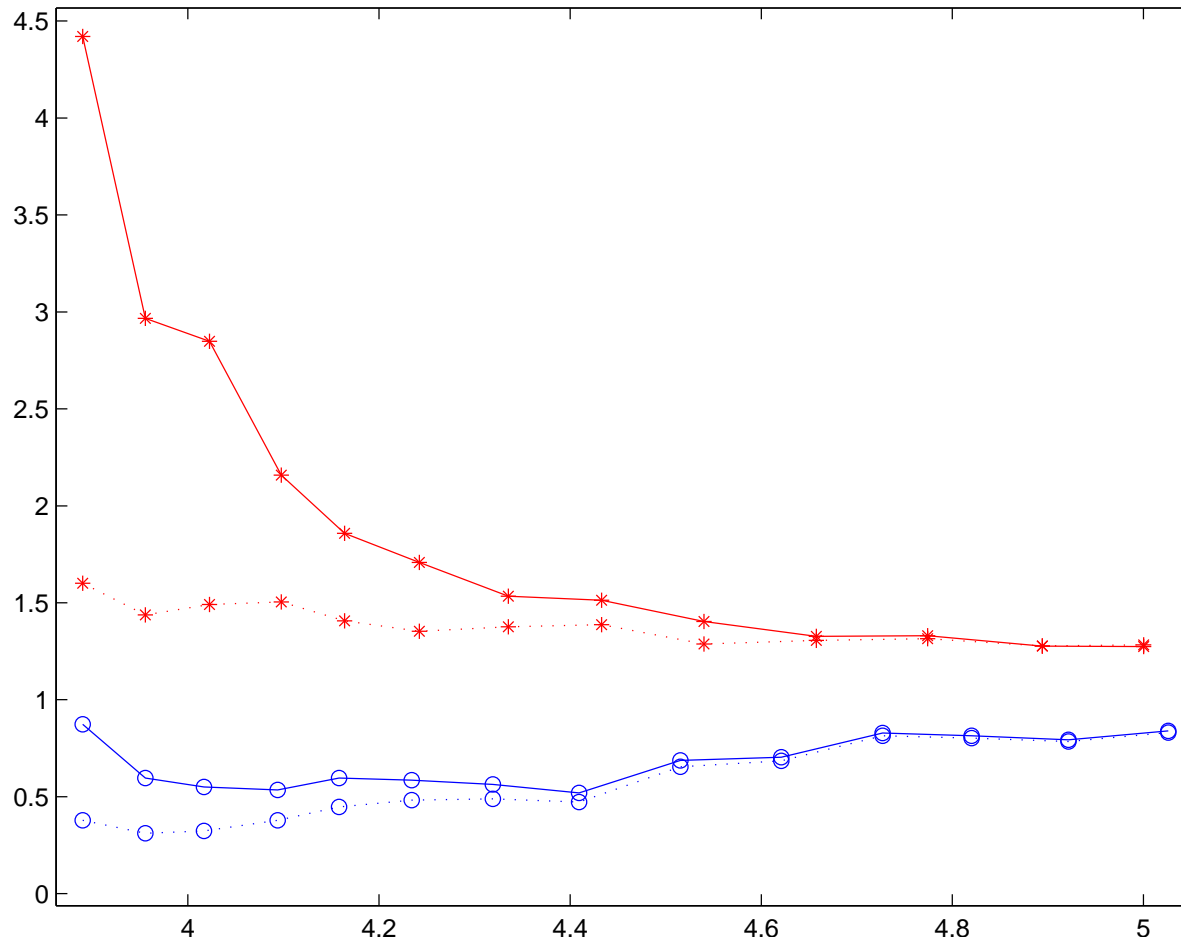
G2: rot vs stat: mesh



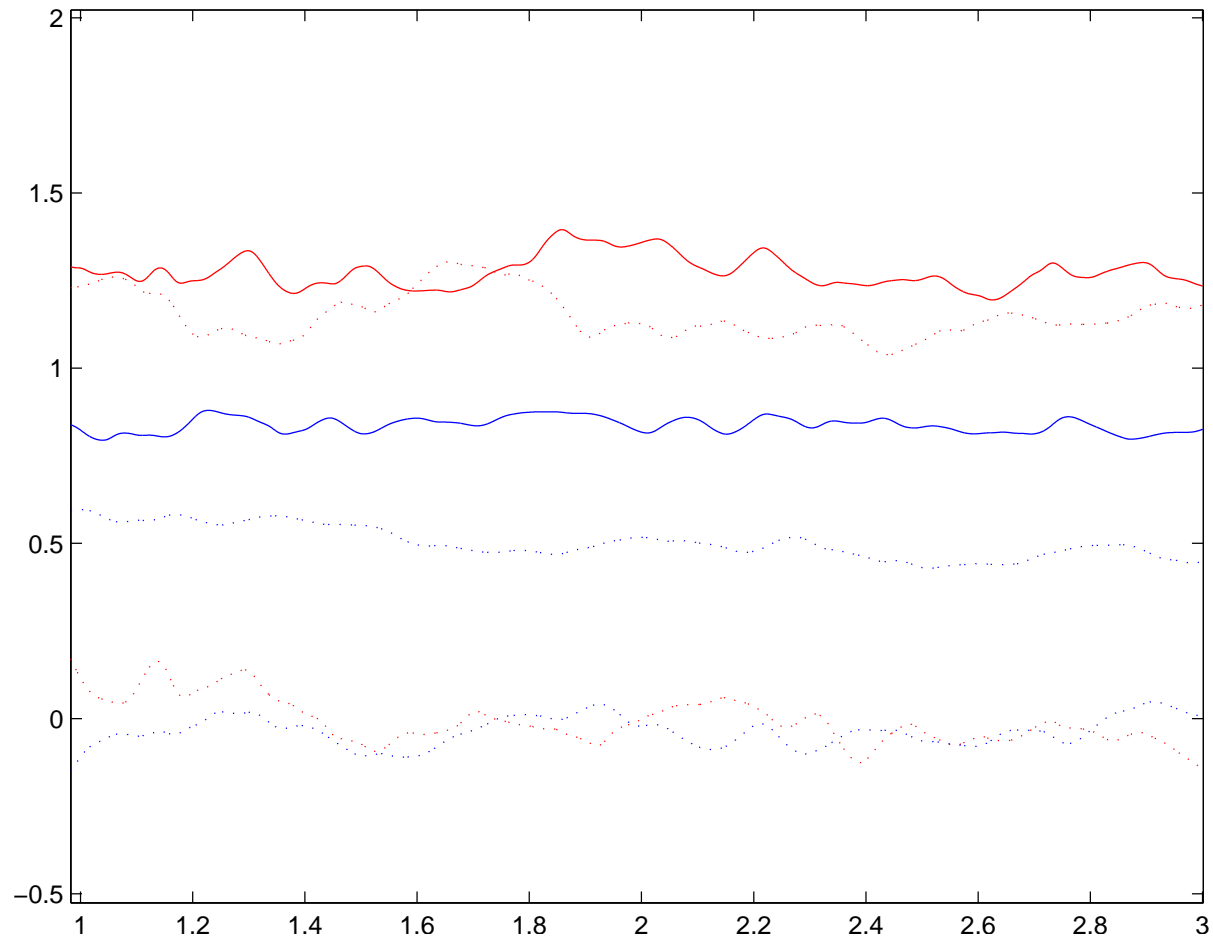
G2: rot vs stat: mesh



G2: rot vs stat: drag c_D : 1.3 vs 0.8



G2: rot vs stat: forces



Turbulent boundary layers

So far: No slip boundary conditions seems ok for modeling laminar boundary layers (separation and skin friction)

For high Re boundary layer undergoes transition

Extremely expensive to resolve turbulent boundary layer:

⇒ We need wall-modeling

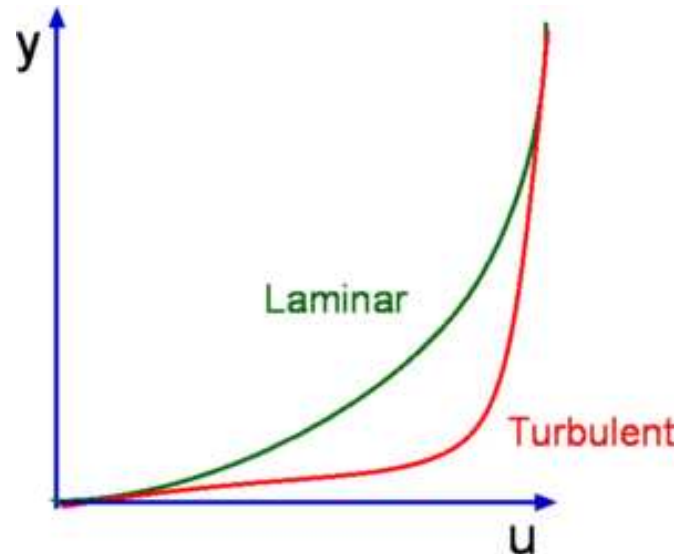


Turbulent boundary layers

Separation given by: $\dot{u} + u \cdot \nabla u - \nu \Delta u = -\nabla p$

Retarding flow near the boundary: $\nabla p > 0$

Turbulent boundary layer \Rightarrow higher momentum near the boundary \Rightarrow delayed separation



Friction boundary condition

Slip with friction boundary condition [Maxwell, Navier,....]

Friction coefficient β ; $\beta = 0$: slip b.c., $\beta = \infty$: no slip b.c.

[LES + Boundary Layer theory: Layton, John, Illiescu,....]

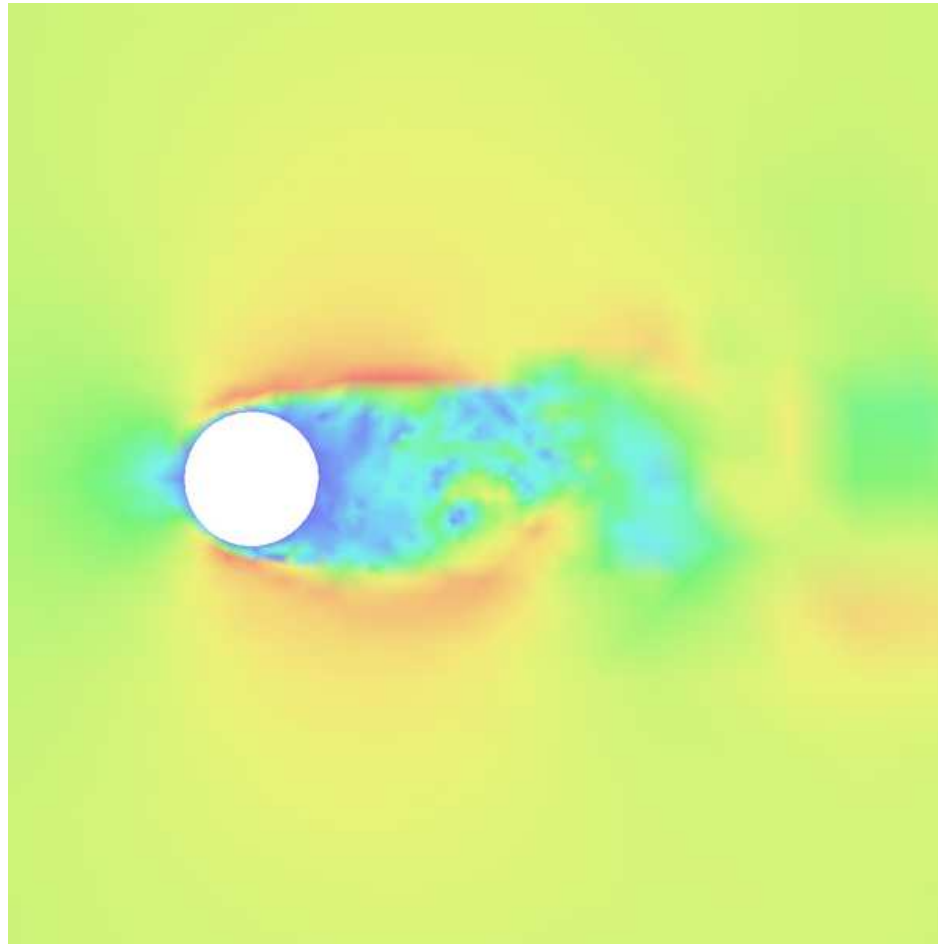
Simple wall model: $\beta \approx c_f$ (skin friction $c_f \sim Re^{-0.2}$)

$$\beta = \beta(Re, h); \lim_{h \rightarrow 0} \beta = \infty, \lim_{Re \rightarrow \infty} \beta = 0 \quad (\lim_{Re \rightarrow 0} \beta = \infty)$$

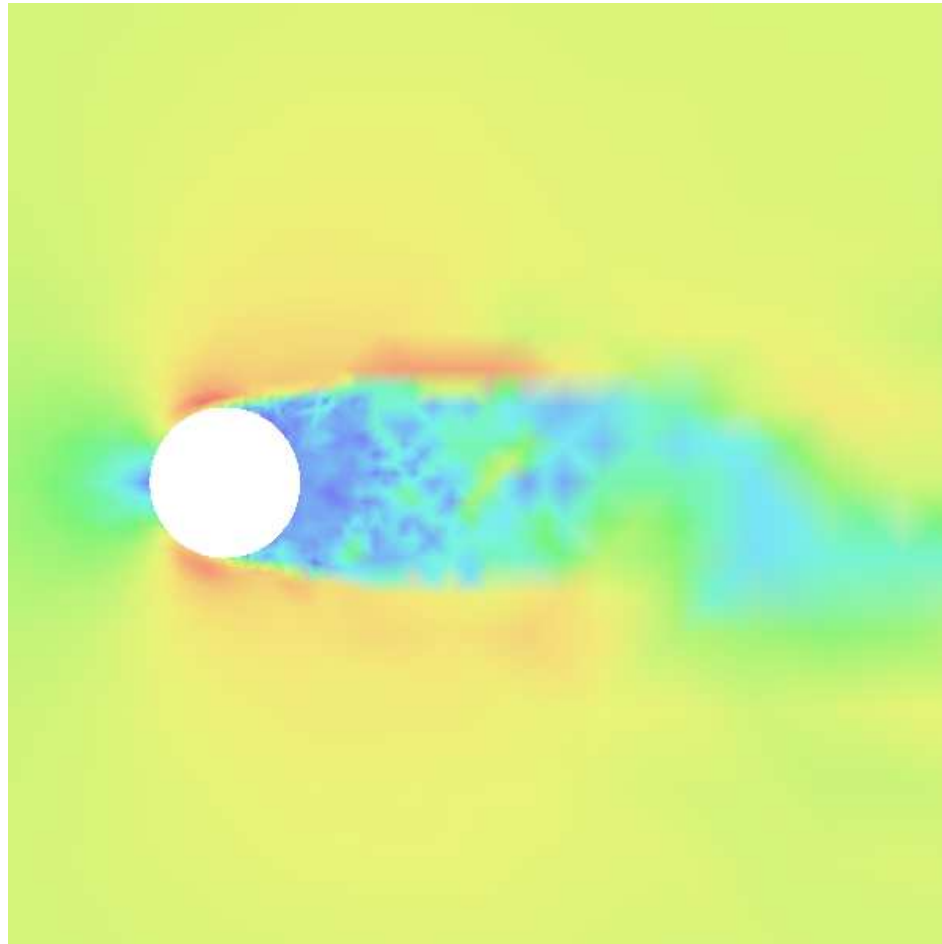
$$\frac{1}{2} \|U(t)\|^2 + \sum_{i=1}^2 \|\sqrt{\beta} u \cdot \tau_i\|_{\Gamma \times I}^2 + \|\sqrt{h} R(\hat{U})\|_Q^2 = \frac{1}{2} \|U(0)\|^2$$

(with ν small)

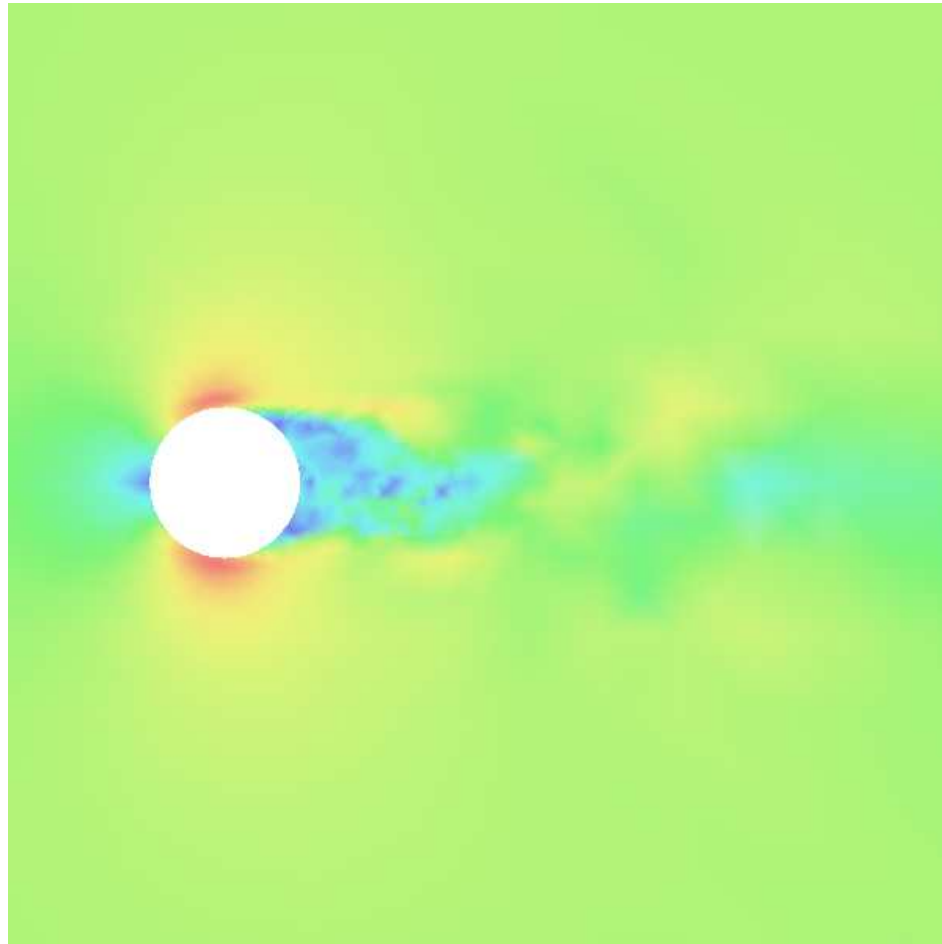
drag crisis; $\beta = 1$: $c_D \approx 1.0$



