

# FEniCS Course

## Lecture 9: Incompressible Navier–Stokes equations

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# The incompressible Navier–Stokes equations

$$\begin{aligned} \dot{u} + u \cdot \nabla u - \nu \Delta u + \nabla p &= f && \text{in } \Omega \times (0, T] \\ \nabla \cdot u &= 0 && \text{in } \Omega \times (0, T] \\ u &= g_D && \text{on } \Gamma_D \times (0, T] \\ \nu \frac{\partial u}{\partial n} - pn &= g_N && \text{on } \Gamma_N \times (0, T] \\ u(\cdot, 0) &= u_0 && \text{in } \Omega \end{aligned}$$

- $u$  is the fluid velocity and  $p$  is the pressure divided by the density  $\rho$
- $\nu = \mu/\rho$  is the kinematic viscosity,  $\mu$  dynamic viscosity
- $f$  is a given body force per unit mass
- $g_D$  is a given boundary velocity
- $g_N$  is a given boundary function for the natural boundary condition
- $u_0$  is a given initial velocity

## Variational problem

Multiply the momentum equation by a test function  $v$  and integrate by parts:

$$\begin{aligned} \int_{\Omega} (\dot{u} + u \cdot \nabla u) \cdot v \, dx + \nu \int_{\Omega} \nabla u : \nabla u \, dx - \int_{\Omega} p \nabla \cdot v \, dx \\ = \int_{\Omega} f \cdot v \, dx + \int_{\Gamma_N} g_N \cdot v \, ds \end{aligned}$$

Short-hand notation:

$$(\dot{u} + u \cdot \nabla u, v) + \nu(\nabla u, \nabla v) - (p, \nabla \cdot v) = (f, v) + (g_N, v)_{\Gamma_N}$$

Multiply the continuity equation by a test function  $q$  and sum up: find  $(u, p) \in V$  such that

$$\begin{aligned} (\dot{u} + u \cdot \nabla u, v) + \nu(\nabla u, \nabla v) - (p, \nabla \cdot v) - (q, \nabla \cdot u) \\ = (f, v) + (g_N, v)_{\Gamma_N} \end{aligned}$$

for all  $(v, q) \in \hat{V}$

# Discrete mixed variational form of Navier–Stokes

Time-discretization leads to a *saddle-point* problem on each time step:

$$\begin{bmatrix} M + \Delta t A + \Delta t N(U) & \Delta t B \\ \Delta t B^\top & 0 \end{bmatrix} \begin{bmatrix} U \\ P \end{bmatrix} = \begin{bmatrix} b \\ 0 \end{bmatrix}$$

- Efficient solution of the saddle-point problem relies on the efficiency of special-purpose preconditioners (Uzawa iteration, Schur complement preconditioners, ...)
- We will use another approach (simpler and often more efficient)

# The classical Chorin-Teman projection method

**Step 1:** Compute *tentative velocity*  $u^\star$  solving

$$\begin{aligned}\frac{u^\star - u^n}{\Delta t} - \nu \Delta u^\star + (u^\star \cdot \nabla) u^\star &= f^{n+1} && \text{in } \Omega \\ u^\star &= g_D && \text{on } \Omega_D \\ \frac{\partial u^\star}{\partial n} &= 0 && \text{on } \Omega_N\end{aligned}$$

**Step 2:** Compute a *corrected* velocity  $u^{n+1}$  and a *new* pressure  $p^{n+1}$  solving

$$\begin{aligned}\frac{u^{n+1} - u^\star}{\Delta t} + \nabla p^{n+1} &= 0 && \text{in } \Omega \\ \nabla \cdot u^{n+1} &= 0 && \text{in } \Omega \\ u^{n+1} \cdot n &= 0 && \text{on } \partial\Omega\end{aligned}$$

## Computing the tentative velocity

In principle, the term  $(u^* \cdot \nabla)u^{**}$  can be approximated in several ways

- Explicit:  $u^* = u^{**} = u^n \Rightarrow$  diffusion-reaction equation
- Semi-implicit  $u^* = u^n$  and  $u^{**} = u^{n+1} \Rightarrow$  convection-diffusion-reaction equation
- Fully-implicit  $u^* = u^{**} = u^{n+1}$  retaining the basic non-linearity in the Navier-Stokes equations

The natural outflow condition  $\nu \partial_n u - pn = 0$  is artificially enforced by requiring

- $\partial_n u^\star = 0$  on  $\partial\Omega_N$  in step 1
- $p^{n+1} = 0$  on  $\partial\Omega_N$  in step 2