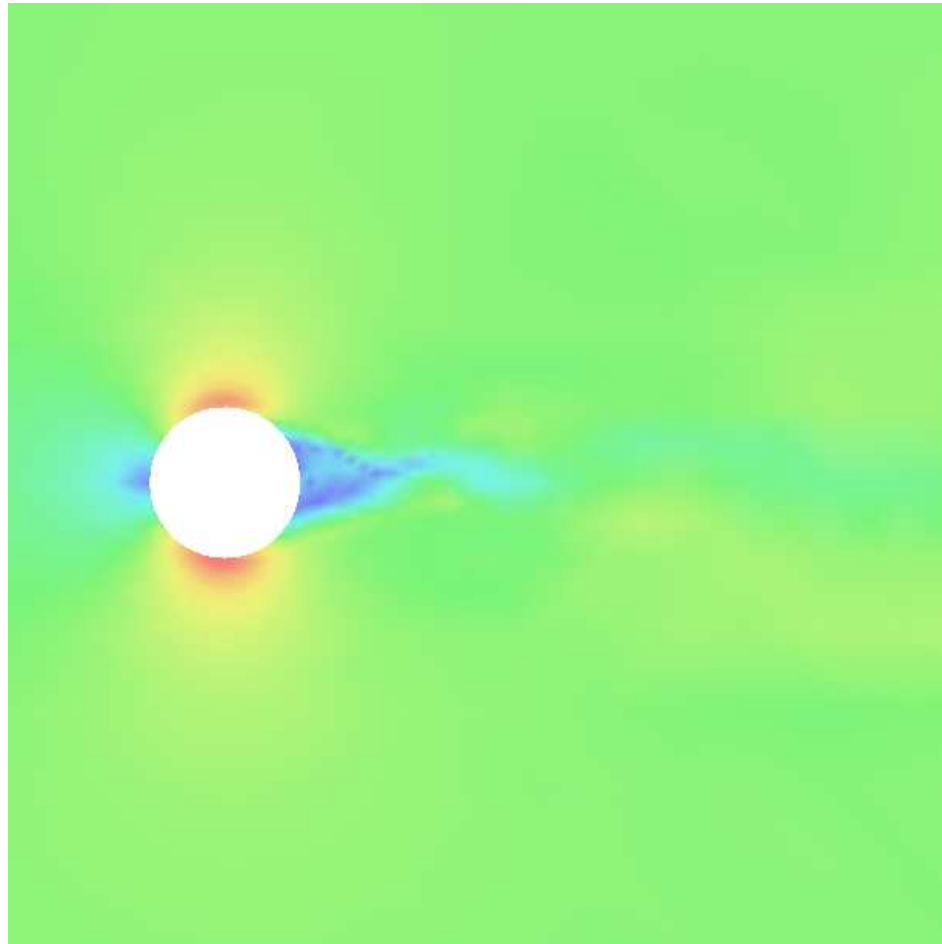
drag crisis; $\beta = 2 \times 10^{-2}$: $c_D \approx 0.7$



drag crisis: $\beta = 1 \times 10^{-2}$; $c_D \approx 0.5$



drag crisis; $\beta = 5 \times 10^{-3}$: $c_D \approx 0.45$

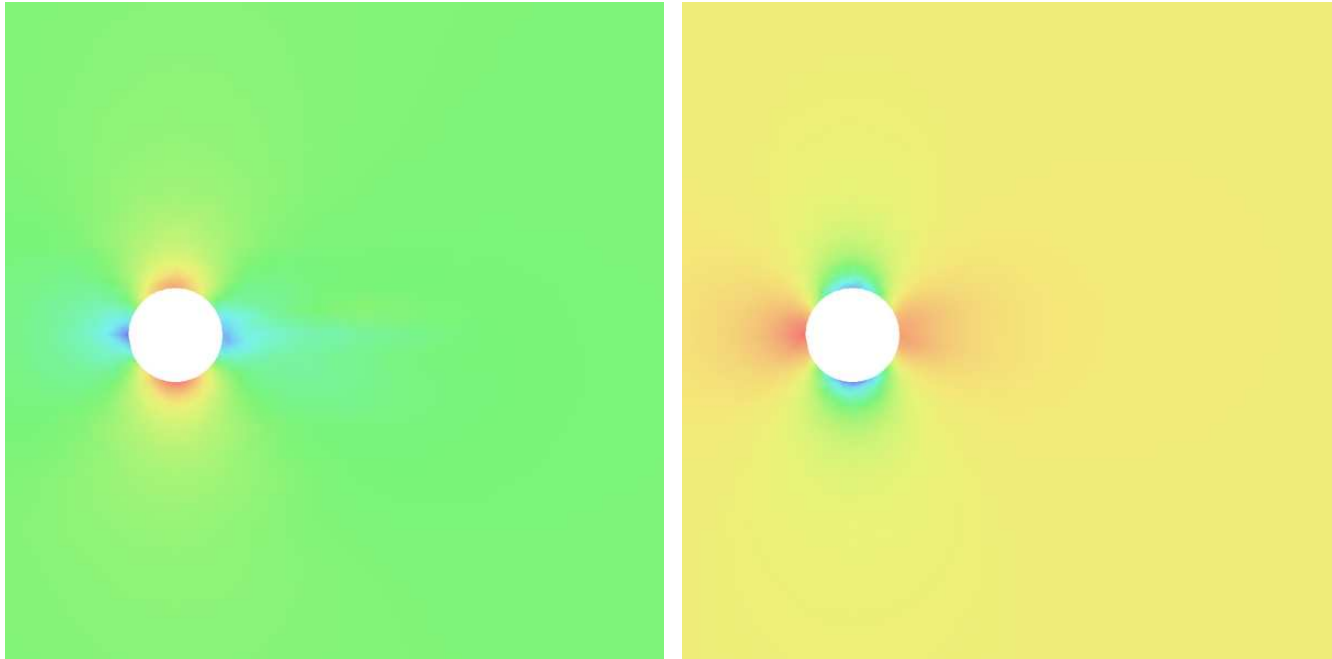


EG2 and Turbulent Euler solutions

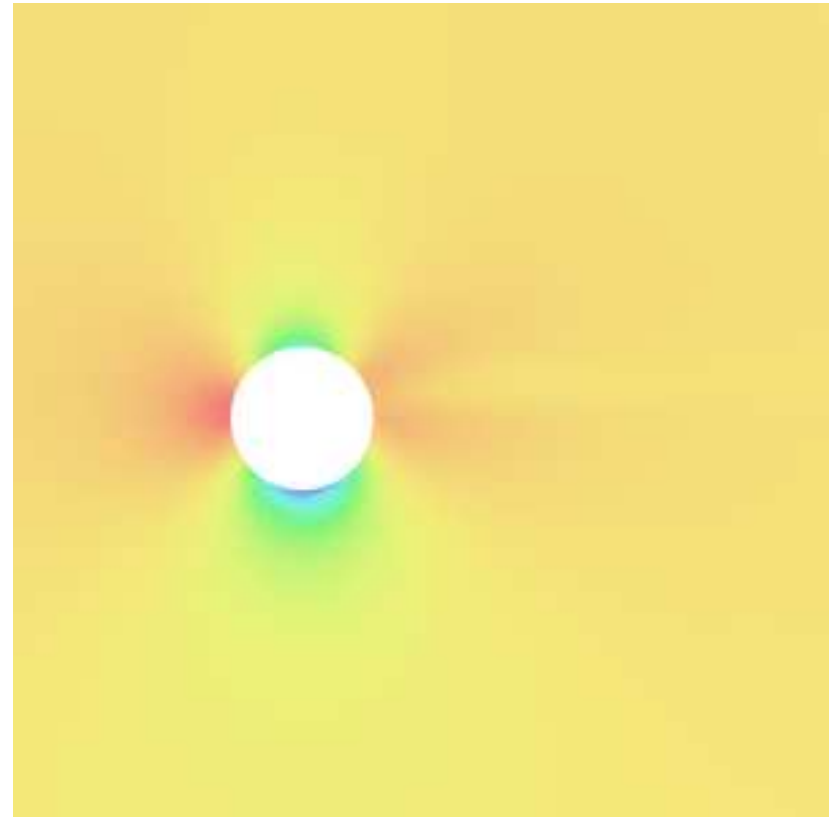
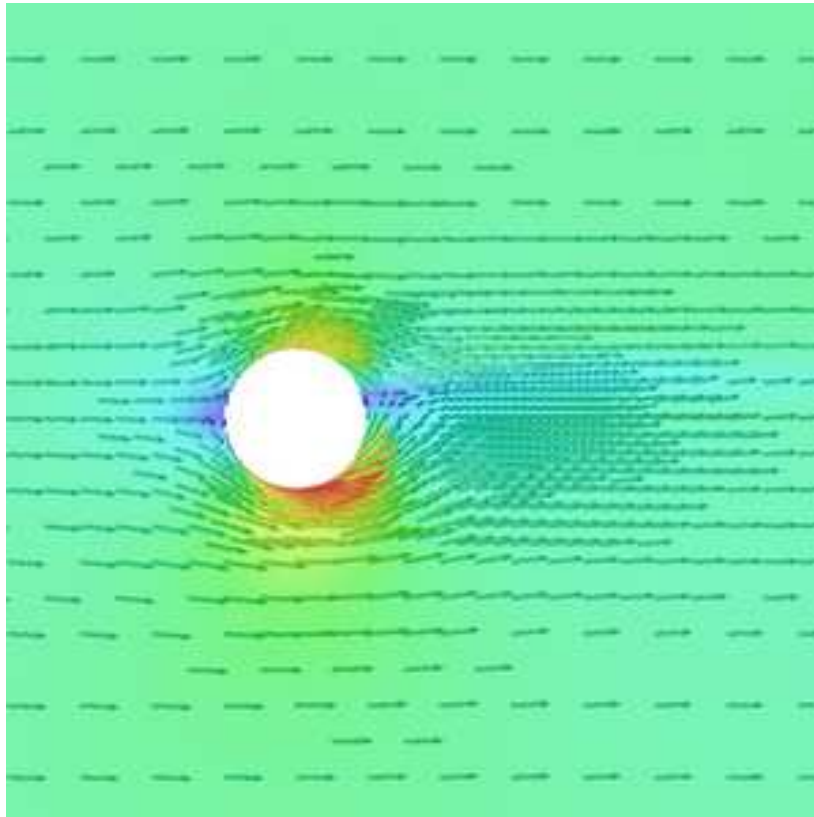
$\beta \rightarrow 0; Re \rightarrow \infty (\nu \rightarrow 0) \Rightarrow$ Euler/G2 + slip b.c. (EG2)

EG2: no empirical parameters; only h (very general...)

What happens in the limit? The potential solution ($c_D = 0$)?

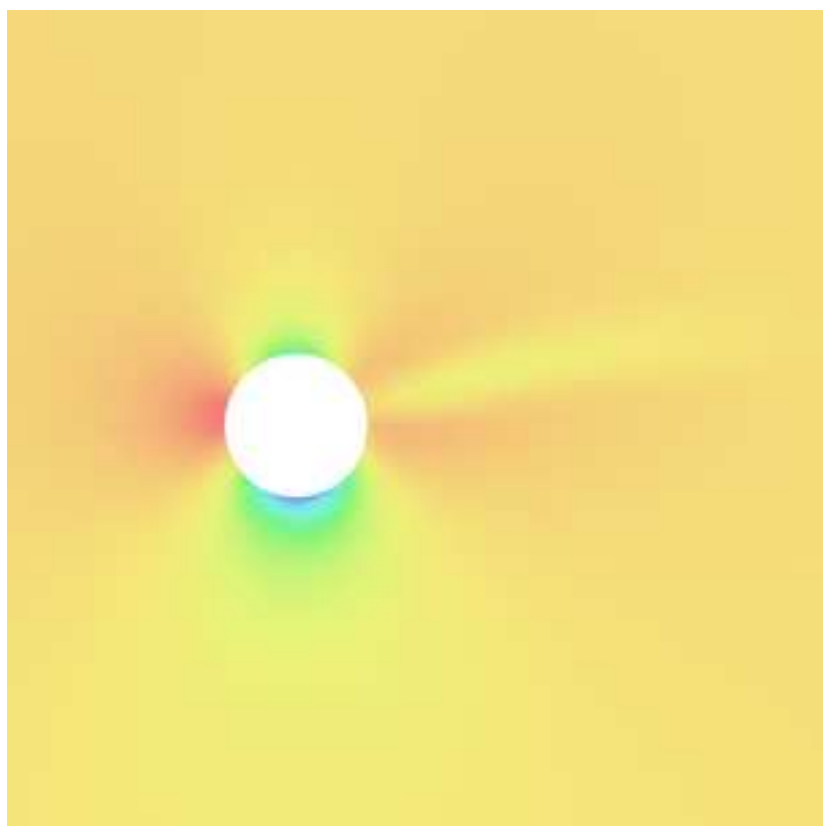
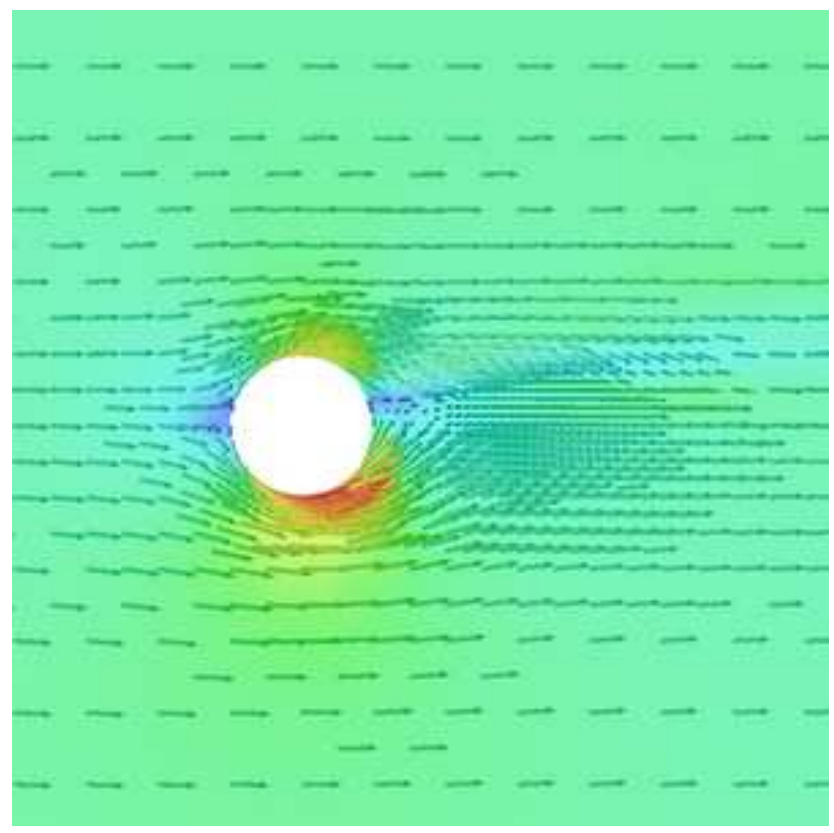


Velocity & pressure: $t=1.0$: $c_D = 0.22$



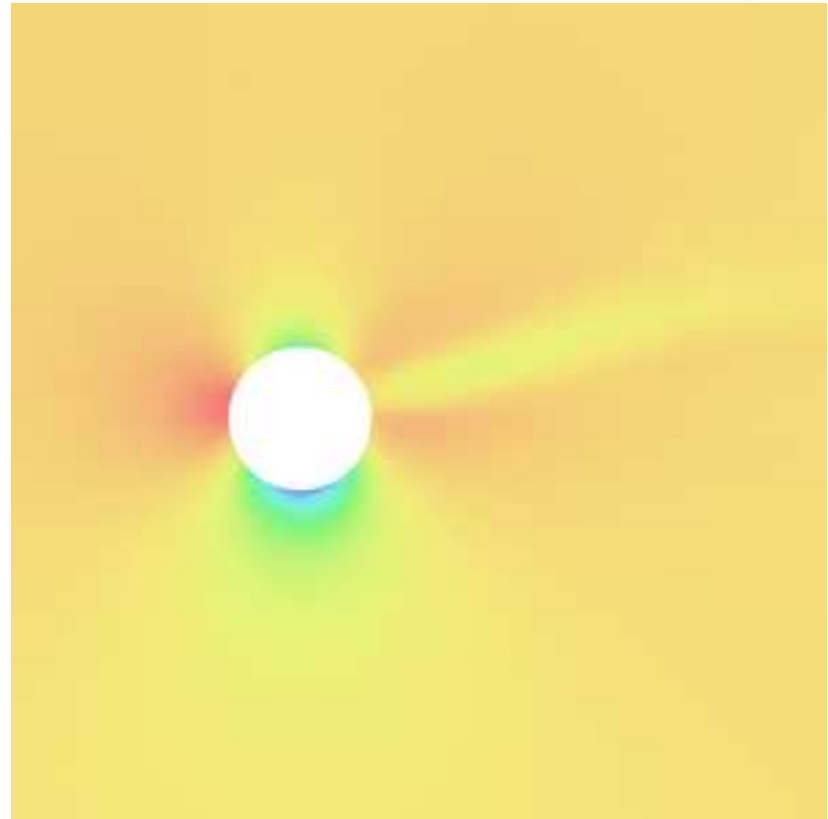
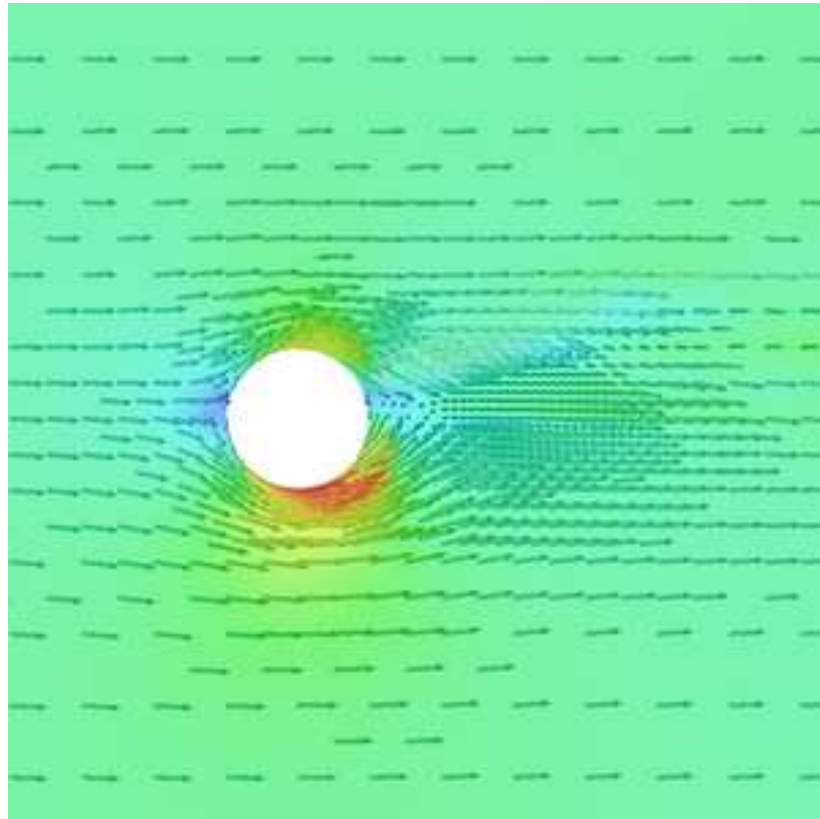
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Velocity & pressure: $t=1.25: c_D = 0.25$



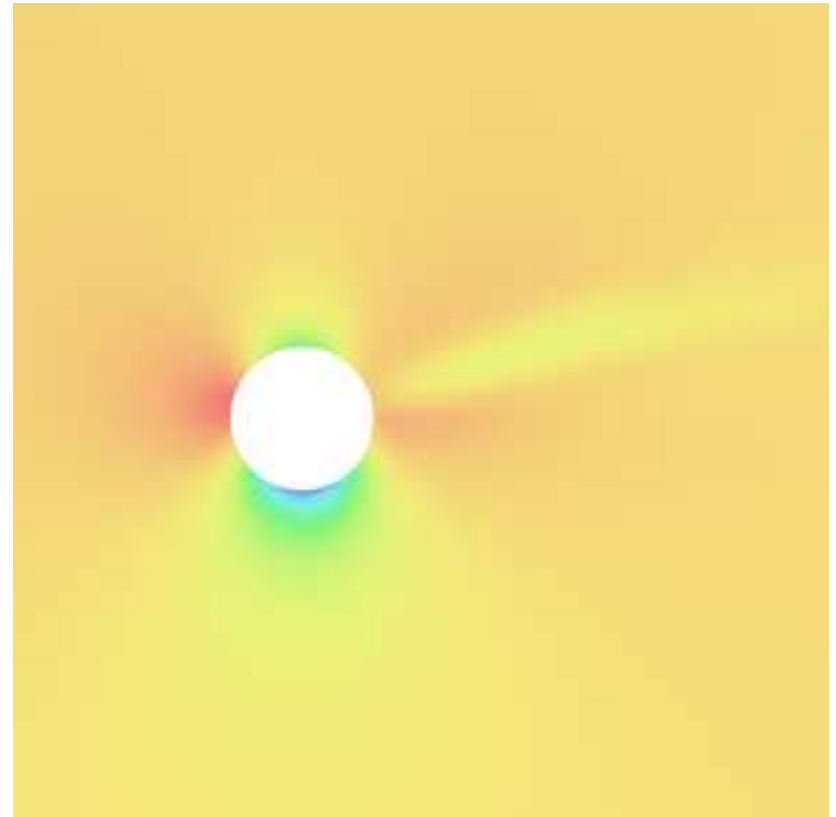
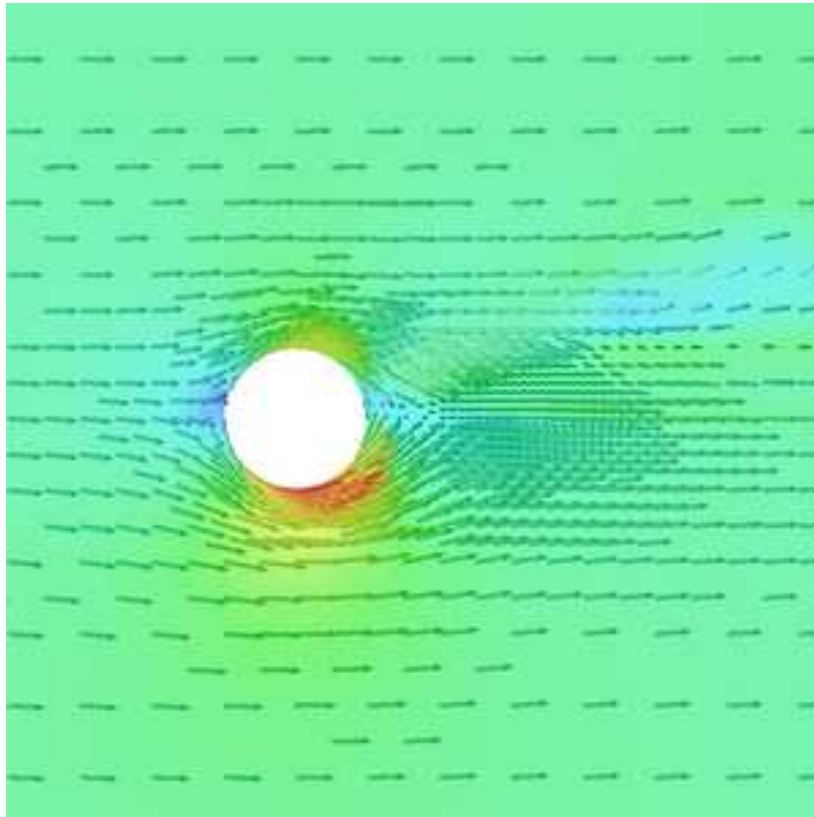
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Velocity & pressure: $t=1.5$: $c_D = 0.28$



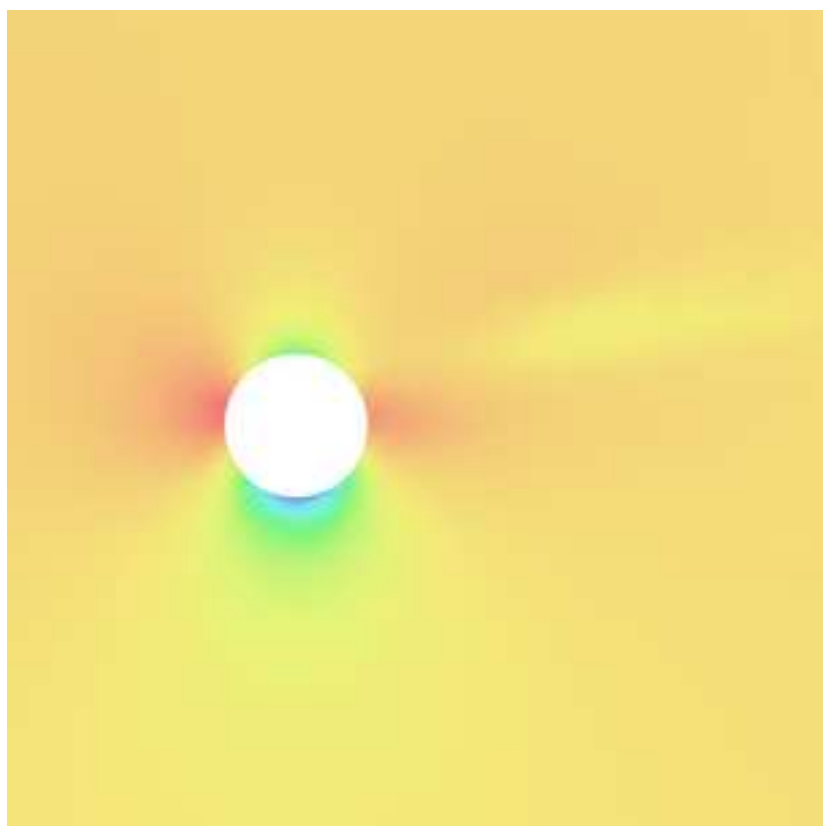
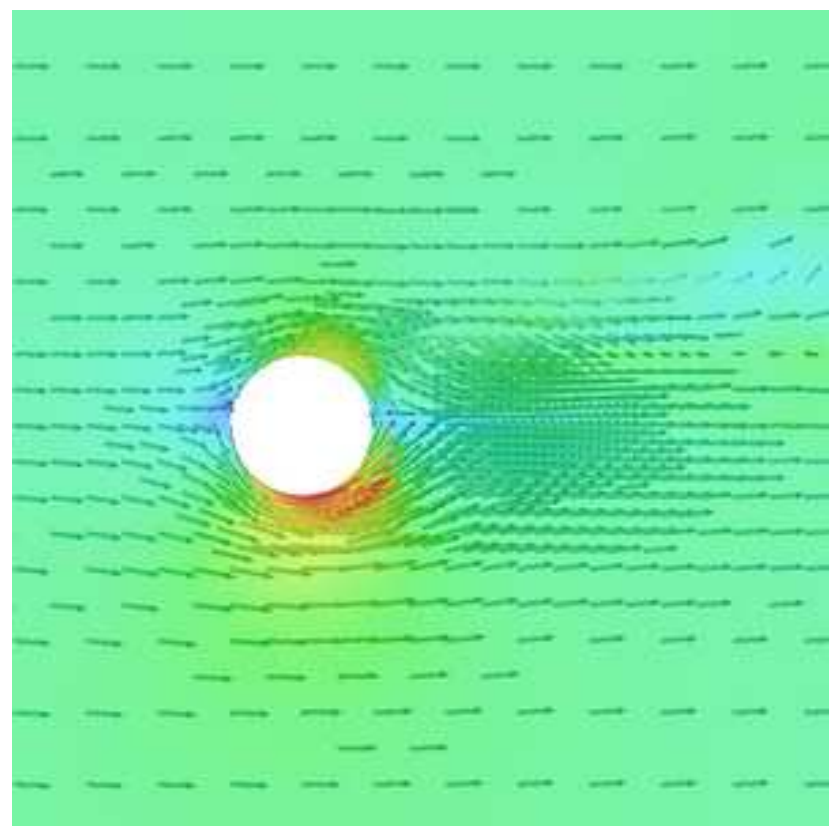
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Velocity & pressure: $t=1.75$: $c_D = 0.36$



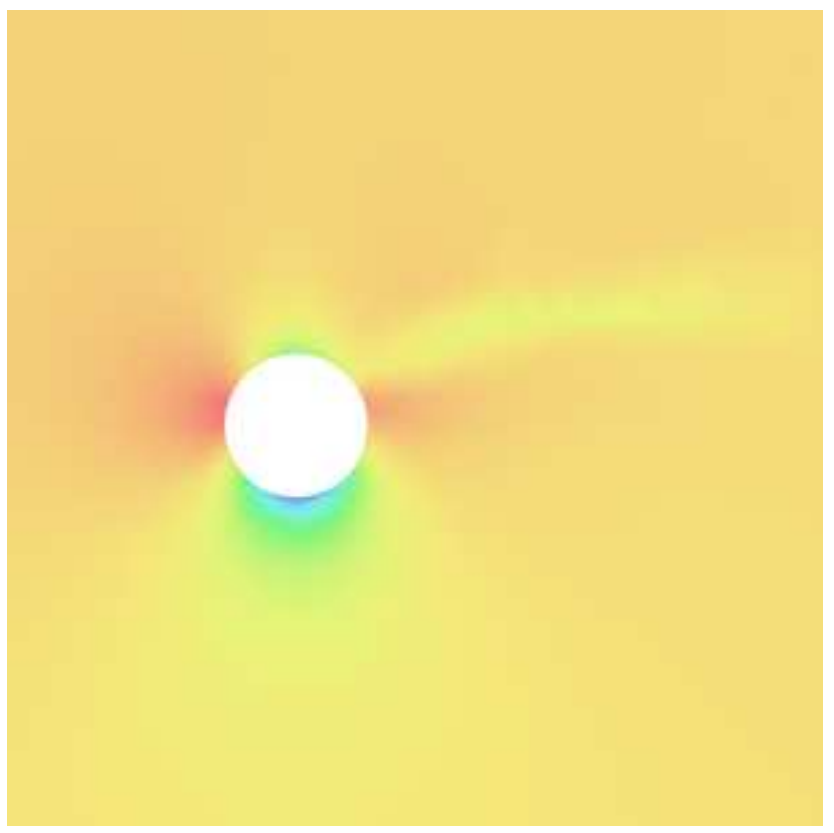
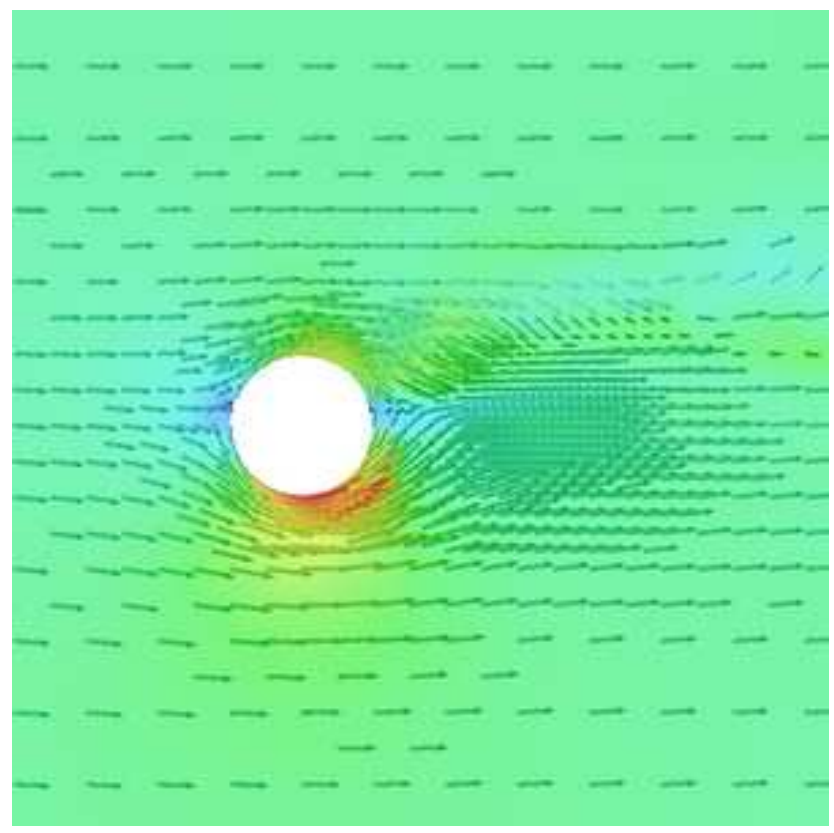
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Velocity & pressure: $t=2.0$: $c_D = 0.51$