## Solving the projection step

Applying  $\nabla \cdot$  to  $\frac{u^{n+1} - u^\star}{\Delta t} + \nabla p^{n+1} = 0$  and using requirement  $\nabla \cdot u^{n+1} = 0$  yields

$$\Delta p^{n+1} = \frac{1}{\Delta t} \nabla \cdot u^\star \quad \text{in } \Omega$$

We already required

$$p = 0 \quad \text{on } \partial\Omega_N$$

Multiplying  $\frac{u^{n+1} - u^\star}{\Delta t} + \nabla p^{n+1} = 0$  with  $n$  and restricting to  $\partial\Omega_D$  gives

$$\frac{\partial p^{n+1}}{\partial n} = 0 \quad \text{on } \partial\Omega_D$$

Compute  $u^{n+1}$  by

$$u^{n+1} = u^\star - \Delta t \nabla p^{n+1}$$

including boundary conditions for  $u$  at  $t = t^{n+1}$

## Chorin-Teman projection method – Summary

- ❶ Compute tentative velocity  $u^\star$  by

$$\left(\frac{u^\star - u^n}{\Delta t}, v\right) + ((u^* \cdot \nabla)u^{**}, v) + \nu(\nabla u^\star, \nabla v) - (f, v) = 0$$

including boundary conditions for the velocity.

- ❷ Compute new pressure  $p^{n+1}$  by

$$(\nabla p^{n+1}, \nabla q) + \frac{1}{\Delta t}(\nabla \cdot u^\star, q) = 0$$

including boundary conditions for the pressure.

- ❸ Compute corrected velocity by

$$(u^{n+1} - u^\star, v) + \Delta t(\nabla p^{n+1}, v) = 0$$

including boundary conditions for the velocity.



# Useful FEniCS tools (I)

Note grad vs.  $\nabla$ :

*Python code*

```
dot(grad(u), u)
dot(u, nabla_grad(u))
```

Solving linear systems:

*Python code*

```
solve(A, x, b)
solve(A, x, b, "gmres", "ilu")
solve(A, x, b, "cg", "amg")
```

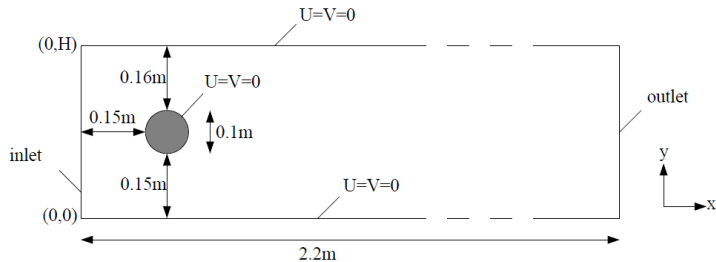
Defining  $a$  and  $L$  based on residual formulation:

*Python code*

```
F1 = ( (1/k)*inner(u - u0, v) + inner(grad(u0)*u0, v)
+ nu*inner(grad(u), grad(v)) - inner(f, v) ) * dx
a1 = lhs(F1)
L1 = rhs(F1)
```

# The FEniCS challenge!

Implement a famous benchmark simulating a laminar flow around a cylinder. The geometry is described by

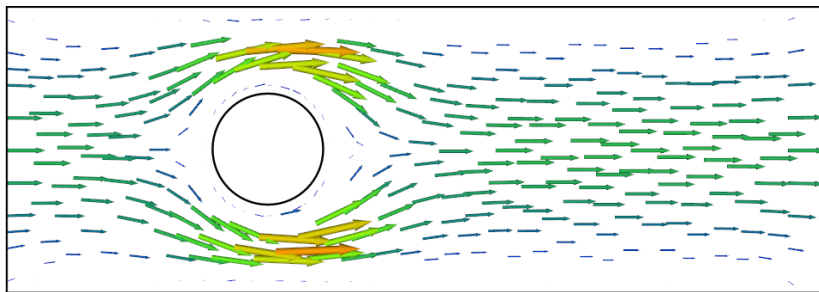


Set the kinematic viscosity  $\nu = 0.001 \text{ m}^2/\text{s}$  and  $\rho = 1.0 \text{ kg}/\text{m}^3$ . A “do-nothing” boundary condition is assumed at the outlet. Defining  $U_m = 1.5 \text{ m}/\text{s}$ , the time-dependent inflow condition is given by

$$U = 4U_m y(H - y) \sin(\pi t/8)/H^2, \quad V = 0.$$

## The FEniCS challenge!

The inflow boundary lies at  $x = -0.2$  and the outflow boundary at  $x = 2.0$ . Compute the flow on the time interval  $[0, 8]$  with time-step  $dt = 0.001$ . Test your implementation first for a larger time-step  $dt = 0.01$  and the same channel problem but with the cylinder removed. If everything goes fine you should get something like



Happy coding!