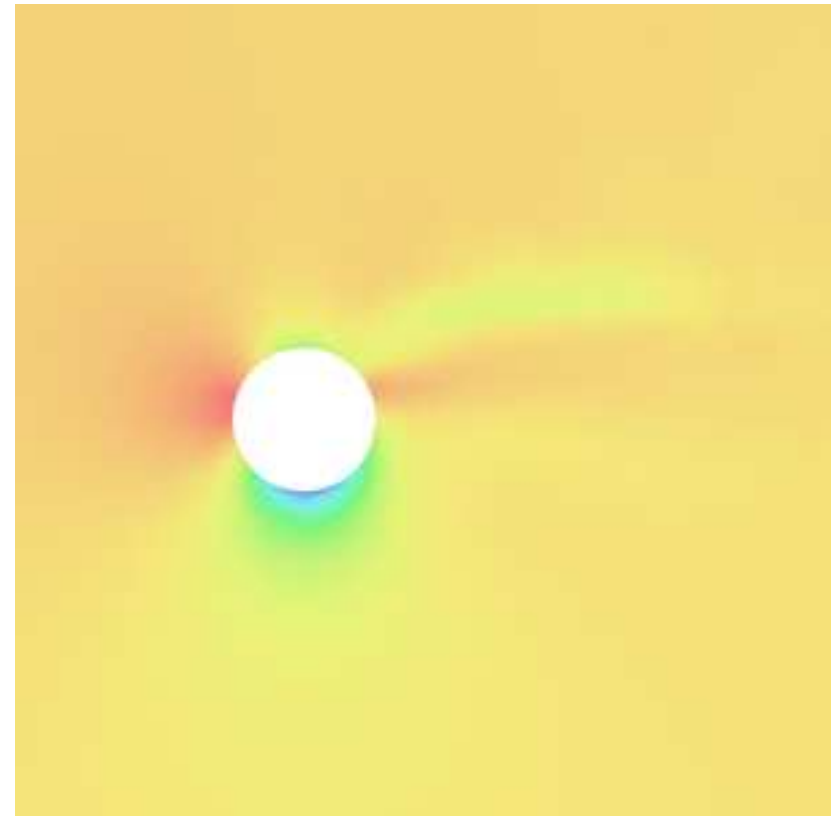
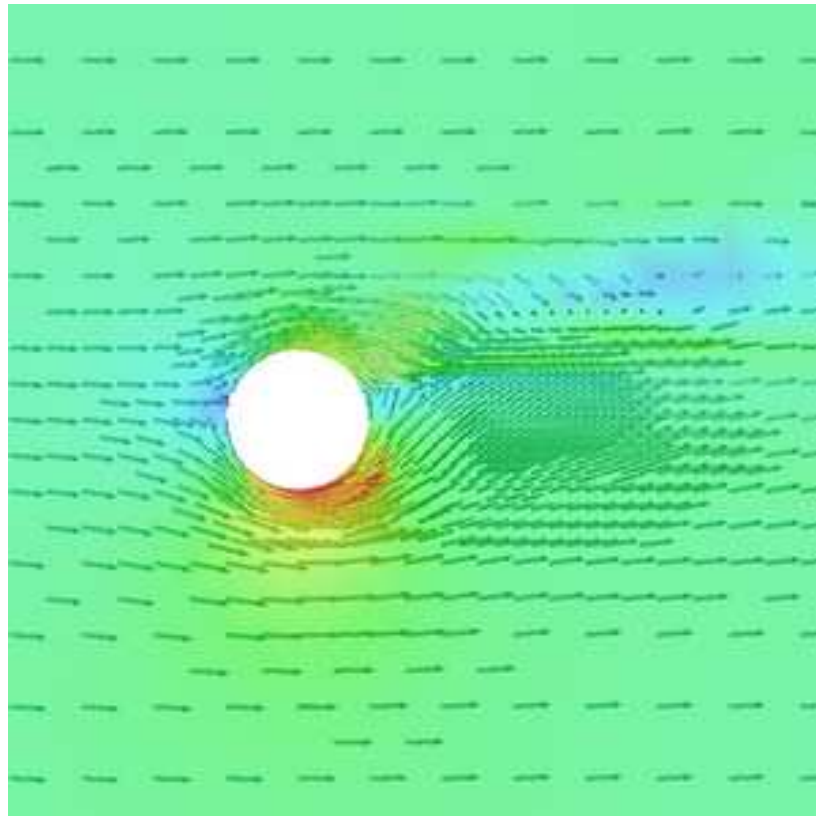
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Velocity & pressure: $t=2.25$: $c_D = 0.78$

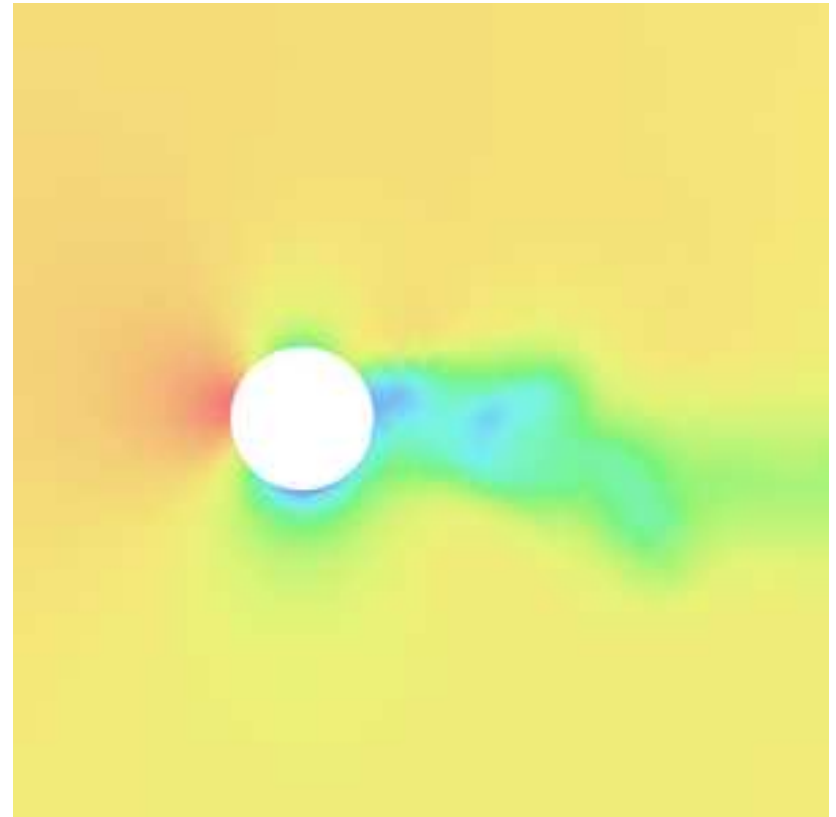
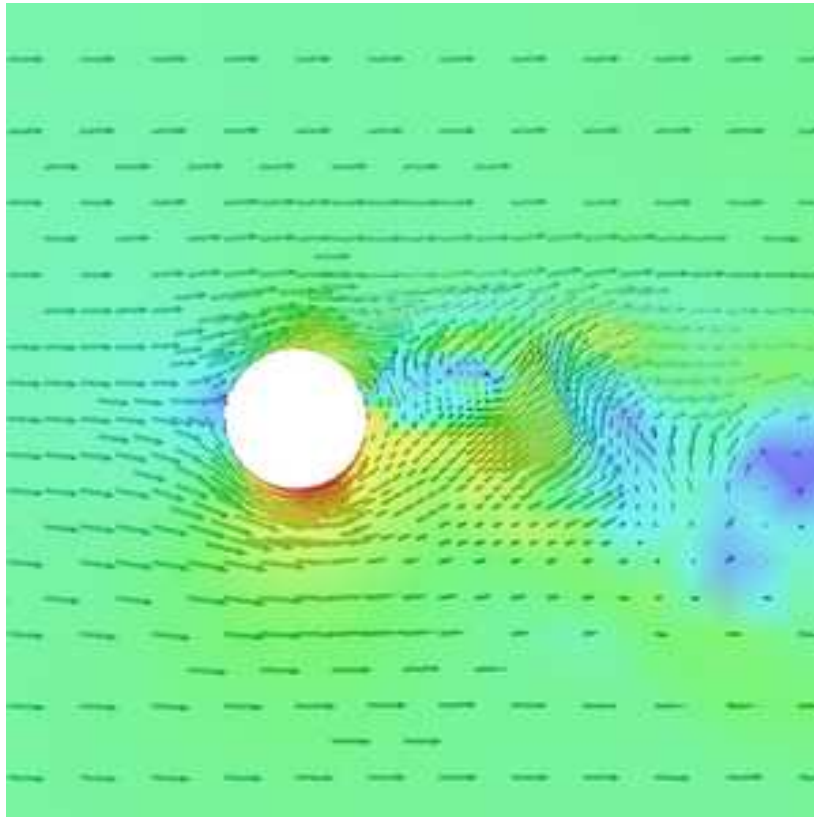


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Velocity & pressure: $t=2.5$: $c_D = 1.14$

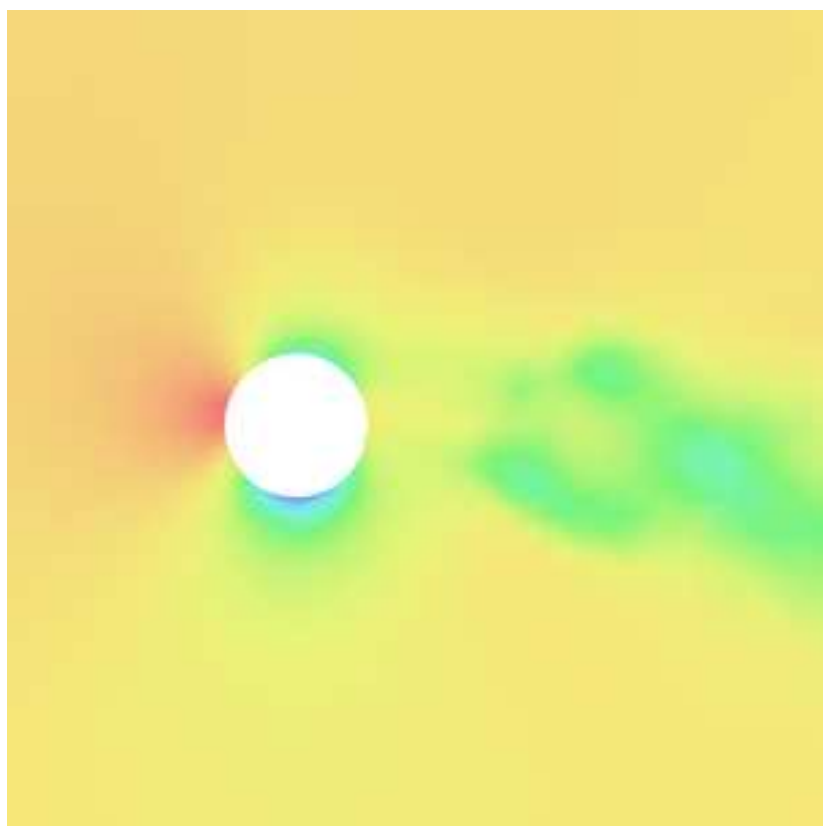
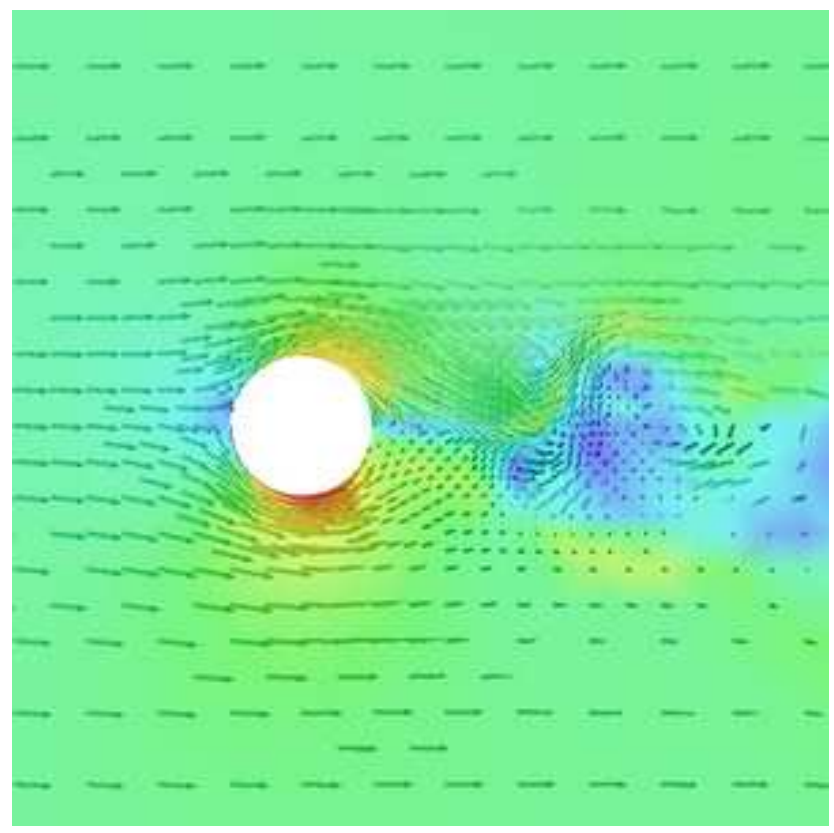


Velocity & pressure: $t=4.5$: $c_D = 1.63$



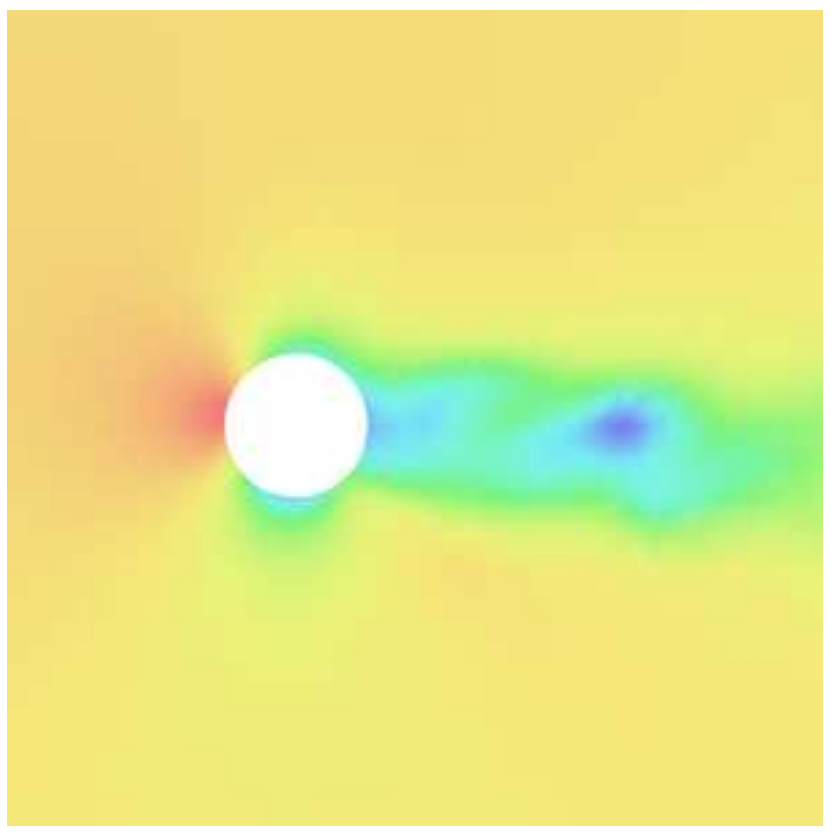
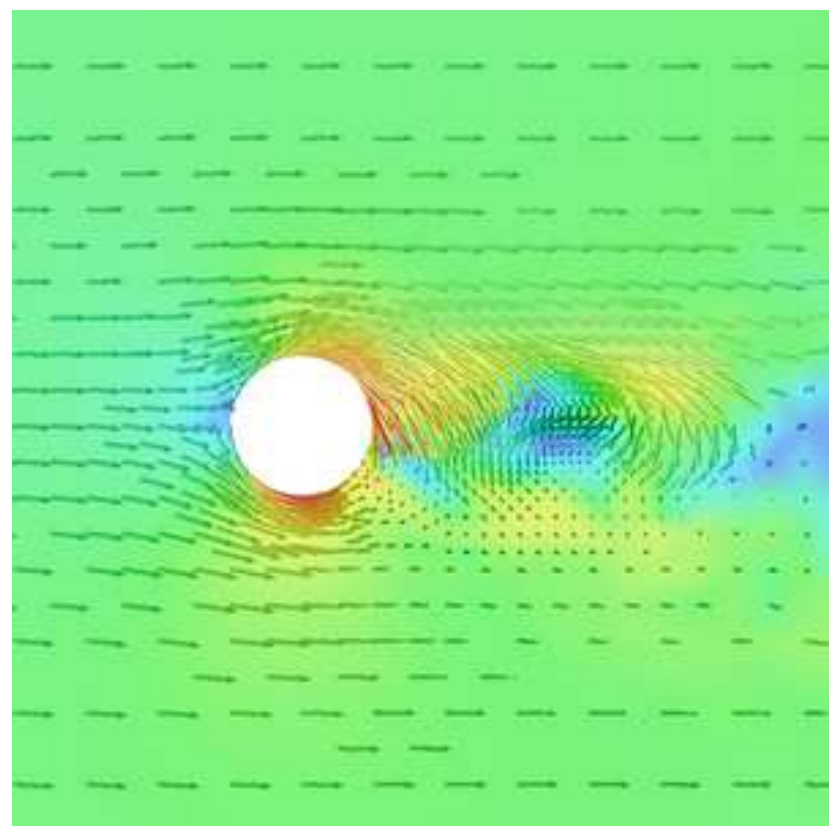
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Velocity & pressure: $t=5.0$: $c_D = 1.79$



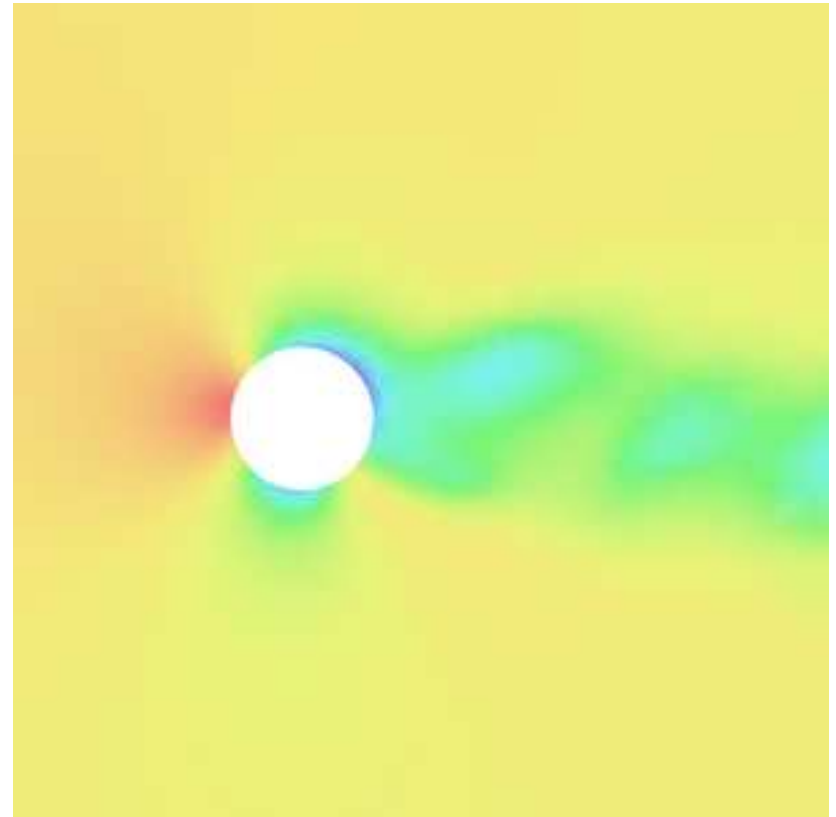
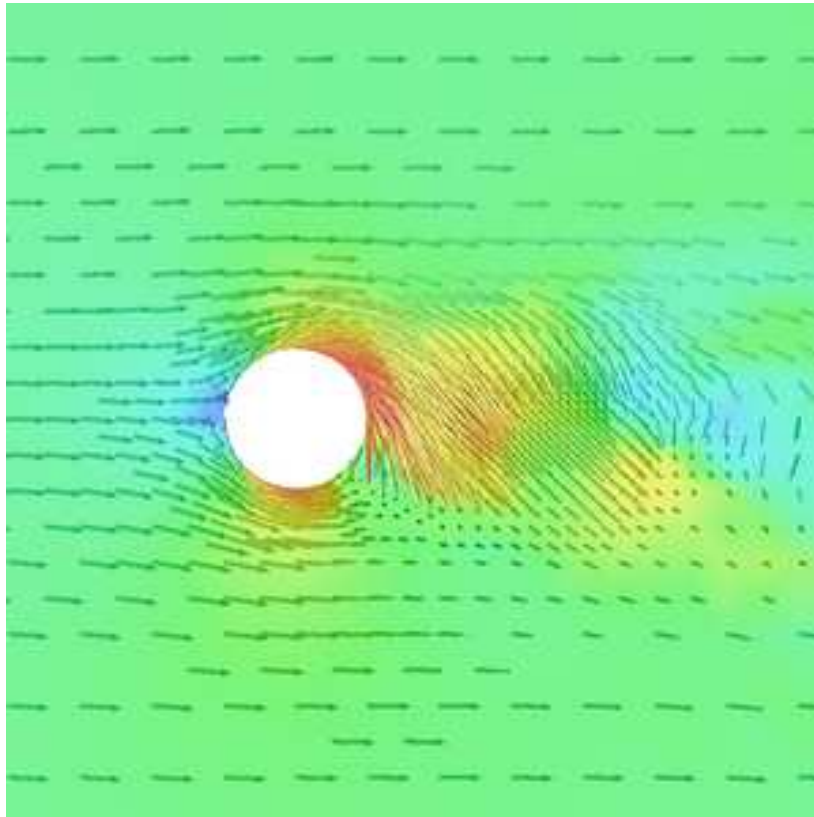
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Velocity & pressure: $t=5.5$: $c_D = 1.96$



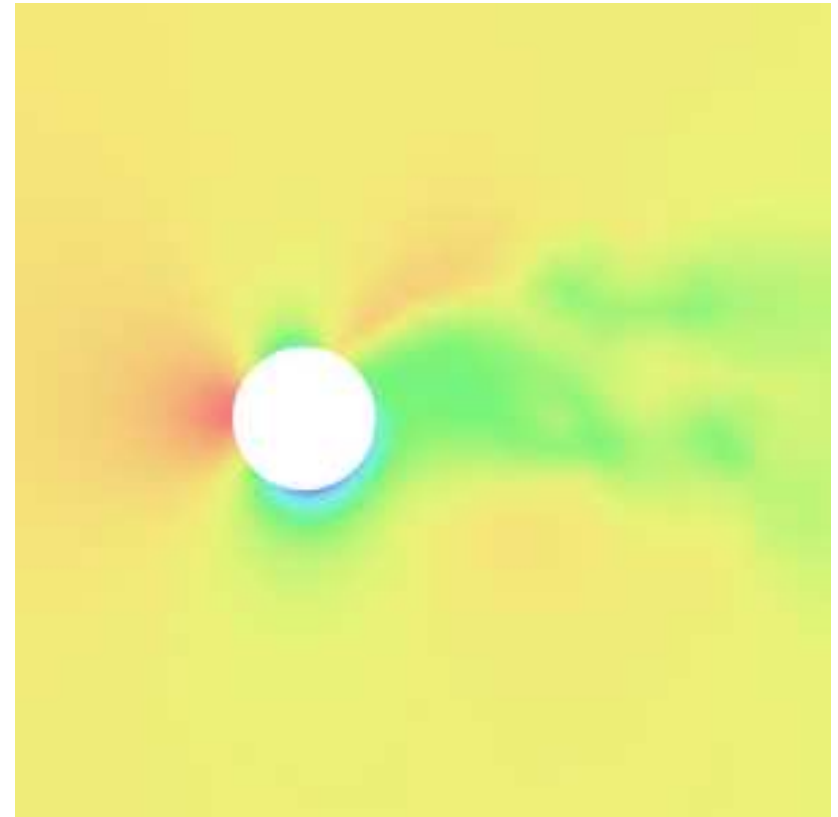
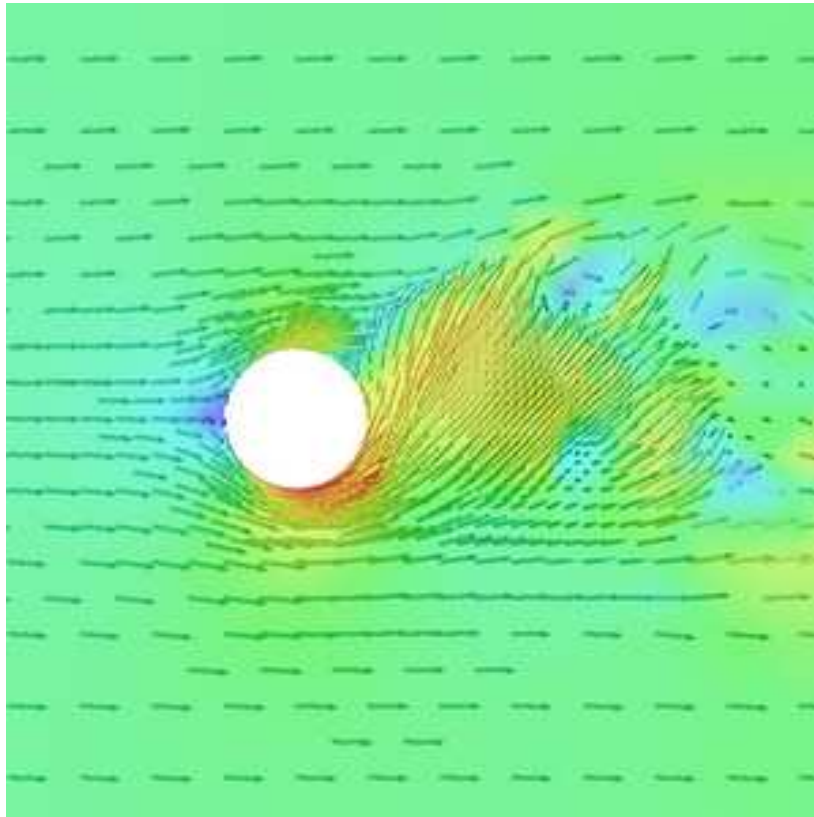
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Velocity & pressure: $t=5.75$: $c_D = 1.90$



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Velocity & pressure: $t=11.0$: $c_D = 1.82$



Vorticity: $\omega_1, \omega_2, \omega_3$: $t=1.0$: $c_D = 0.22$



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Vorticity: $\omega_1, \omega_2, \omega_3$: $t=1.25$: $c_D = 0.25$

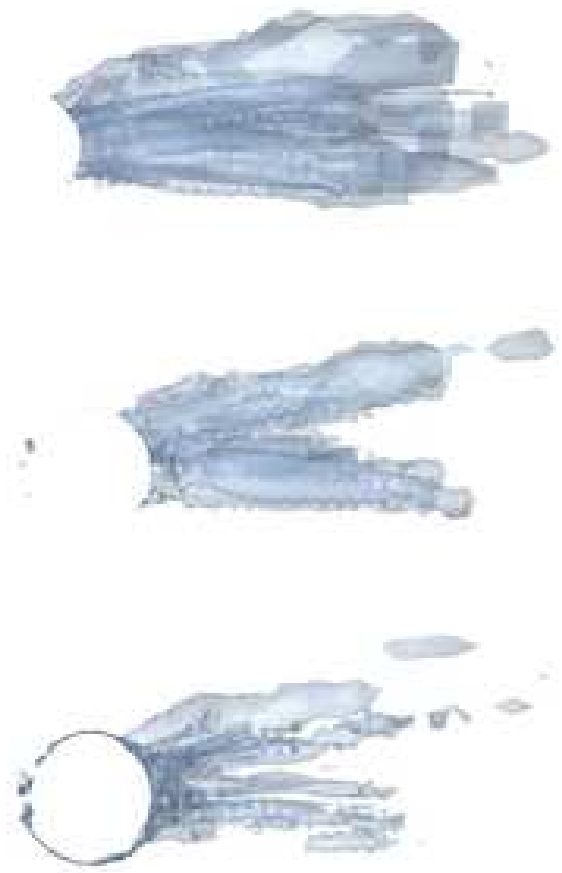


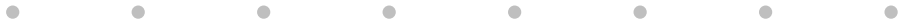
Vorticity: $\omega_1, \omega_2, \omega_3$: $t=1.5$: $c_D = 0.28$



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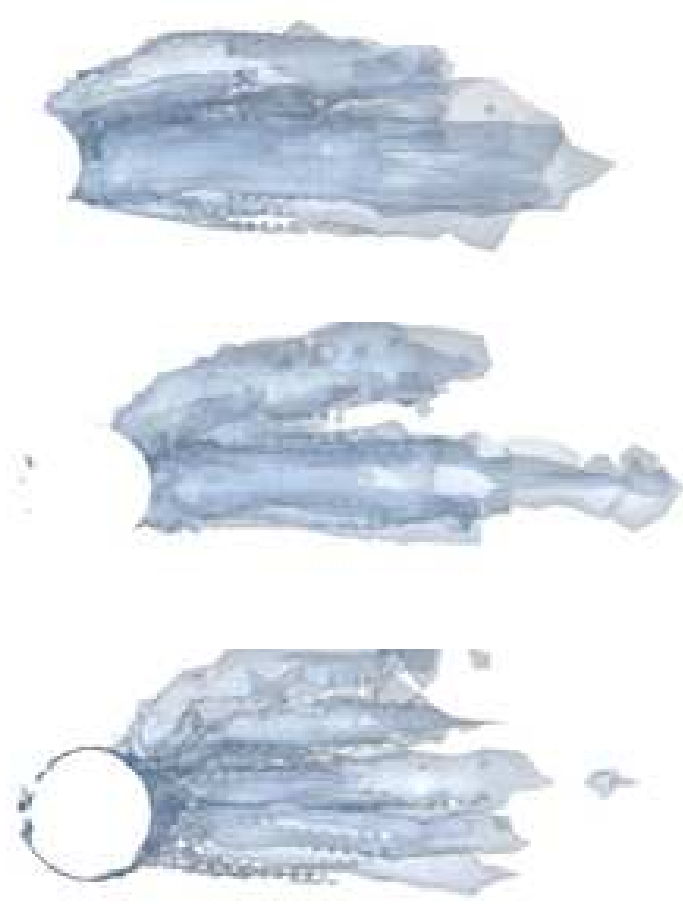
Vorticity: $\omega_1, \omega_2, \omega_3$: $t=1.75$: $c_D = 0.36$



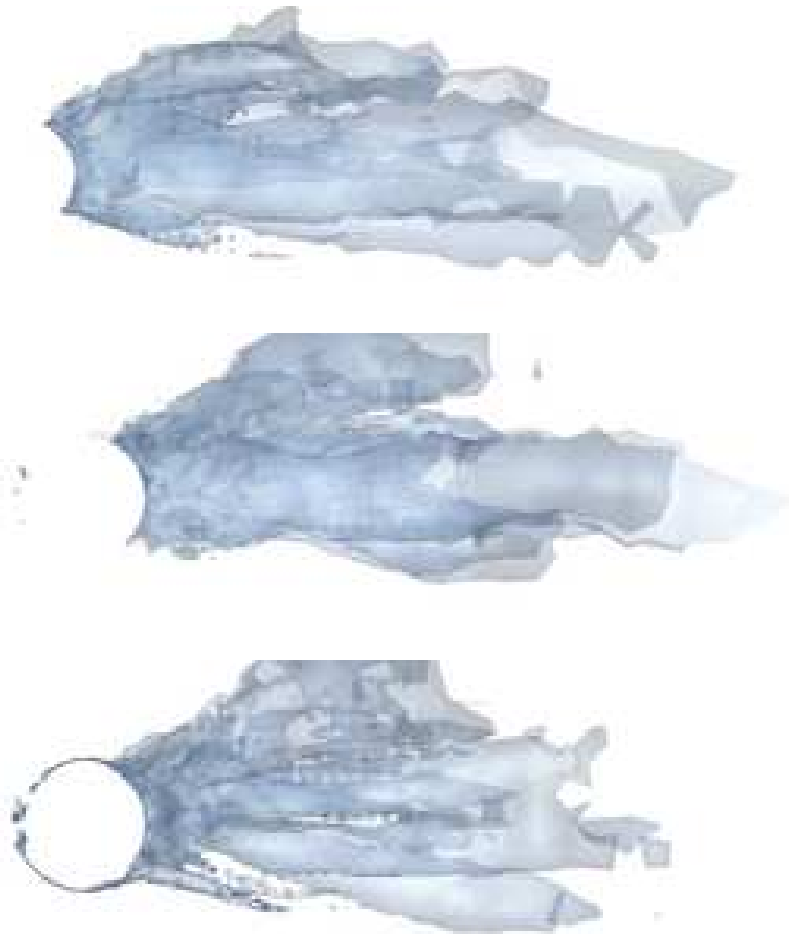


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Vorticity: $\omega_1, \omega_2, \omega_3$: $t=2.25$: $c_D = 0.78$

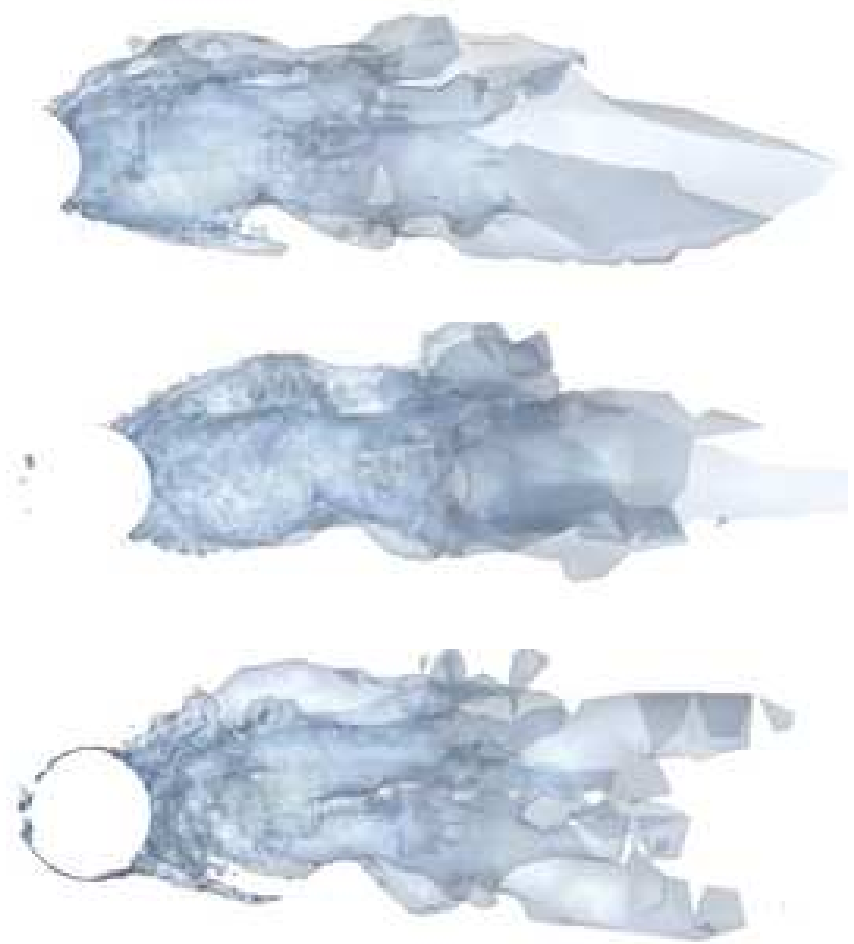


Vorticity: $\omega_1, \omega_2, \omega_3$: $t=2.5$: $c_D = 1.14$



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Vorticity: $\omega_1, \omega_2, \omega_3$: $t=2.75$: $c_D = 1.04$



Vorticity: $\omega_1, \omega_2, \omega_3$: $t=1.0$: $c_D = 0.22$

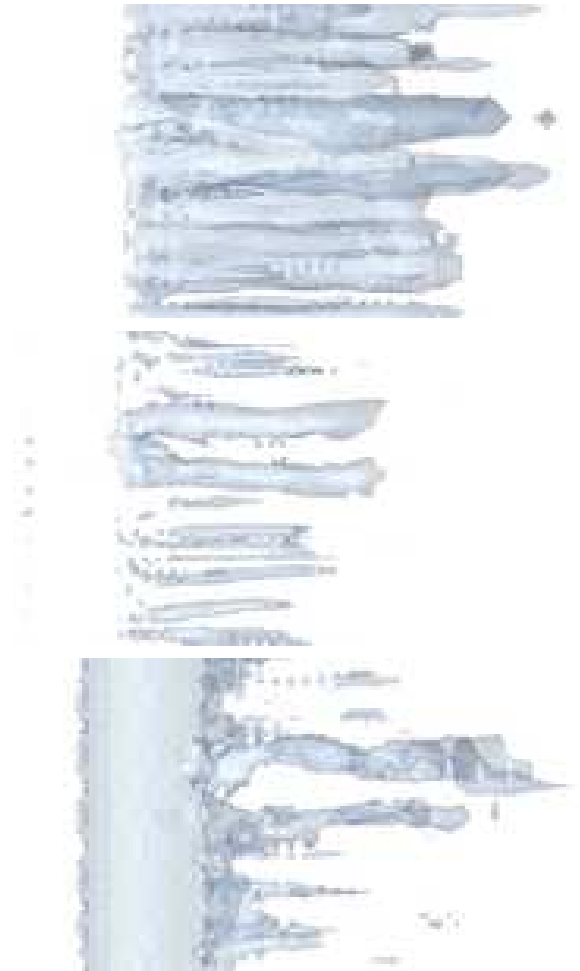


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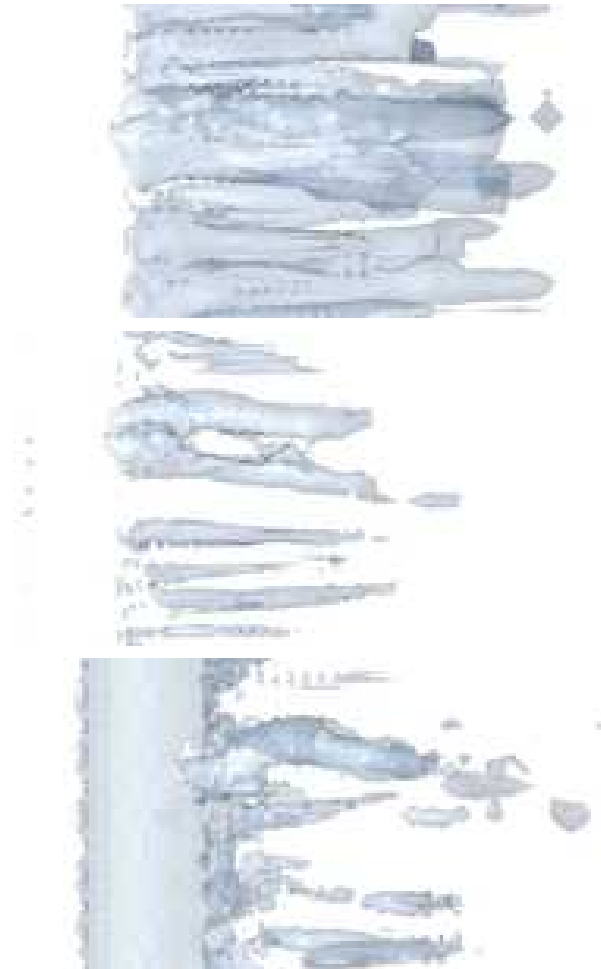
Vorticity: $\omega_1, \omega_2, \omega_3$: $t=1.25$: $c_D = 0.25$



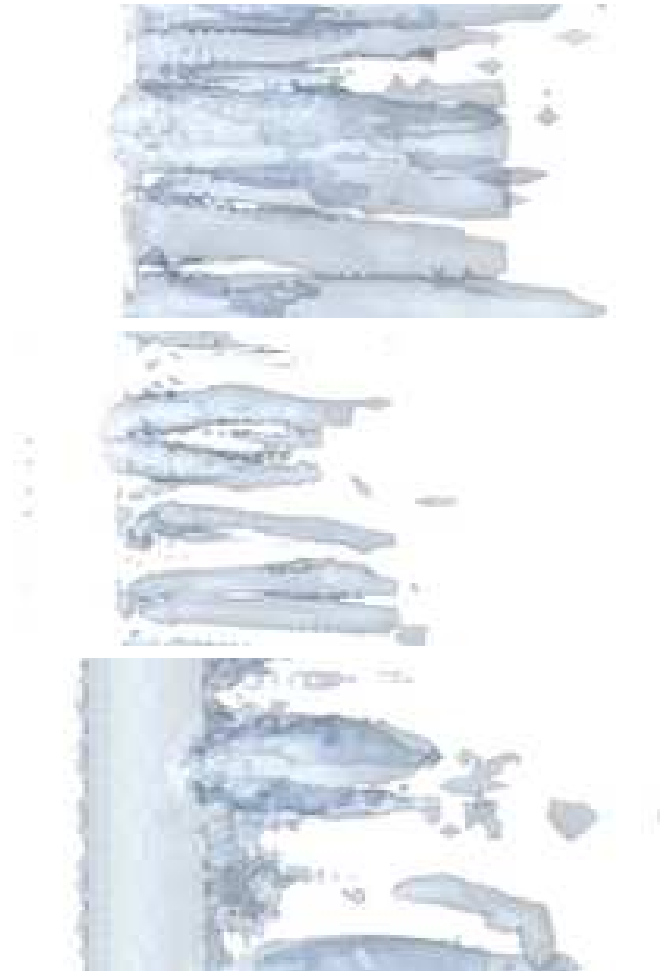
Vorticity: $\omega_1, \omega_2, \omega_3$: $t=1.5$: $c_D = 0.28$



Vorticity: $\omega_1, \omega_2, \omega_3$: $t=1.75$: $c_D = 0.36$

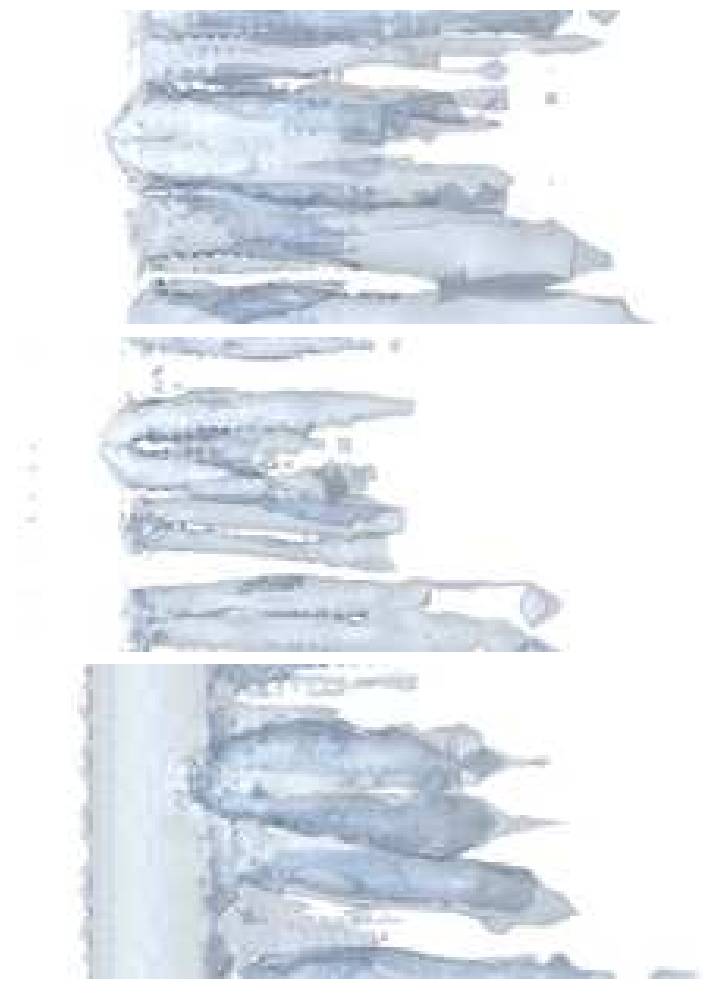


Vorticity: $\omega_1, \omega_2, \omega_3$: $t=2.0$: $c_D = 0.51$



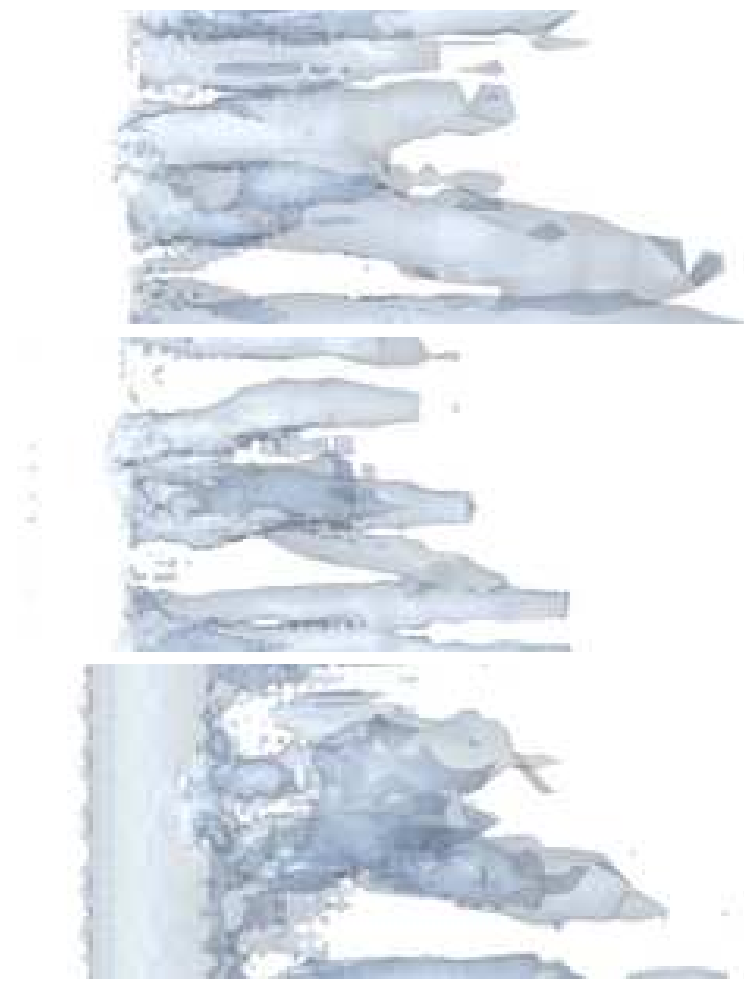
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Vorticity: $\omega_1, \omega_2, \omega_3$: $t=2.25$: $c_D = 0.78$



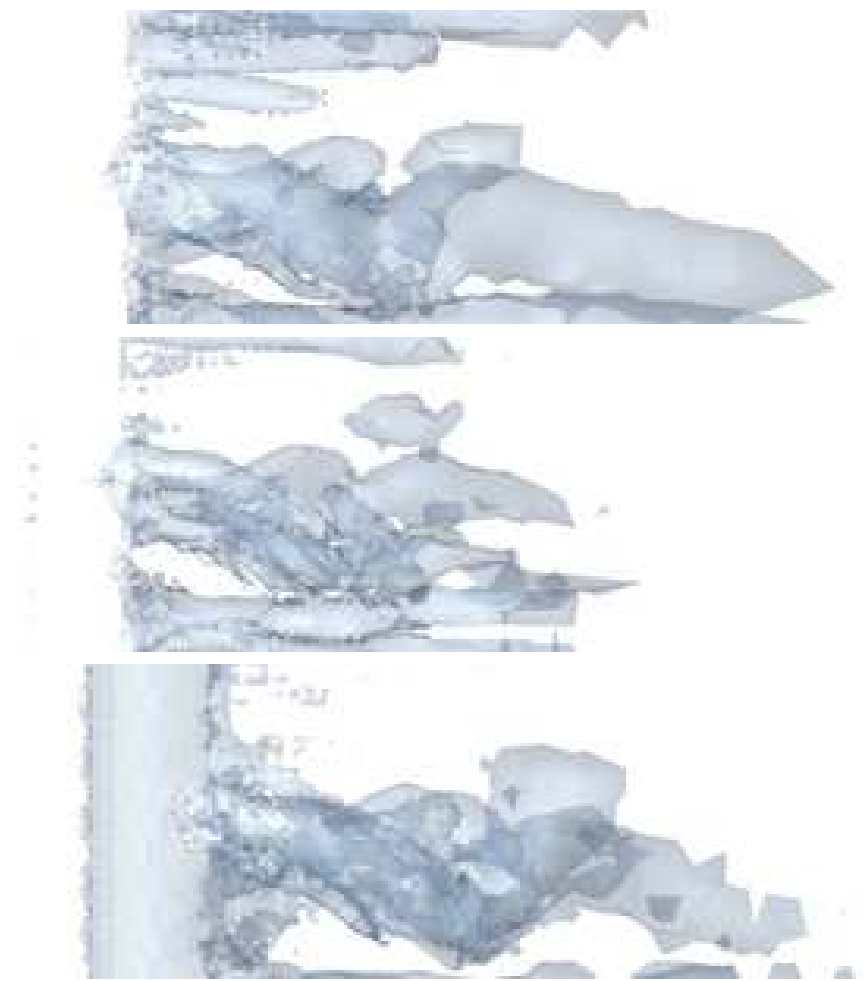
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Vorticity: $\omega_1, \omega_2, \omega_3$: $t=2.5$: $c_D = 1.14$

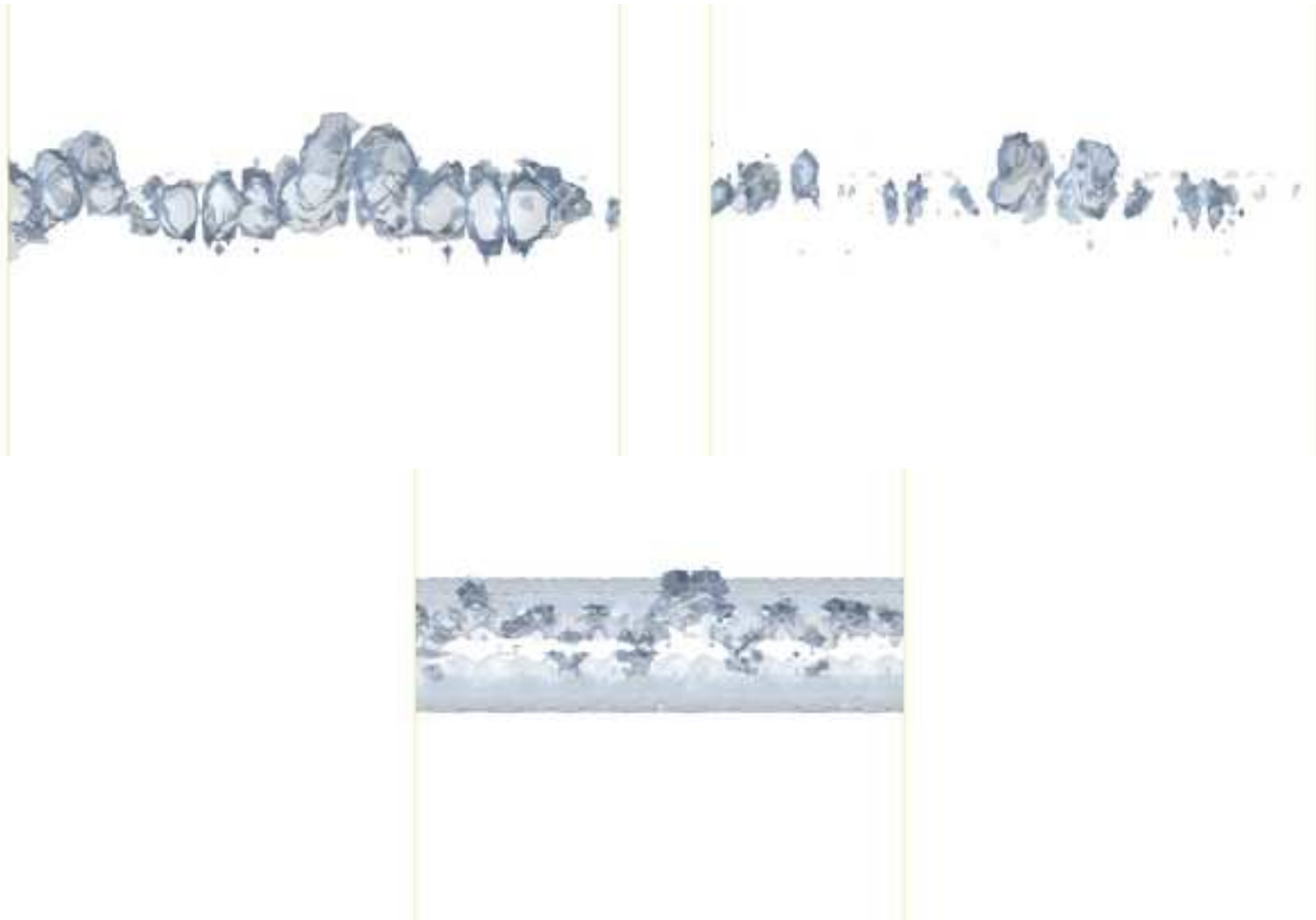


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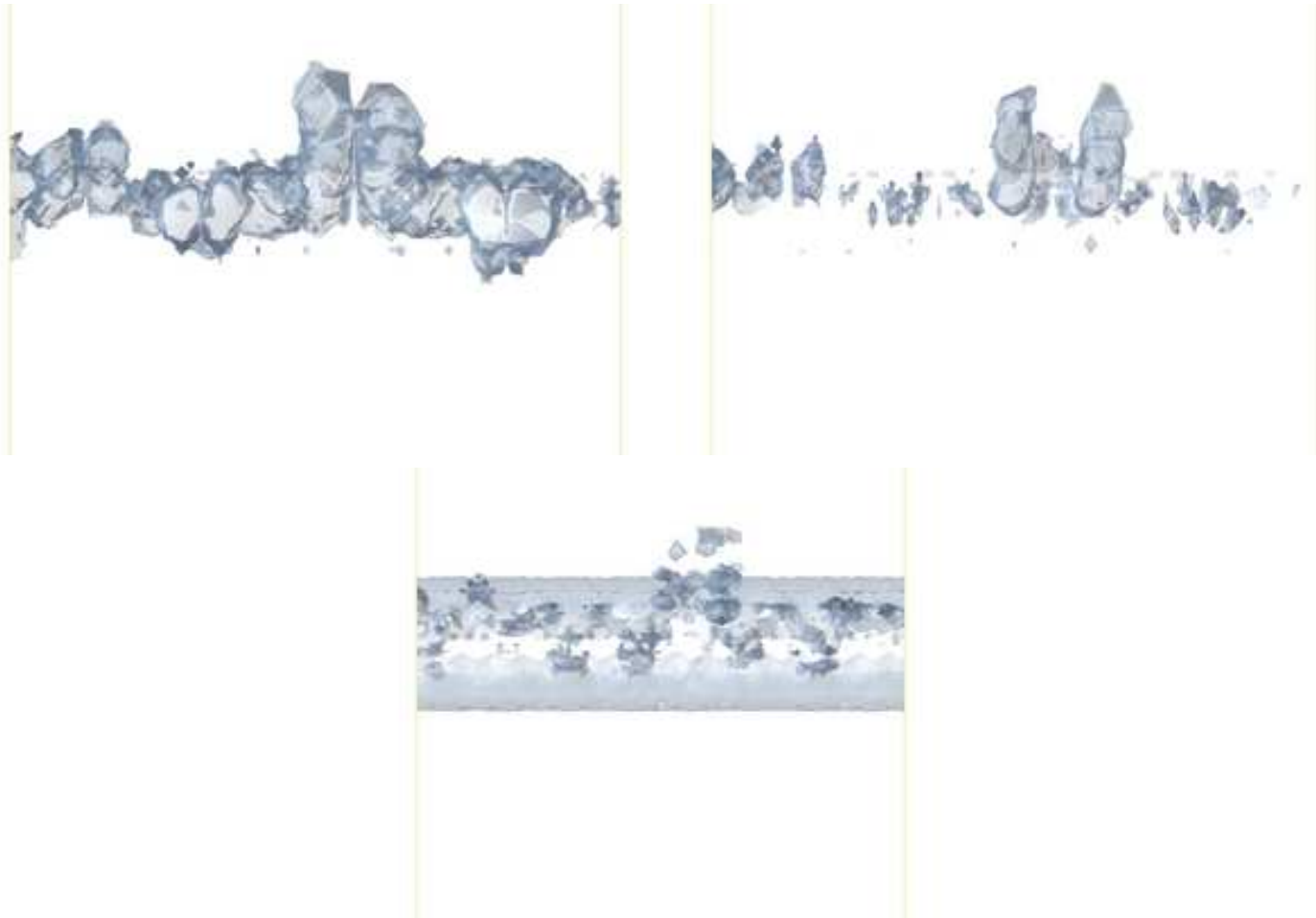
Vorticity: $\omega_1, \omega_2, \omega_3$: $t=2.75$: $c_D = 1.04$



Vorticity: $\omega_1, \omega_2, \omega_3$: $t=1.0$: $c_D = 0.22$

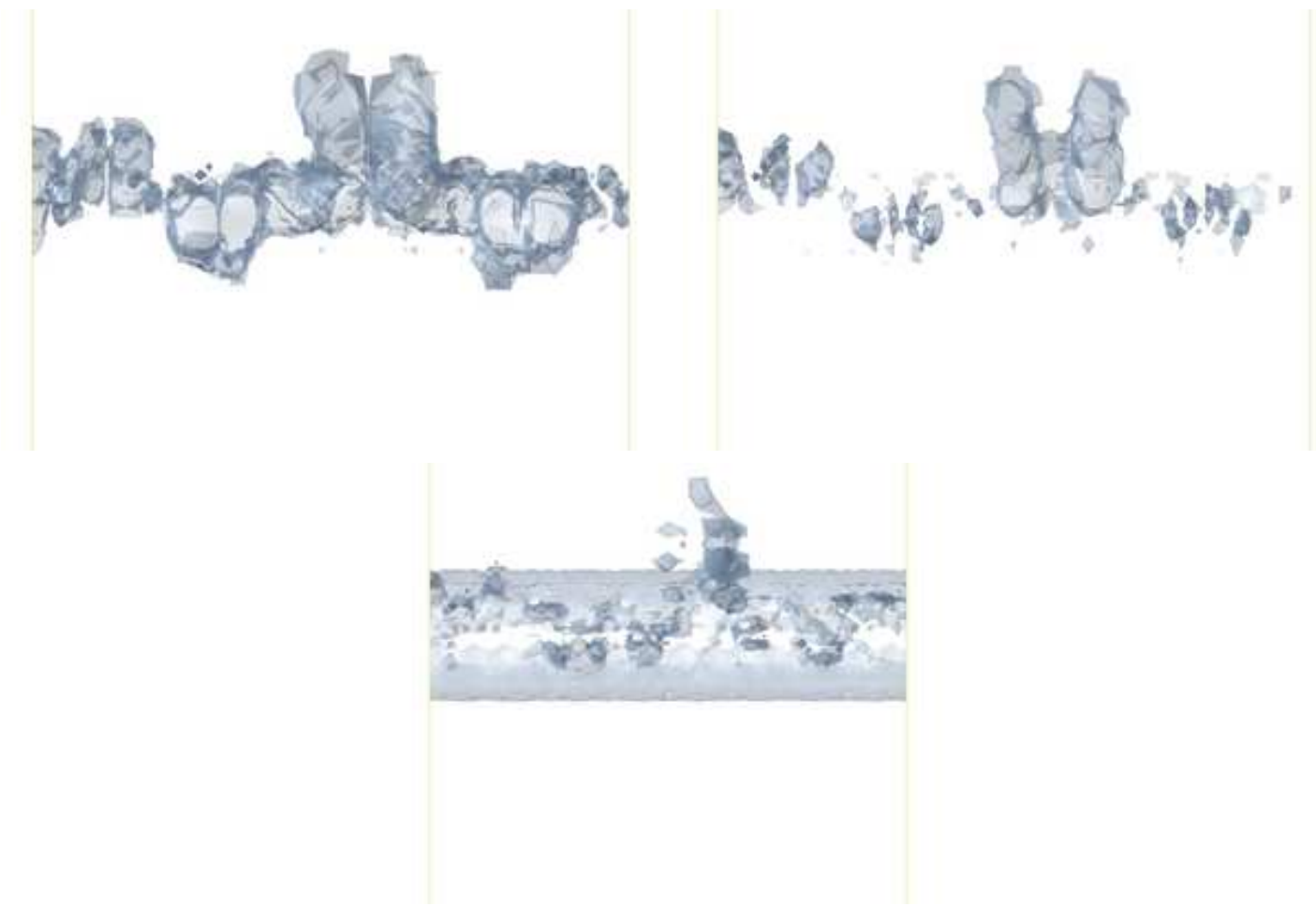


Vorticity: $\omega_1, \omega_2, \omega_3$: $t=1.25$: $c_D = 0.25$



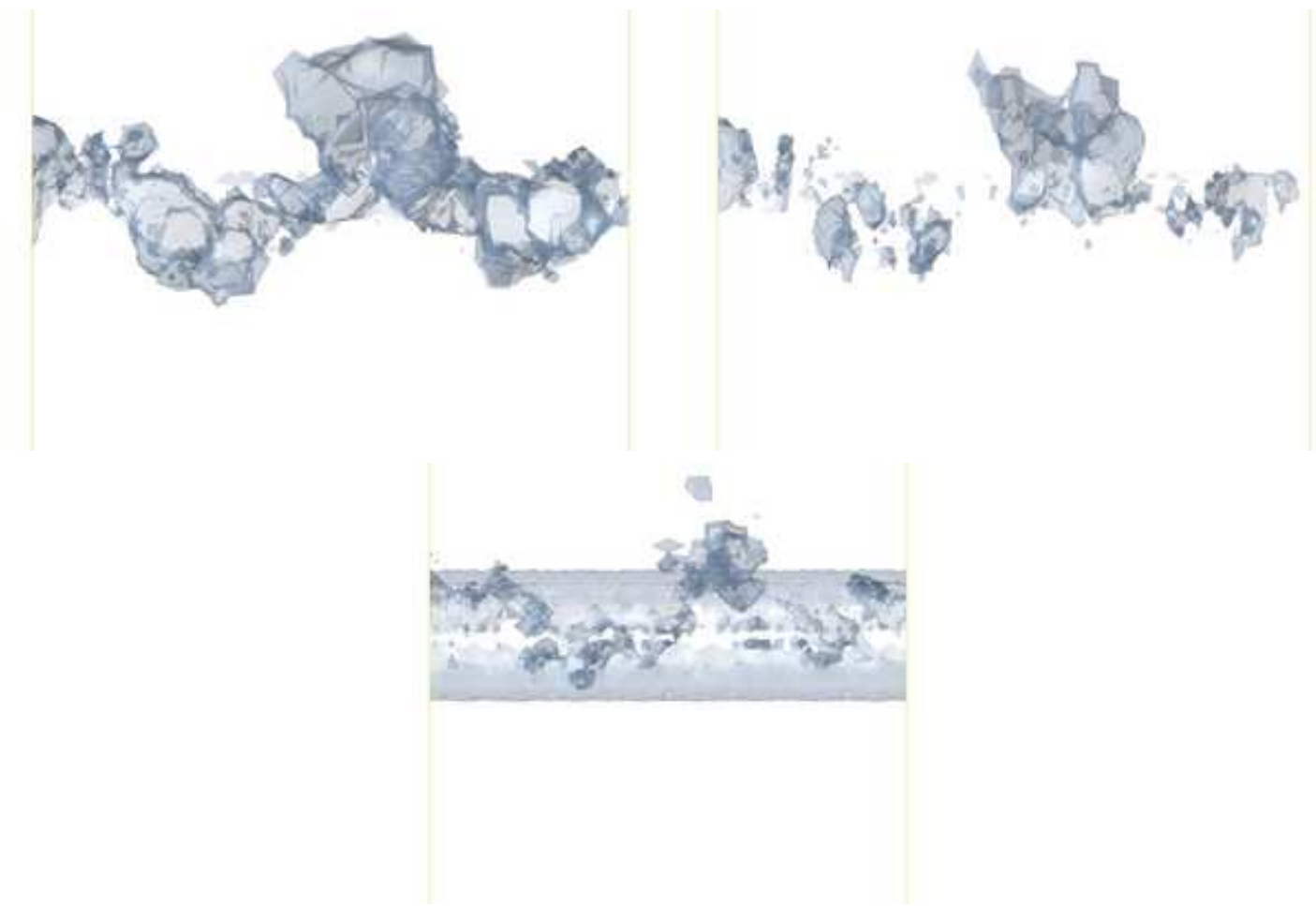
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Vorticity: $\omega_1, \omega_2, \omega_3$: $t=1.5$: $c_D = 0.28$



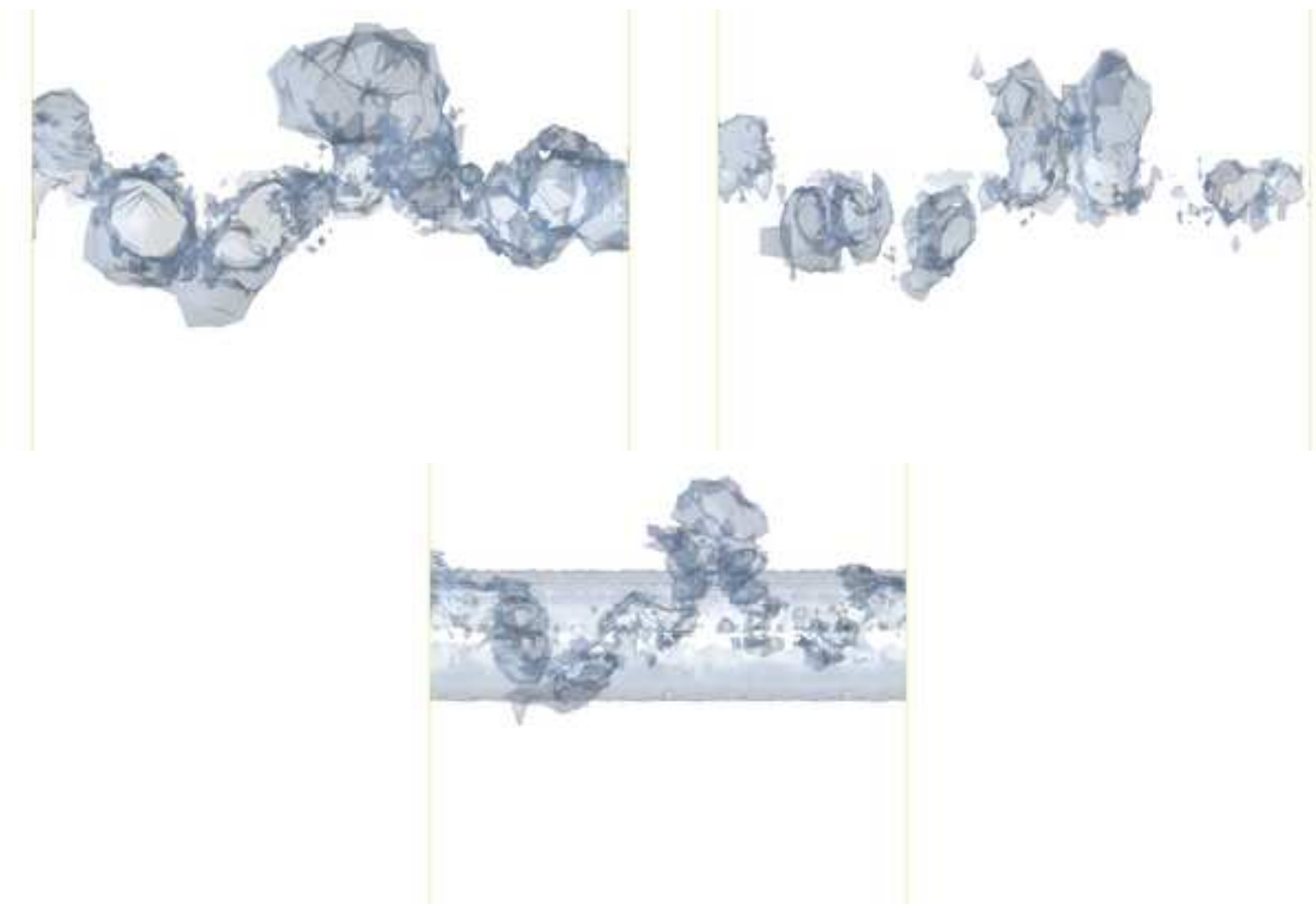
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Vorticity: $\omega_1, \omega_2, \omega_3$: $t=1.75$: $c_D = 0.36$



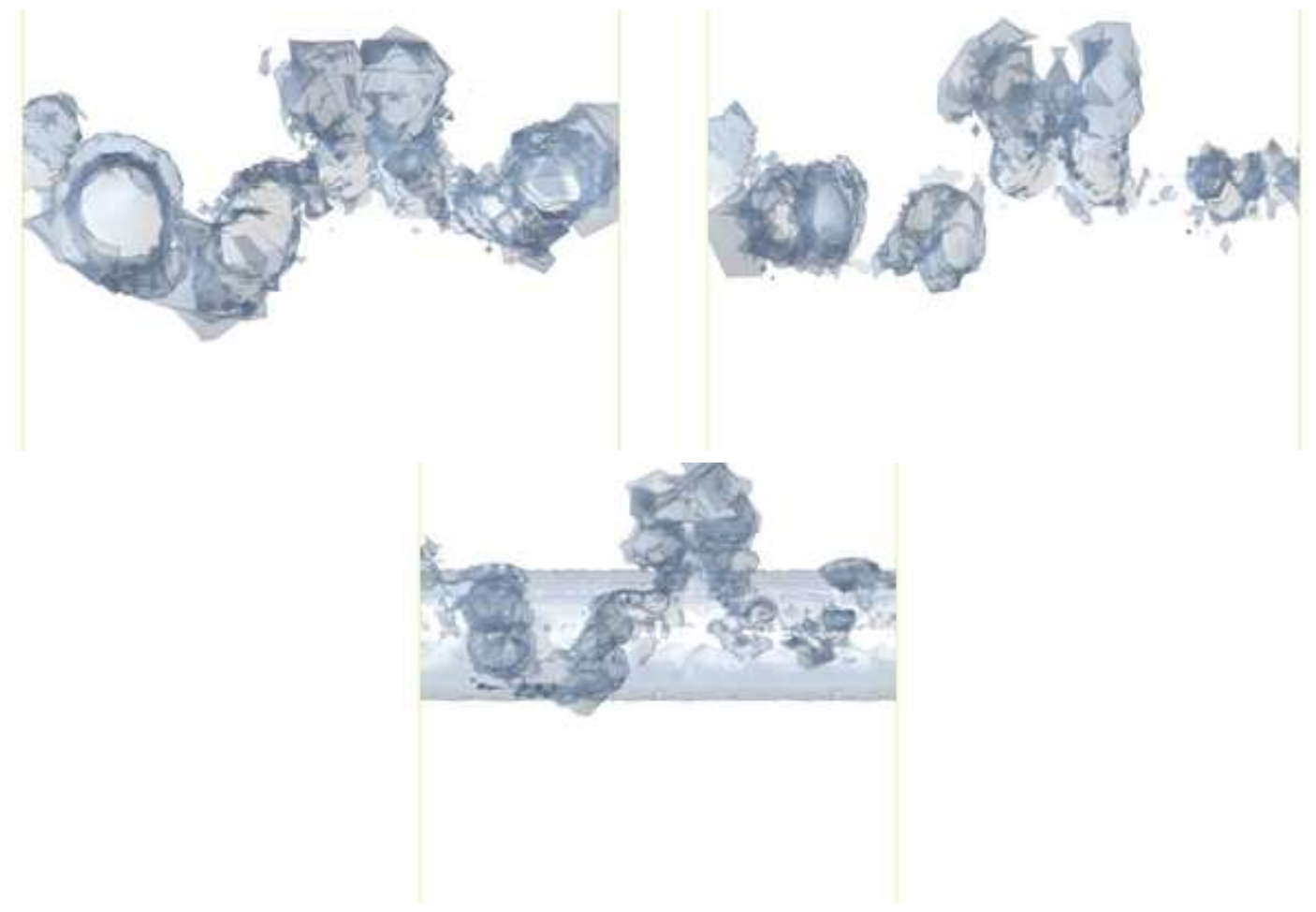
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Vorticity: $\omega_1, \omega_2, \omega_3$: $t=2.0$: $c_D = 0.51$



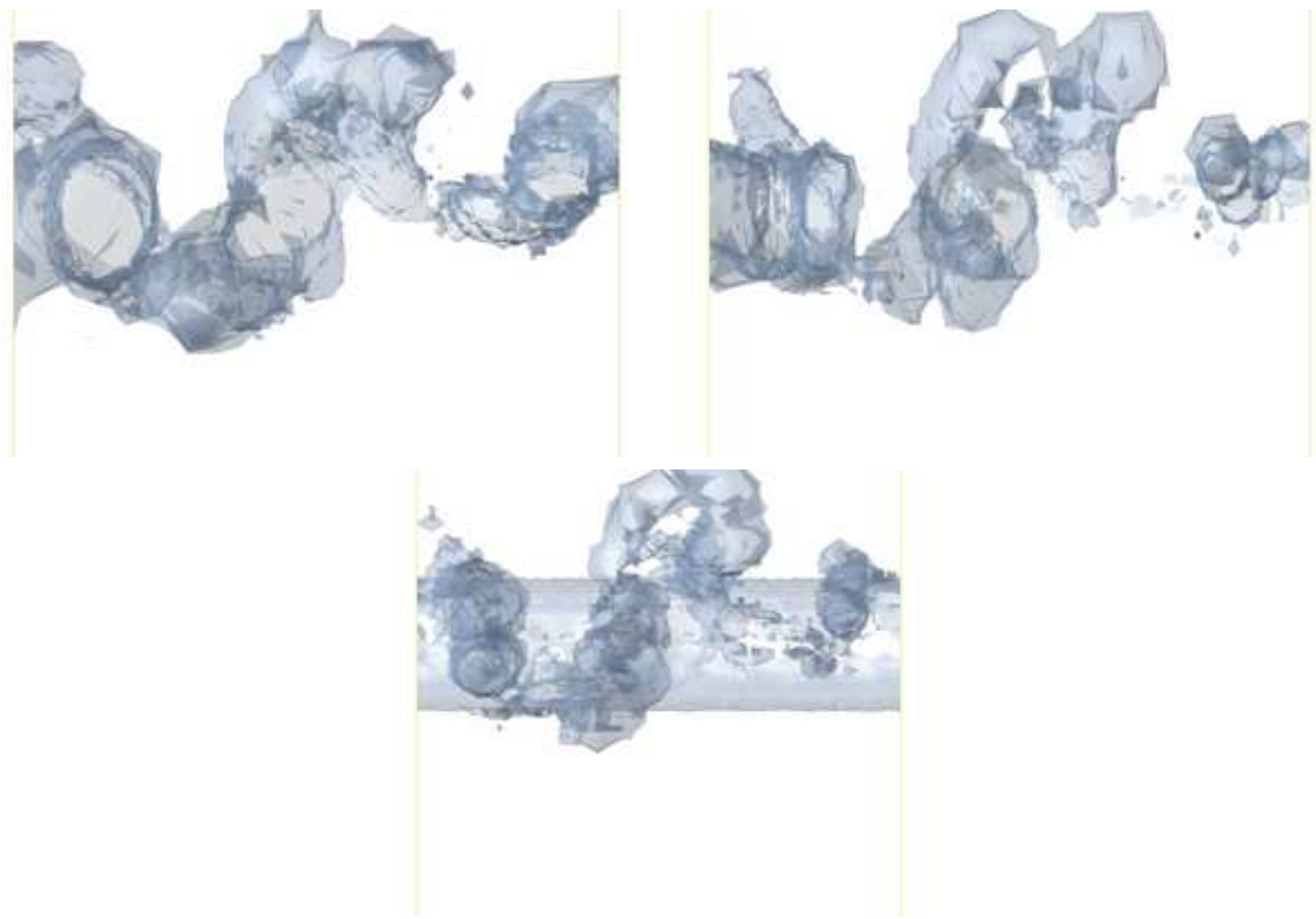
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Vorticity: $\omega_1, \omega_2, \omega_3$: $t=2.25$: $c_D = 0.78$



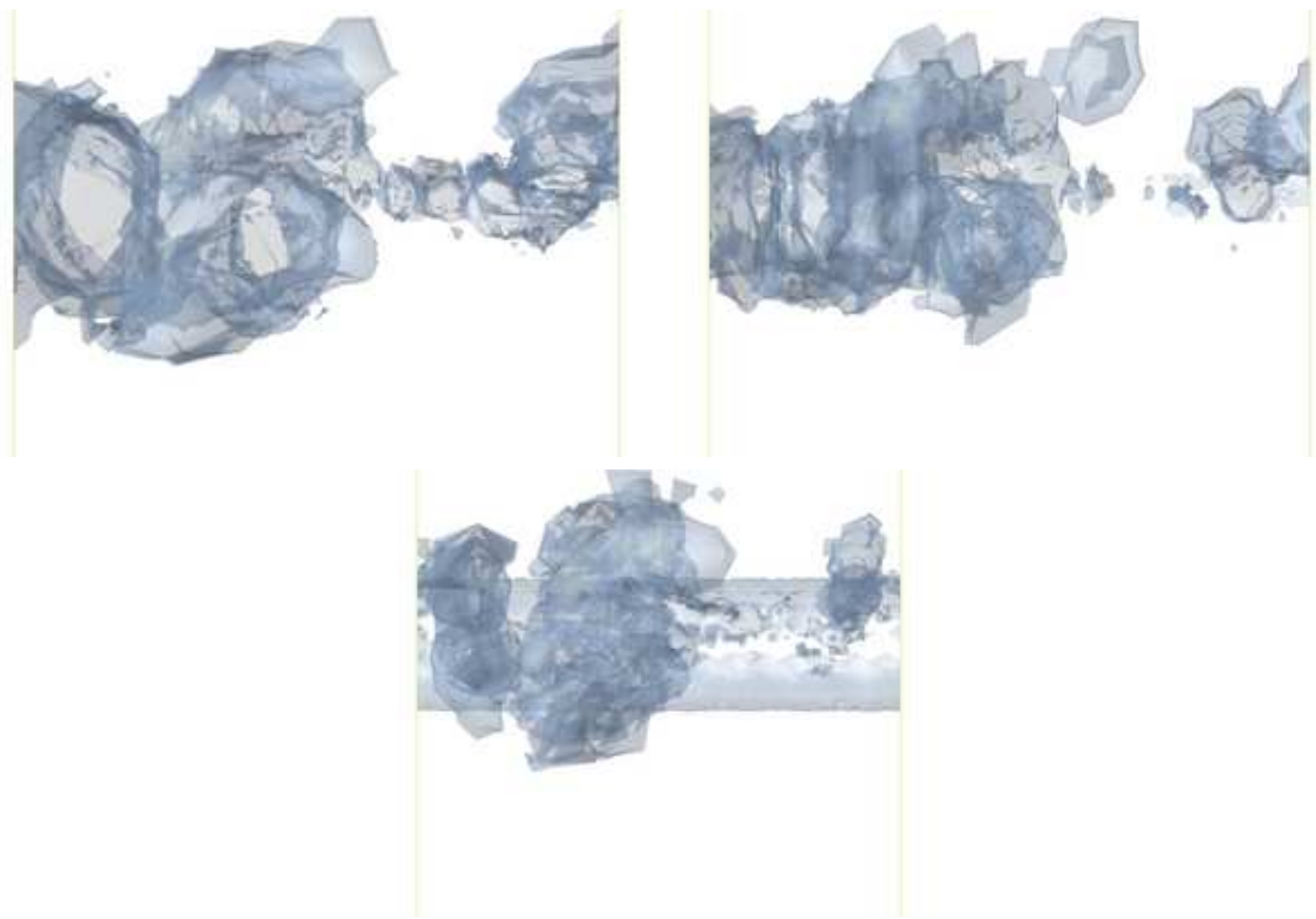
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Vorticity: $\omega_1, \omega_2, \omega_3$: $t=2.5$: $c_D = 1.14$

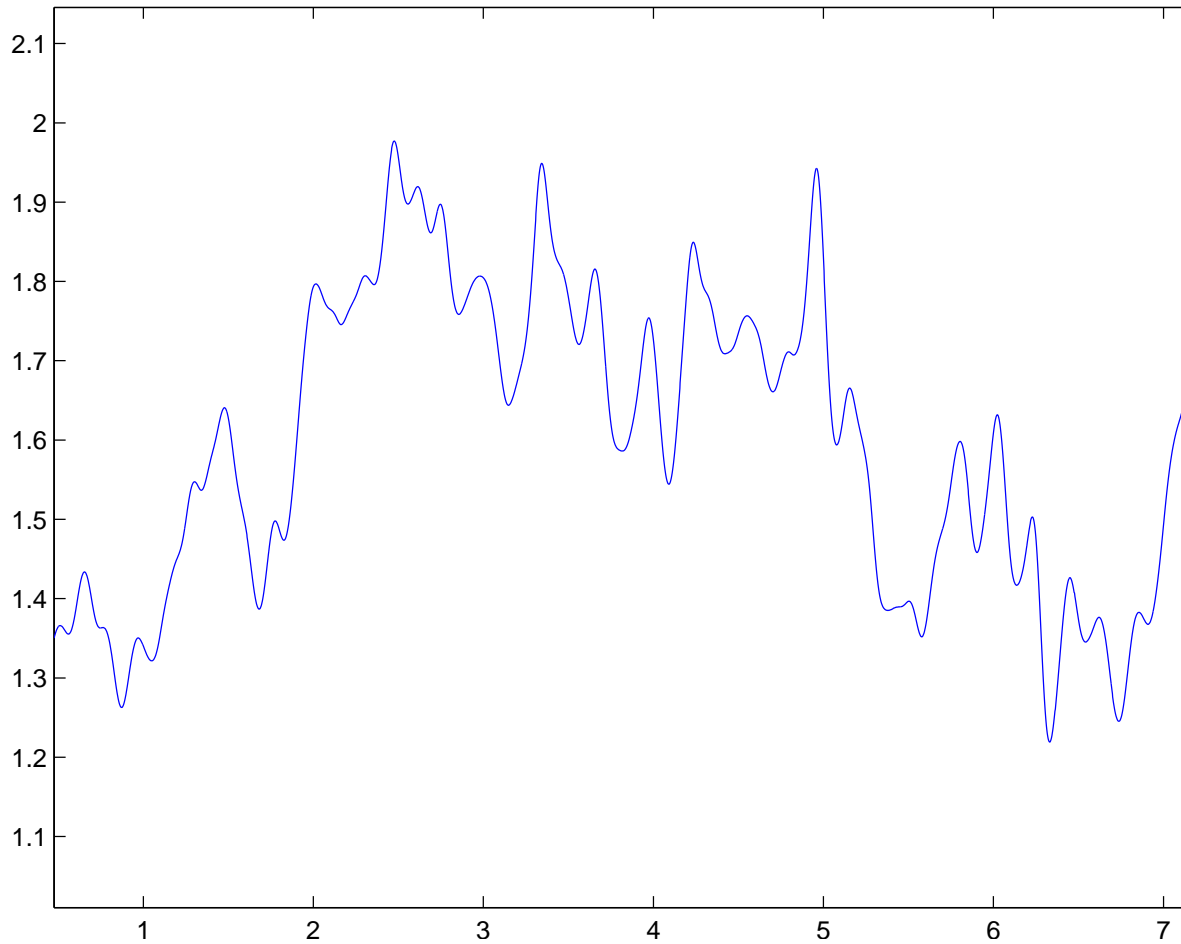


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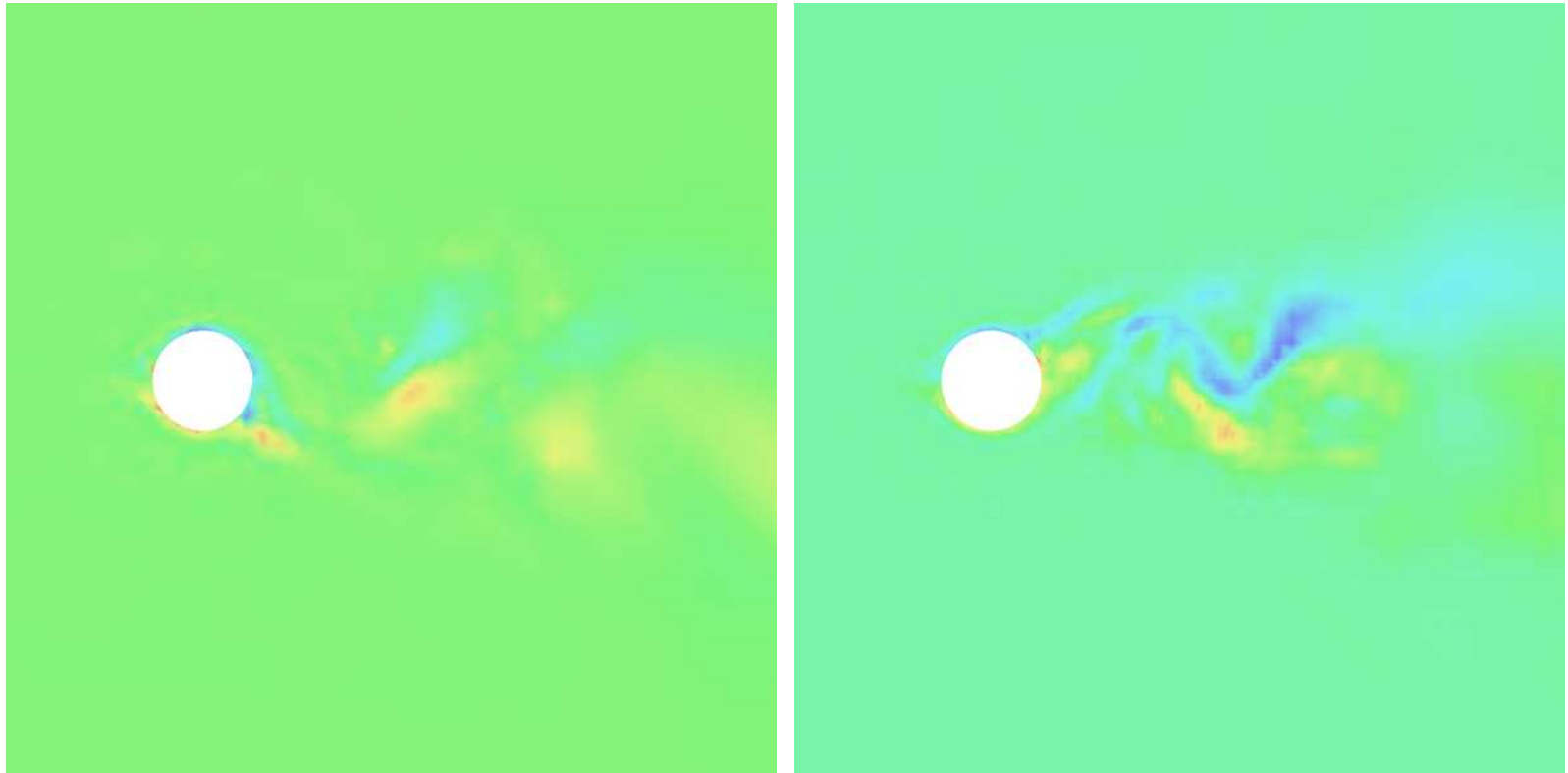
Vorticity: $\omega_1, \omega_2, \omega_3$: $t=2.75$: $c_D = 1.04$



Oscillating EG2 solution: c_D



Oscillating EG2 solution



Note: there is only one separation point (in each section)

Guadalupe, Canaries: $Re \approx 10^9$



Note: only one separation point (consistent with EG2)

EG2 and Turbulent Euler solutions

- No experimental results for cylinder at $Re > 10^7$
- Therefore not much is known about this fbw regime.
- For geophysical fbw ($Re \approx 10^9$) we observe “von Karman vortex shedding”; but with separation in one point only (no wake!) consistent with the EG2 solution.
- EG2: no empirical parameters (only parameter is h)

“In a reasonable theory there are no dimensionless numbers whose values are only empirically determinable.”
(Einstein)

Summary: NSE & G2

- Approximate weak solutions: existence by G2
- Weak uniqueness: duality: computation of $S_\epsilon(\hat{\psi})$
- Computational method: G2 (Adaptive DNS/LES)
 - Adaptivity \rightarrow very low computational cost
 - Generality \rightarrow quick & easy to adjust to new problem
 - Quantitative error control \rightarrow reliable results
- No filtering \Rightarrow no Reynolds stresses
- Turbulent boundary layer by simple friction b.c.
- EG2: parameter-free model of high Re flow